ALMA Cycle 6 Technical Handbook
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Revision History

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Chapter 1

Introduction

The Atacama Large Millimeter/Submillimeter Array (ALMA) is an aperture synthesis telescope consisting of 66 antennas that can be positioned in a number of different configurations. It operates over a broad range of observing frequencies in the millimeter and submillimeter regime of the electromagnetic spectrum.

ALMA Early Science Operations started with Cycle 0 in September 2011 and the official inauguration took place in March 2013.

Cycle 6 operations will include standard and non-standard observing modes. Standard observing modes are those well characterized, the observations can be calibrated, and in many cases, imaged by the ALMA data reduction pipeline. Non-standard modes are not as well characterized and may require manual calibration and imaging by ALMA staff. Up to 20% of the principal investigator (PI) observing time in Cycle 6 will be allocated to proposals requesting non-standard modes. Users should refer to Appendix A (ALMA Capabilities) in the ALMA Cycle 6 Proposer’s Guide, for the latest information and a description of standard and non-standard modes.

The Technical Handbook concentrates on the technical aspects of the Cycle 6 observing capabilities as well as on the hardware and software available for ALMA users. It should, however, not be necessary to use the Technical Handbook to prepare an ALMA proposal.

The Technical Handbook is divided into three main sections: The concepts of interferometry and the ALMA hardware components (Chapters 2–5), the observing concepts and software (Chapters 6–9), and finally the data quality and handling (Chapters 10–13). It also includes a number of appendices with expanded information on specific ALMA hardware components and calibration, and an acronym list widely used by ALMA staff.

Chapter 2 describes the ALMA Array components: The 12-m Array and the Atacama Compact Array (ACA), also known as the Morita Array which is comprised of the 7-m and Total Power (TP) Array. A general description of the different array elements and hardware components is provided.

Chapter 3 gives a brief introduction to interferometry, including a description on the concepts of basic radio astronomy and the principles of aperture synthesis.

Chapter 4 describes the details of the eight receiver bands offered for Cycle 6 observations. The general technical specifications and a brief explanation on local oscillators and intermediate frequency (IF) range is presented. Plots with the atmospheric transmission and the typical system temperatures per receiver band are also included.

Chapter 5 describes the correlators and the data processing taking place in these special purpose supercomputers. A description of the 64-input Correlator (used for the 12-m Array) and the ACA Correlator (used by the 7-m and the TP Arrays) is provided. Observing modes available for continuum and spectral line observations are presented.

Chapter 6 describes how the spectral setup is done in the correlators. A description of the signal path and local oscillator (LO) chain used between the frontends and the correlators is covered and how the hardware is
configured to define spectral setups for an observation.

Chapter 7 describes several aspects of imaging to consider in ALMA observations. A short description on the different configurations proposed for Cycle 6 is included. Concepts of shadowing, beam shape, and spatial scale filtering are revisited. Mosaicing and 12-m and 7-m Array data combination are also presented.

Chapter 8 describes the observing modes offered for Cycle 6 and the observing sequence of projects. Single field interferometry, mosaics, single-dish observations, polarization, multiple region modes, solar, VLBI, astrometry and ephemeris observations are detailed in this section.

Chapter 9 gives a brief overview how sensitivities and integration times are calculated at ALMA.

Chapter 10 describes how calibration is performed at ALMA, providing a description on how to calibrate long-term and short-term effects as well as how calibrators are selected.

Chapter 11 describes the data quality assurance process. A description of the criteria used for passing quality assurance as well as the pipeline heuristics used in data reduction and calibration which are used in determining the overall quality of the data are presented.

Chapter 12 describes the data flow and structure of ALMA data. A description of all the main software subsystems involved from data acquisition to archiving is provided as well as a description of the ALMA Science Data Model (ASDM) which defines the metadata structure adopted by ALMA.

Chapter 13 describes how the data are stored, the data flow and the user interface to the ALMA Archive.

The Technical Handbook concludes with three appendices, which contain supplemental material about the antenna design and ALMA transporter (Appendix A), the Local Oscillator (LO) and Intermediate Frequency (IF) system (Appendix B) and a list of acronyms widely used by ALMA staff (Appendix C).

Figure 1.1: ALMA antennas on the Chajnantor Plateau.
Credit: ALMA (ESO/NAOJ/NRAO), O. Dessibourg
Chapter 2

Array Components

This chapter describes the main characteristics of each ALMA array. Unless otherwise noted, the description is appropriate for the fully completed ALMA telescope.

2.1 The ALMA Telescope

ALMA is composed of 66 high-precision antennas. Fifty of these antennas are 12 meter antennas in the 12-m Array, used for sensitive, high-resolution imaging. These fifty 12 m antennas are complemented by the Atacama Compact Array (ACA), also known as the Morita Array\(^1\), composed of twelve closely spaced 7 m antennas (the 7-m Array), and four 12 m antennas for single-dish (or Total Power) observations (the TP Array), to enhance wide-field imaging of extended structures. In Cycle 6, ALMA will cover most of the wavelength range from 8.5 to 0.32 mm (35–950 GHz), and when ALMA is completed the coverage will be 50 GHz.

The array is located on the Chajnantor plain of the Chilean Andes (lat.=–23.02917°, long.=–67.754649°), a site that normally offers the exceptionally dry and clear sky conditions required to observe at millimeter and sub-millimeter wavelengths\(^2\). The ALMA antennas, weather stations, the two correlators and their computer interfaces, Local Oscillator generation hardware, timekeeping hardware, and the related array Real-Time Machine computer are all located at the 5000-meter altitude site referred to as the Array Operations Site (AOS). This site is connected via Gigabit fiber links to the Operation Support Facility (OSF), located at an altitude of 2900 meters, about 22 km from the AOS and 40 km from the town of San Pedro de Atacama. Science operations are conducted from the OSF and coordinated from the JAO Central office in Santiago. All three ALMA arrays are controlled via control software developed on the ALMA Common Software\(^3\) (ACS).

There are 192 antenna foundations (stations) distributed over the Chajnantor and Pampa la Bola plateaus. The antenna foundation distribution yields baselines (distances between two antennas) ranging from 15 m to ~16 km, which are crucial in determining the image quality and spatial resolution of ALMA (see Chapter 7). The antenna foundations provide the stiffness required for precise antenna pointing, as well as electrical power and digital connectivity to the main AOS building (See Appendix A.2). The antennas can be re-configured into the different array configurations (Chapter 7) using the two special purpose ALMA antenna transporters (see Appendix A.3).

The number of antennas in each array component (12-m, 7-m and TP Arrays), and the specific configurations available for an observing season (e.g. Cycle 6) will be published in the Capabilities section of the document ALMA Proposer’s Guide. Complementary background information on ALMA and its capabilities for Cycle 6 can be found in the document A Primer for Cycle 6. Both documents can be found on the link http://www.almascience.org/documents-and-tools/.

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\(^1\)dedicated to the honour of K.-I. Morita

\(^2\)http://www.almascience.org/about-alma/weather

\(^3\)http://www.eso.org/projects/alma/develop/acs/
CHAPTER 2. ARRAY COMPONENTS

Figure 2.1: The ALMA 12-m Array in its compact configuration (left hand side of the image). The ACA with all 7 m antennas (dashed orange circles) and four single-dish 12 m antennas (blue circles) are distributed in the right hand side of the image. A few unoccupied stations can be seen, to which antennas of the 12-m Array can be moved by the transporter as the array is being reconfigured. At its most extended configuration, antennas in the 12-m Array will be about 16 km apart.

2.2 The 12-m Array

The 12-m Array consists of fifty 12 m diameter antennas designed and built by the European and North American ALMA partners (each providing 25 units), according to the stringent ALMA Antenna Performance specifications (see Appendix A). Each antenna contains one front end, including a cryostat (see Appendix A.4), amplitude calibration device (ACD; A.5), water vapor radiometer (WVR; A.6), and back-end electronics (analog and digital racks). The WVRs are used to correct the phase fluctuations caused by water in the atmosphere along the line of sight of each 12-m Array element. The cryostat can contain up to ten cartridges, each covering one frequency band (see Chapter 4). Only one band observes at any time, but up to three can be switched on simultaneously, and rapid switching between them is possible currently only for observatory calibration (see Chapter 4). Each receiver band (Chapter 4) detects two orthogonal linear polarizations and down-converts the signals to an intermediate frequency with eight GHz of bandwidth per polarization. Bands 3-8 cartridges are dual sideband (2SB) and Bands 9 and 10 are double sideband (DSB).

The Local Oscillator (LO) signals (Appendix B) are transmitted to the antennas on optical fibers with a round trip measurement to correct for changes in the fiber length. There are four independent LO reference systems so that the 7-m Array, the TP array and two subsections of the 12-m Array (e.g., two ‘sub-arrays’) can conduct simultaneous independent observations. Please note that the sub-arrays feature of the 12-m Array is not yet a capability that is offered. The 8 GHz total IF bandwidth from the selected receiver is divided into four 2 GHz-wide basebands which are digitized at four Gsamples/s, with three-bit resolution, and transmitted on optical fibres. Total data rates are therefore 96 Gbits/s per antenna. With formatting, the bit rate is 120 Gbits/s.

On arrival at the central building, the data are recovered and processed in one of the two correlators: the 64-input Correlator and the ACA Correlator (see Chapter 5). All antennas can feed either correlator. The 64-input Correlator (Section 5.1) is normally used for the 12-m Array, but it can also take inputs from the 7-m Array or TP antennas. It is an XF correlator (cross-correlates first, then Fourier transforms), but the correlator proper is preceded by Tunable Filter Banks, which makes it a digital hybrid XF correlator or FXF. These can select sub-bands from the 2 GHz-wide basebands in a very flexible manner. From each of these, the correlator then generates 2016 cross-correlations and 64 auto-correlations (requiring $1.7 \times 10^{16}$ operations per second). Either 2-, 3- or 4-bit resolution is used, and the sampling can be Nyquist or twice Nyquist. See

\footnote{See Section 4.1 for the minimum instantaneous bandwidth available for PI science observations}
Table 6.1 for an example table of spectral setups. These correlated data are fed to a group of processors which do the transforms and carry out integration and data compression.

Both correlators have minimum dump rates of 16 ms for cross-correlation and 1 ms for auto-correlation, although these dump rates can only be achieved using a reduced number of channels to prevent exceeding the maximum transmission and storage rates. The systems are designed for a maximum data rate of 64 MB/s, although the mean data rate will be considerably less.

The 12-m Array configurations have been designed so that in the most extended configurations the spatial angular resolution will be as small as 5 milliarcseconds at 950 GHz.

2.3 The Atacama Compact Array

Using an interferometer to obtain images of extended or large-scale structures leads to the well known “zero spacing” problem. This problem arises from the constraint that, to avoid collisions, it is not possible to pack antennas closer than their diameter, leaving a hole in the distribution of baselines at short and zero baseline separations (corresponding to large angular structure). As a result, spatial information from baselines shorter than the closed-packing ratio is not recovered\(^5\). This problem has considerable impact on observations of extended objects, particularly those in which the emitted power is dominated by their large-scale structures.

To achieve high-fidelity imaging of sources with emission on angular scales larger than those corresponding to the minimum spacing of the 12-m Array (the “Maximum Recoverable Scale” for that array - see Section 7.6), ALMA has been designed to include the Atacama Compact Array (ACA), also known as the Morita Array.

The ACA is composed of twelve 7 m antennas for interferometry (the 7-m Array) and four 12 m antennas for single-dish observations (TP Array). The four single-dish antennas provide spatial information samples equivalent from 0 m up to 12 m spacings as auto-correlations. The 7-m Array samples baselines from 9 m to 30 m, bridging the baseline sampling gap between the 12-m Array and the TP Array. The number of array elements available is published for each observing cycle.

The ACA is controlled via control software developed on the ALMA Common Software (ACS) platform and is operated in a similar fashion to the 12-m Array. To achieve this unified operation, the ACA system is as compatible with the 12-m Array as possible at the level of hardware, interface, data, and observing modes. The standard observing modes for the TP Array include spectral line and continuum observations with raster or Lissajous on-the-fly (OTF) scans, or position switching. The raw time-series signals from the ACA antennas are processed in the ACA Correlator (see Section 5.2) to produce the cross-correlated and auto-correlated data.

2.3.1 The 7-m Array

The 7-m Array is composed of twelve 7 m diameter antennas designed and built by East Asia to the ALMA specifications (see Appendix A). Each antenna contains one front end, including a cryostat, amplitude calibration devices, and one back end. Unlike the antennas of the 12-m Array, the 7 m antennas do not contain Water Vapor Radiometers (WVRs). The 7 m antenna cryostats are fitted with receivers nearly identical to those on the 12 m antennas, with small differences in the warm optics. The Local Oscillator signals transmitted to the 7 m antennas are originated identically to the ones sent to the 12-m Array, i.e., from inside the AOS building.

The ACA Correlator is normally used for the ACA, and can work with two sub-arrays. It is an FX correlator (Fourier transform first, then cross-correlate) with 3–bit input and 4 bits in the correlation. The correlator generates 120 cross-correlations and 16 auto-correlations for each baseband. These are passed to a (special-purpose) data processing computer at up to ~0.6 Gbits/s per baseband.

Even in its most compact configuration the 12-m Array does have spacings smaller than 15 m. The array configuration of the 7-m Array is designed to fill in these shorter spacings from about 9 m to ~30 m (see Chapter 7). Each cartridge (see Chapter 4) receives two orthogonal polarizations.

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\(^5\)Strictly speaking, mosaicing with imaging using a joint-deconvolution algorithm allows the recovery of more spatial information than normal synthesis imaging, but the problem caused by absent short and zero spacing information still remains.
2.3.2  The TP Array

The TP Array can fill in baseline coverage from 0 m to about 12 m, complementing the 7-m and 12-m Array’s baseline coverage. It consists of four 12 m diameter antennas built by East Asia (Appendix A). The specifications of the TP antennas are almost identical to the ones for the 12-m Array. The TP antennas are located on stations surrounding the 7-m Array.

The TP Array is usually connected to the ACA Correlator, but its antennas can also be connected to the 64-input Correlator and used for cross-correlation. The ALMA Proposer’s Guide describes the observing modes and capabilities offered for the TP Array for each cycle.

Due to the poorer point-source sensitivity of the 7-m Array, during Full Operations, the TP Array will be routinely used in the calibration observations of the 7-m Array, but this is not yet implemented. Since the 7-m Array is quite compact, atmospheric phase fluctuations will be smaller than for the 12-m Array.

Figure 2.2: The Morita Array - In remembrance of Professor Koh-Ichiro Morita. Koh-ichiro Morita, a professor at the NAOJ Chile Observatory, was one of the world’s renowned scientists in the field of aperture synthesis. He made a great contribution to designing the configuration of 16 antennas composing the Atacama Compact Array (ACA) manufactured by Japan, as well as to realizing high-resolution and high-quality imaging at millimeter/submillimeter wavelengths to further enhance the performance of ALMA. The picture above shows Professor Koh-Ichiro Morita taken at his office in the Joint ALMA Observatory.
Chapter 3

Principles and Concepts of Interferometry

3.1 Introduction

Interferometry is the technique ALMA uses to obtain very high angular resolution observations of astronomical phenomena. Figure 3.1 shows examples of two other observatories, the IRAM Plateau de Bure Interferometer (now expanding as NOEMA) in France and the Submillimeter Array in Hawaii, which also use this technique. In this Chapter, the principles and concepts behind interferometry are described, so that ALMA users can plan and understand their observations better.

Interferometry involves the combination of signals received from the sky by two or more physically separated antennas. The signals are interfered, allowing a sky brightness distribution to be sampled on an angular scale smaller than possible with a single antenna. The interference modifies the angular sensitivity of the antennas to include a sinusoid of constructive and destructive nodes. In this sense, the only emission measured by the interferometer is that from the scale defined by the angular extent of the sinusoidal wavelength, equivalently, the "spatial frequency". This wavelength is inversely proportional to the projected distance between the two antennas. Each datum from the interferometer is called a visibility, and is a measure of the brightness of the emission on the angular scale sampled, i.e., related to the amplitude of the sinusoid, and the relative position of that brightness on the sky, i.e., related to the phase of the sinusoid.

A range of discrete angular scales can be sampled by including many pairs of antennas in an array. Importantly, by tracking a source across the sky, the rotation of the Earth can be used to change the projected separations of the antenna pairs, allowing more angular scales to be sampled. An ensemble of the data, i.e., sinusoids of various amplitude and phase, can be then "summed" via the Fourier transform to produce an image of the sky brightness distribution. How well this image reflects the actual sky brightness distribution depends on how completely the relevant angular scales have been sampled. With its emphasis on delivering images at high angular resolution, interferometry works extraordinarily well for observing intrinsically compact targets.

The following sections expand upon these basic ideas. We begin by introducing the basic concepts of radio astronomy, and then move to the principles of aperture synthesis.

3.2 Single-dish Response

The brightness, or equivalently specific intensity, $I_\nu$, is defined as the electromagnetic (EM) power $\delta P$ within a range of frequencies (a bandwidth) $\delta \nu$ received from a solid angle $\delta \Omega$ and intercepted by surface area $\delta A$:

$$I_\nu = \frac{\delta P}{\delta \Omega \delta A \delta \nu}, \quad (3.1)$$

where $I_\nu$ has typical units of W m$^{-2}$ Hz$^{-1}$ sr$^{-1}$. In addition, the flux density, $S_\nu$, is defined as the integration
CHAPTER 3. PRINCIPLES AND CONCEPTS OF INTERFEROMETRY

Figure 3.1: The Plateau de Bure Interferometer (top) and the Submillimeter Array (SMA) (bottom) are the precursors to the ALMA telescope; both are still in full operation and pioneered the science of millimeter wave interferometry.

of brightness over the solid angle of the emitting source:

\[ S_\nu = \int I_\nu \, d\Omega, \]  

(3.2)

where \( S_\nu \) has typical units of W m\(^{-2}\) Hz\(^{-1}\). In millimeter/radio astronomy, the power received is typically so weak that a convenient unit to use for \( S_\nu \) is the Jansky (Jy), where 1 Jy = \( 10^{-26} \) W m\(^{-2}\) Hz\(^{-1}\). A radio telescope with effective area \( A_e \) receives power \( P_{\text{rec}} \) per unit frequency from an unpolarized source:

\[ P_{\text{rec}} = \frac{1}{2} I_\nu A_e \delta \Omega. \]  

(3.3)

The coefficient of 1/2 in Equation 3.3 comes from the fact that a receiver is generally sensitive to only one mode of polarization. (ALMA’s receivers, however, have been constructed with two independent receptors so that both modes of polarization can be detected simultaneously, so the coefficient is actually 1 for ALMA.) The antennas bring incident EM power to a focus after reflecting it off a primary surface. The antenna response, i.e., its relative sensitivity, is a summation of all EM power brought to the focus.

Antenna response is actually dependent on the angle from the on-axis pointing direction of the antenna due to diffraction, i.e., self-interference. To demonstrate the angular dependence, Figure 3.2 shows in the top panel the case for EM power of wavelength \( \lambda \) arriving along the axis of an unobstructed antenna of diameter \( D \). Since the source of the EM power is very distant, the EM power arrives at the primary surface essentially as plane-parallel wavefronts. Note that the antenna surface is parabolic in shape, so the path that each part of the front travels to the focus is constant. With zero path difference, the EM power arriving on-axis is coherently summed at the focus. This arrangement is only true, however, along the axis of the antenna. In the lower panel of Figure 3.2, the case for EM power arriving from an off-axis direction is shown. In this situation, the EM power does not add as constructively. In addition, the diameter of the antenna projected along the off-axis direction is less than the true diameter, decreasing the amount of power received from that direction. As a result, the antenna power response, i.e., its relative sensitivity, will be less than that found on-axis. In particular, at the off-axis angle of \( \lambda/D \) radians, where \( \lambda \) is the wavelength of observation and \( D \) is the diameter.
3.2. SINGLE-DISH RESPONSE

Figure 3.2: Schematic of an incoming plane-parallel wavefront reflecting off an antenna of diameter $D$ and being brought to a focus. The top panel shows the case for a wavefront arriving on-axis. The bottom panel shows the case for a wavefront arriving off-axis. Note that the paths of incident EM power in the first case are all of equal length, and hence the power is summed constructively at the focus. In the second case, the path lengths differ, leading to less constructive summations at the focus.

As an illustration, Figure 3.3 shows an example of a one-dimensional antenna power response with angle for a 12-m diameter parabolic antenna uniformly illuminated by emission of wavelength $\approx 0.85$ mm (350 GHz). The power response is largest on-axis but it declines to zero in $\sim 18$ arcseconds. The central Gaussian-like feature is called the primary beam or the antenna beam size and it has a Half Power Beam Width (HPBW) given by:

$$\text{HPBW Primary Beam} = 1.02 \times \lambda/D.$$ \hspace{1cm} (3.4)

For example, the HPBW of a uniformly illuminated antenna of 12 m diameter at $\lambda = 0.85$ mm is $14.95^\circ$. HPBW is sometimes referred to as Full Width at Half Power or FWHP. In one dimension, the FWHP is equivalent to the Full Width at Half Maximum or FWHM, a quantity sometimes used to describe the width of a Gaussian approximation to the central feature.

Note that the antenna power response rises and declines repeatedly at ever larger angles. The constructive and destructive interference at larger angles leads to successive sidelobes (whose maxima decline with increasing
angle) and *nulls* respectively. The first sidelobes have a relative response of only 1.74% that of the primary beam. Nevertheless, incident emission, if bright enough, coming in at angles well beyond those of the primary beam can make a large contribution to the received EM power. The angular distance between the first nulls is termed the Full Width Between Nulls (FWBN), and is given by:

$$\text{FWBN Primary Beam} = 2.44 \times \lambda/D.$$  \hfill (3.5)

Half the FWBN of the primary beam, \(\sim 1.22 \lambda/D\), is considered the Rayleigh resolution of the antenna, i.e., its ability to distinguish objects on the sky separated by some angular distance. For convenience, the antenna power response is typically normalized to 1.0 along the axis. Figure 3.3 illustrates the antenna power response in one dimension (in log units); on the actual sky, the antenna power response is two-dimensional, and is obtained by rotating the function shown in Figure 3.3 about its central axis.

Up until now, an idealized antenna was described. The power response of an actual antenna, however, can be altered by various effects, including the degree by which the secondary is illuminated, diffraction by the arms supporting the secondary, and surface imperfections. For ALMA, each 12-m antenna has a secondary reflector and support arms that block an effective area of 0.75 m diameter on the primary surface. The actual ALMA feedhorns were designed to provide an antenna power response with a nearly Gaussian primary beam and low sidelobes, preserving as much resolution and sensitivity as possible. The actual ALMA 12 m antennas have measured primary beam HPBW values of \(\sim 1.13 \lambda/D\).
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An equivalent way to consider the antenna power response is in terms of the voltage response, $V(\theta)$, where $P(\theta) \propto V^2(\theta)$. In the far-field, i.e., under the Fraunhofer approximation, a diffraction pattern at the point of observation is the Fourier transform of the field distribution at an aperture. Hence, the voltage response at the focus is the Fourier transform of the aperture shape. For an unobstructed antenna, the aperture is a uniform circle, and $V(\theta) = J_1(\theta)/\theta$, where $J_1(\theta)$ is the Bessel function of the first kind. $P(\theta)$ is correspondingly proportional to $(J_1(\theta)/\theta)^2$. The normalized version of the antenna power response, $P_N$, is also known as the Airy function.

Defining $\theta$ and $\varphi$ as orthogonal directional variables (e.g., sky coordinates), $I_\nu(\theta, \varphi)$ and $P_N(\theta, \varphi)$ can be defined as the directional functions of the sky brightness and the normalized antenna power, respectively. The total received power of an antenna at a given pointing is the integration over the sky of the product of the sky brightness distribution and the antenna power response:

$$ P_{\text{rec}} = \frac{1}{2} A_e \int_{4\pi} I_\nu(\theta, \varphi) P_N(\theta, \varphi) \delta \Omega. \quad (3.6) $$

In addition, the solid angle of the antenna power response $P_N(\theta, \varphi)$ is defined by:

$$ \Omega_A = \int_{4\pi} P_N(\theta, \varphi) \delta \Omega. \quad (3.7) $$

### 3.3 Visibilities and Aperture Synthesis

Observing at millimeter/radio wavelengths is essentially diffraction-limited, as the angular diameter of the Airy disk is the FWHM for a circular aperture of diameter $D$ is $\sim 1.22 \lambda/D$. Millimeter/radio wavelength (single-dish) observations have lower resolutions than optical wavelength (single-dish) observations because $\lambda$ is larger by many orders of magnitude. Though the $D$ of millimeter/radio single-dish telescopes can also be much larger than those of optical wavelengths, the increase in $D$ possible for single-dish telescopes is generally never enough to obtain angular resolutions comparable to ground-based optical telescopes, e.g., 1" or better. For example, the JCMT 15-m diameter antenna has an angular resolution at 850 $\mu$m of $\sim 14''$, and the Arecibo 300-m diameter antenna has an angular resolution at 21 cm of $\sim 3'$. To obtain higher angular resolution images at millimeter/radio wavelengths, signals from physically separated antennas can be combined through interferometry. With this technique, sometimes called aperture synthesis, it is possible achieve a resolution that emulates the effect of having a radio telescope with a larger diameter. Observers, however, must contend with the reality that only certain angular scales, i.e., those determined by the projected separations of each pair of antennas, will be sampled. This section builds on the concepts introduced previously to discuss aperture synthesis in more detail. This powerful technique of aperture synthesis was pioneered by Sir Martin Ryle, who shared the Nobel Prize in Physics in 1974 in part for its development.

As there are no path differences in the case of a plane-parallel wavefront arriving on-axis an antenna, EM power from across the antenna is brought together in phase as the focus. Now imagine that the parabolic surface is divided into $N$ smaller contiguous areas, i.e., elements. In this situation, the received voltage $V(t)$ is the sum of contributions $\Delta V_i(t)$ from each of element:

$$ V(t) = \sum_i \Delta V_i(t) \quad (3.8) $$

The power received by the antenna is proportional to the running time average of the square of the contributions from each element. Assuming illumination is the same for each element, the expression for received power in terms of the sum of time averages of the products of voltages from element pairs can be rewritten:

$$ \langle P \rangle \propto \langle (\sum \Delta V_i)^2 \rangle = \sum \sum \langle (\Delta V_i \Delta V_k) \rangle. \quad (3.9) $$
Next, this expression can further be rewritten in terms of the sums of element pairs which are the same and those which are not, i.e.,

$$\langle P \rangle \propto \sum (\Delta V_i^2) + \sum_{i \neq k} (\Delta V_i \Delta V_k).$$  \hspace{1cm} (3.10)

The first and second sets of terms in Equation 3.10 are called auto-correlation and cross-correlation terms, respectively, since the voltages multiplied in each term are from either the same or different elements, respectively.

From Equation 3.10, any measurement with a large filled-aperture telescope can be understood as being a sum in which each term depends on contributions from only two of the \(N\) elements. As long as the contributions from each element arrive at the focus in phase, there is no need for the elements to be physically contiguous. Generalizing, each cross-correlation term \(\Delta V_i \Delta V_k\) in Equation 3.10 can be measured with two smaller, physically separated antennas (at locations \(i\) and \(k\)) by measuring the average product of their output voltages with a correlating, i.e., multiplying, receiver. Moreover, if the source properties do not change, there is no need to measure all pairs at the same time. A given parabolic surface with \(N\) elements has \(N(N-1)/2\) pairs of elements, and these could be observed sequentially to "synthesize" a measurement by a large filled-aperture telescope. Alternatively, numerous pairs of antennas, with each antenna considered an element, can be distributed to positions at distances much larger than it is possible to build a single filled-aperture telescope, and the signals received by these antennas can be combined in phase to approximate the resolving power of a single filled-aperture telescope.

The above situation only describes the emission received on-axis from antenna pairs. Of course, as noted above, emission also arrives at the antennas from other directions, leading to phase differences. To understand the power response expected from a pair of antennas, let’s look at the ideal 1-D situation of a two-antenna interferometer.

Figure 3.4 shows a schematic picture of a two-antenna interferometer separated by distance \(b\), known as a baseline. This distance can be measured in units of the observing wavelength, \(\lambda\). In terms of familiar units of length, \(b = L/\lambda\), where \(L\) is the distance between antennas and \(\lambda\) is the wavelength in the same unit, e.g., meters. Both antennas observe a common position \(s_o\) located at an angle \(\theta\) from the meridian. The projected separation of the two antennas towards \(s_o\) from the perspective of the source is \(u = b \cos \theta\). In this example, an on-axis wavefront incident to both telescopes reaches antenna 1 first and the wavefront reaches antenna 2 a little later, having traversed an extra path length of \(b \cdot s_o = b \sin \theta\). In other words, emission received by antenna 1 experiences a geometrical delay relative to that received by antenna 2, where the time equals \(\tau_g = b \cdot s_o / c\). To compensate for the geometrical delay, an artificial delay can be inserted into the signal path of antenna 2 (e.g., electronically) so that the signals from both antennas arrive at the correlator with the same phase.

Moving slightly off-axis, a small angle from the axis can be described as \(\alpha\), and its 1-D sky position as \(l = \sin \alpha\), i.e., the direction cosine. At angle \(\alpha\), an off-axis signal reaching antenna 1 will have to travel a slightly longer path than an on-axis signal reaching antenna 2, even with the geometrical delay introduced to compensate for an on-axis signal. This extra path length is \(x = u \sin \alpha = ul\). All distances can be considered in this situation in units of the wavelength of the emission, \(\lambda\), so that \(x\) is the number of wavelengths within a given distance. The extra path lengths result in phase differences with \(\alpha\) that can be characterized where the voltage response of antenna 2, \(V_2\), can be written in terms of the product of the voltage response of antenna 1, \(V_1\), and a phase delay factor sinusoidally varying as a function of angle:

$$V_2 = V_1 e^{2\pi i(ul)}.$$  \hspace{1cm} (3.11)

Expanding to two dimensions introduces \(\beta\), a direction on the sky orthogonal to \(\alpha\). Also, \(m = \sin \beta\) is defined as the small angle analog to \(l\) in this new direction. Note that the baseline here is actually a 2-D vector with components in both dimensions, i.e., \(b_1\) and \(b_2\). Hence, we define \(u = b_1 \cos \theta\) and \(v = b_2 \cos \phi\) where \(\phi\) is the angle of the position \(s_o\) on the sky from the reference position orthogonal to \(\theta\). Finally, \(y\) is defined as the extra path length introduced in this new direction, in units of the wavelength of emission, i.e., \(y = v \sin \beta = vm\). With these changes, the two-dimensional voltage response of antenna 2 is:
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Figure 3.4: An ideal 1-D two-antenna interferometer consisting of two antennas, 1 and 2, separated by physical distance (i.e., a baseline) \( b \). The antennas are both pointed towards a sky location given by \( s_0 \), which is at an angle \( \theta \) from the meridian. The projected distance between the two antennas in that direction is thus \( u = b \cos \theta \). The two antennas are connected to a correlator where the voltages detected from each are combined.

\[
V_2 = V_1 e^{2\pi i (ul + vm)},
\]  

(3.12)

Where \( u \) and \( v \) are identified as specific spatial frequency components of the sinusoid in the E-W and N-S directions respectively, and these are the projected lengths of the antenna separations measured in units of the wavelength at the time of observation. Also, \( l \) and \( m \) are identified as direction cosines relative to a reference position in the E-W and N-S directions, respectively. Typically, the on-axis position \( s_0 \) has \( l = 0 \) and \( m = 0 \) and is called the phase center.

The correlator acts as a multiplying and time-averaging device for the incoming signals from antennas 1 and 2. Hence, its output is:

\[
\langle V_1 V_2 \rangle = \langle \iint V_1(l, m)dl dm \iint V_2(l, m)dl dm \rangle
\]  

(3.13)

Under the assumption that signals emanating from different parts of the sky are incoherent, i.e., they have no similarities in phase, the time averages of the correlation of those signals will be zero. Thus, the product of the integrals in Equation 3.13 can be simplified to:

\[
\langle V_1 V_2 \rangle = \langle \iint V_1(l, m)V_2(l, m)dl dm \rangle
\]  

(3.14)
\[ \langle V_1 V_2 \rangle = \int \int (V_1(l, m) V_2(l, m)) dldm \quad (3.15) \]

\[ \langle V_1 V_2 \rangle = \int \int (V_1(l, m)^2 e^{2\pi i (ul + vm)}) dldm. \quad (3.16) \]

As \( V^2 \propto P \) (see Equation 3.9), and \( P \propto I_0 \) (see Equation 3.3),

\[ \langle V_1 V_2 \rangle \propto \int \int I(l, m) e^{2\pi i (ul + vm)} dldm \quad (3.17) \]

where \( I(l, m) \) is the intensity distribution on the sky. The correlator therefore measures a quantity known as the complex visibility, \( \mathcal{V} \), which is formally the Fourier transform of the intensity distribution on the sky:

\[ \mathcal{V}(u, v) = \int \int I(l, m) e^{2\pi i (ul + vm)} dldm = A e^{i\phi}. \quad (3.18) \]

Note that \( \mathcal{V} \) is a complex number, and can be described by an amplitude, \( A \), and a phase, \( \phi \). The amplitude and phase contain information about the source brightness and its location relative to the phase center, respectively, at spatial frequencies \( u \) and \( v \).

### 3.4 The Visibility or \((u, v)\) Plane

The relationship between the sky brightness distribution and the complex visibility distribution is governed by the van Cittert-Zernike theorem and it is the basis of aperture synthesis. Given that the complex visibility is the Fourier transform of the sky brightness distribution in the image plane, it follows that the sky brightness distribution is in turn the inverse Fourier transform of the complex visibility distribution in the visibility plane:

\[ \mathcal{V}(u, v) = \int \int I(l, m) e^{2\pi i (ul + vm)} dldm \quad (3.19) \]

\[ I(l, m) = \int \int \mathcal{V}(u, v) e^{-2\pi i (ul + vm)} dudv \quad (3.20) \]

By measuring the distribution of complex visibilities (in the visibility or \((u, v)\) plane), in principle the sky brightness distribution can be recovered. In essence, an image is a "sum" (i.e., the Fourier transform) of the visibilities where each visibility has an amplitude and phase representing the brightness and relative position of emission on a specific angular scale. The image and its Fourier transform are conjugates of each other, and each contains the same amount of information.

Two antennas at a given physical distance \( b \) can have signals interfered to sample the sky brightness distribution on a scale inversely proportional to the projection of that distance on the sky. As shown above, the response of the interferometer is sinusoidal, and is sometimes referred to as a fringe, with spacing on the sky in the 1-D case of:

\[ \text{Fringe Spacing} = 1/u \text{ (radians)} = 1/(b \cos \theta) = \lambda/(L \cos \theta). \quad (3.21) \]

In effect, the interference of the signals modifies the angular response of the antennas and the antennas can "see" the true sky brightness distribution only on the scale defined by the wavelength of the sinusoid. As the fringe spacing depends inversely on the projected distance, antennas closer together measure emission on larger scales. Conversely, those antennas spaced further apart measure emission on smaller scales. Since fringe spacing also depends on the wavelength of emission, as \( b \) is measured in numbers of wavelengths, observing shorter or
longer wavelengths also can sample smaller or larger scales, respectively. These ideas can be easily generalized to two dimensions, with the fringe spacing and on-sky orientation depending on the relative magnitudes of $u$ and $v$.

A given pair of antennas will only instantaneously sample a single scale of the sky brightness distribution. Given the E-W and N-S separations of the pair, a visibility in the $(u, v)$ plane is measured. Since visibilities are samples of a complex-valued function with Hermitian symmetry, a single sampling gives two visibilities, one at $(u, v)$ and its complex conjugate at $(-u, -v)$. To recover the true sky brightness distribution, however, knowledge of the distribution of visibilities across the $(u, v)$ plane would be needed. Improving coverage of visibilities over the $(u, v)$ plane can be done in several ways. First, multiple antennas can be incorporated into an array, with each at a different distance from the others to prevent redundancy. An array of $N$ antennas will have $N(N-1)/2$ independent baselines, with each pair providing a single pair of samples in the $(u, v)$ plane. Second, a target can be observed repeatedly by the array as it appears to move across the sky due to the Earth’s rotation. Though the physical distances between the antennas do not change, their projected distances do change depending on the altitude and azimuth of the target. Hence, repeated observations by all the pairs in an array can sample many visibilities across the $(u, v)$ plane. Finally, antennas in the array may be arrangeable in several configurations so that pairs of antennas have different distances and can sample different parts of the $(u, v)$ plane. Assuming the source emission is not variable, the combination of these schemes can reasonably sample the $(u, v)$ plane, yielding an image that can resemble the true sky brightness distribution.

### 3.5 Fields-of-view and Mosaics

Each antenna of an actual interferometer has finite diameter. As noted before, such antennas have their own power responses on the sky $P_N = A(l, m)$ (e.g., the Airy function for an unobstructed, uniformly illuminated aperture). Indeed, the individual antenna response fundamentally limits the extent of an interferometric image made with a single pointing. In practice, the HPBW of the primary beam serves as the “field-of-view” of the single-pointing interferometric image. Moreover, $A(l, m)$, the Airy function, is actually included formally in the correlator output:

$$\mathcal{V}(u, v) = \int \int A(l, m)I(l, m)e^{-2\pi i(u\ell + v\ell)} d\ell dm. \quad (3.22)$$

Hence, an interferometer actually measures the Fourier transform of the sky brightness distribution multiplied by the antenna power response. To recover $I(l, m)$, the image resulting from the Fourier transform of the complex visibilities must be divided by $A(l, m)$ as the last step of image processing. This so-called primary beam correction should be performed if the image contains scientifically relevant emission off the phase centre. Otherwise, intensities and flux densities measured at those locations will not be accurate.

To counteract the angular fall-off of sensitivity due to the primary beam response, or even to sample emission over areas on the sky larger than the primary beam, an interferometer can observe adjacent positions, producing a mosaic. Sensitivity across a mosaic depends on the spacing of the individual positions observed. A mosaic can have close to uniform sensitivity with a minimum number of pointings if the positions observed are arranged in a grid of equilateral triangles spaced by $\lambda/(\sqrt{3}D)$, where $D$ is the diameter of the antenna. With this spacing, the fall-off of the primary beam response at one pointing is made up by the responses of the primary beam at adjacent pointings, except of course at the edge of the mosaic. Mosaics can be made with adjacent positions that are spaced either closer or more distant, with non-uniform sensitivities. Mosaics provide increased areal coverage but at a cost of more observing time. For example, uniform sensitivity can be obtained across the primary beam of a single pointing but requires a minimum of six other pointings around the single pointing. These six pointings are arranged in a hexagonal pattern with each vertex spaced $\lambda/(\sqrt{3}D)$ from each other (and the single pointing). A single image is produced by combining the visibilities obtained at all pointings into a single ensemble that is simultaneously Fourier transformed.
3.6 Spatial Filtering

Through the principles described in Sections 3.3-3.5, the true sky brightness distribution can be recovered. It is, however, impossible in practice to sample completely the \((u, v)\) plane and obtain all visibilities. The incomplete \((u, v)\) plane sampling effectively provides a fundamental limit to the level of detail discernible in the sky brightness distribution, i.e., down to a minimum scale defined as the resolution. In addition, incomplete sampling results in spatial filtering of the true sky brightness distribution, i.e., the resulting images do not contain information on angular scales unobserved by the interferometer. In particular, the lack of coverage at the shortest baselines (i.e., lower than those sampled by the smallest baselines) results in an intrinsic lack of sensitivity to large-scale emission. It is crucial for ALMA users to understand these limitations. Below, we describe the ideas behind spatial filtering and provide a few illustrative examples. Further examples can be found in Chapter 7 (“Imaging with ALMA”), Sections 7.6 and 7.7.

First let’s discuss resolution. The resolution of any interferometric image depends on the distribution of visibilities sampled. Assuming a finite number of \(M\) visibilities has been obtained, the \((u, v)\) plane has been sampled at \(2M\) discrete points. The sampling distribution can then be characterized as an ensemble of \(2M\) (Dirac) delta functions:

\[
B(u, v) = \sum_{k=1}^{2M} \delta(u - u_k, v - v_k). \tag{3.23}
\]

The inverse Fourier transform of this ensemble of visibilities can then be written as:

\[
I^D(l, m) = FT^{-1}\{B(u, v)V(u, v)\}. \tag{3.24}
\]

Following the convolution theorem, the Fourier Transform of a convolution of two functions is the product of the Fourier Transforms of those functions. Hence, Equation 3.24 can be rewritten as:

\[
I^D(l, m) = b(l, m) * I(l, m)A(l, m). \tag{3.25}
\]

In effect, the image obtained is the convolution of the true sky brightness distribution (modified by the antenna power response \(A(l, m)\)) with the point spread function, \(b(l, m) = FT^{-1}\{B(u, v)\}\), the Fourier transform of the \((u, v)\) plane sampling distribution. The point spread function is sometimes called the synthesized beam or the dirty beam. It is important to distinguish this beam from the single-dish response function \(A(l, m)\), which in the interferometry context is called the primary beam. The image resulting from the Fourier transform of a finite number of visibilities, \(I^D(l, m)\), is sometimes referred to as the dirty image.

The measure of how similar an image is to the true sky distribution is sometimes referred to as image fidelity. Image fidelity depends on the specifics of coverage of the \((u, v)\) plane sampled by the interferometer. Since the numbers of samples are necessarily finite and discrete, there are invariably gaps in any practical sampling of the \((u, v)\) plane. These gaps mean that no information is obtained about the true sky brightness distribution on those specific angular scales. Note that visibilities corresponding to those unobserved scales can have any value. With no information, however, it is typically assumed that \(V(u, v) = 0\) at unsampled locations in the \((u, v)\) plane. Including these visibility domain gaps through the Fourier transform produces aliased features in the resulting image, the magnitude of which depends on the extents and locations of gaps in the \((u, v)\) plane and the brightness of emission on sampled scales. If the \((u, v)\) plane has been reasonably well sampled, the synthesized beam will consist of a compact positive feature surrounded by positive and negative features of lower relative amplitude. These latter features, also called sidelobes, can complicate the image since brightness is distributed via the point spread function throughout the image. The resulting image can have significant artefacts depending on the sky brightness distribution and the sampling of the \((u, v)\) plane. A dirty image, however, can be improved through deconvolution techniques to minimize the effect of incomplete spatial frequency sampling (e.g., CLEAN and its variants; see Chapter 7).

Though speaking generally about true sky brightness distributions so far, a special note should be made for the case of point sources. Obviously, a point source is a distribution of emission that is not extended relative
3.6. SPATIAL FILTERING

Figure 3.5: Imaging concepts. Panel a (upper left): Example of a dirty beam, \( b(l,m) \). Panel b (upper right): The related ensemble of discrete points sampled in the \((u,v)\) plane, \( B(u,v) \). The black points were obtained from a compact configuration while the red ones were obtained from an extended configuration. Panel c (lower left): Example of a true sky distribution, \( I(l,m) \). Panel d (lower right): The dirty image \( I^D(l,m) \) resulting from observing \( I(l,m) \) over the baselines of \( B(u,v) \), or equivalently the convolution of \( I(l,m) \) by \( b(l,m) \). The antenna power response, \( A(l,m) \), has been ignored in this illustration since it is much wider than the true sky brightness distribution. (Figure courtesy of D. Wilner.)

to the resolution of the observation. In this case, the morphology of the dirty image will equal that of the dirty beam (e.g., see Equation 3.25). Moreover, the complex visibilities of the point source have the same amplitudes on all observed angular scales. Of course, sources may appear point-like at low resolutions but may appear extended in higher-resolution observations.

Figure 3.5 illustrates the concepts of dirty beam and dirty image and their impact on the recovered image. Panels a and b (upper pair) show respectively the dirty beam and the related ensemble of locations observed in the \((u,v)\) plane, i.e., the \((u,v)\) coverage. Note in panel a the positive feature in the center of the dirty beam distribution and the surrounding positive and negative features of lower amplitude. These latter features arise from the incomplete sampling of the \((u,v)\) plane seen in panel b. As an aside, note that two sets of \((u,v)\) plane samples are identified in panel b; these result from observations by the same antennas in two different configurations, a compact one (black) and an extended one (red). Panels c and d (lower pair) show respectively an example of a true sky distribution (here, a model of a ring of emission) and the dirty image. The dirty image is the convolution of the true sky distribution by the dirty beam, and one can easily see how incomplete sampling of the \((u,v)\) plane leads to the appearance of significant artefacts in the resulting dirty image.

The resolution of the dirty image is defined effectively by the compactness of the central feature of the dirty beam, e.g., half its FWBN. Since the structure of the dirty beam is generally more complicated than that of a single-dish antenna, e.g., the beam from a uniformly illuminated antenna shown in Figure 3.3, it is not so easy to measure FWBN. Instead, the resolution is typically approximated to first order by the FWHM of a
Gaussian fit to the central feature of the dirty beam. The resolution of the dirty image depends ultimately on how the interferometer antennas are arranged. In general, distributions connected via a Fourier transform scale inversely to each other. For example, narrow distributions in one domain have wide ones in the other, and vice versa. By analogy, an ensemble of discrete points, \( B(u, v) \), clustered around the \( (u, v) \) plane origin provided by a compact configuration yields a low-resolution image since the central beam feature \( b(l, m) \) is wide. Conversely, an ensemble of discrete points distributed more widely from the \( (u, v) \) plane origin yields a high-resolution image since the central beam feature is narrow. Indeed, resolution is fundamentally limited by the extent of the longest baselines in a given configuration. The minimum scale discernible in the image is limited by these maximum baselines. A handy formula for the approximate resolution provided by an interferometer is:

\[
\text{Interferometer Resolution} = \theta_{\text{res}} = k \lambda / L_{\text{max}},
\]

where \( k \) is a factor that depends on how the visibilities are weighted during inversion (typically \( k \approx 1 \); see Figure 7.6) and \( L_{\text{max}} \) is the longest baseline in the array.

Another important limitation of interferometric array observations is insensitivity to large angular scales. This insensitivity arises because interferometric arrays alone cannot sample spatial frequencies lower than those sampled by a baseline equal to an antenna diameter. In effect, visibilities at locations on the \( (u, v) \) plane at or near its origin are not sampled, leading to the so-called zero-spacing problem. The lack of sensitivity to larger scale emission due to the zero-spacing problem biases the resulting image to the compact, small-scale emission of the true sky brightness distribution. As a guideline, the interferometer image has a maximum recoverable scale given roughly by:

\[
\text{Maximum Recoverable Scale} = \theta_{\text{MRS}} \approx 0.6 \lambda / L_{\text{min}},
\]

where \( L_{\text{min}} \) is the minimum baseline in the array configuration. Strictly speaking, for an input Gaussian visibility distribution of FWHM \( \theta_{\text{MRS}} \), the ratio of the brightness at source center of an image made by an array with \( L_{\text{min}} \) to the same made with no central hole in its visibility sampling (i.e., no zero spacing problem) is \( 1/e^2 \). The smallest baseline possible in an array occurs when two antennas are adjacent to each other. Of course, the antennas cannot be moved physically closer together than their diameters. Note that in projection antennas can appear closer than their diameters. In those cases, found typically when observing low elevation sources in compact configurations, the antenna in front partially blocks the reception by the antenna in the rear. The resulting visibilities are distorted, e.g., the antenna in the rear receives less power than the one in front. This situation is called shadowing and affected data are typically removed from the ensemble of observed visibilities.

Figure 3.6 illustrates the idea of spatial filtering using simulations of actual ALMA configurations. The simulations were performed using the Common Astronomy Software Applications (CASA\(^2\)) package, the suite of data reduction and analysis software used by ALMA. Here, panel \( a \) shows an optical image of the galaxy M51 used as an example of a true sky brightness distribution, after changing the frequency of the image to 100 GHz and placing the image center at \( \delta = -40^\circ \), a declination easily observable with ALMA. Panels \( b, c, \) and \( d \) show the recovered images obtained by simulating 43-antenna observations of the galaxy using the CASA task \textit{simobserve} in extended (C43-6), somewhat extended (C43-4), and compact (C43-1) configurations, respectively, and CLEANing. Angular resolutions of \( \sim 0.47^\prime\prime, 1.3^\prime\prime, \) and \( 4.3^\prime\prime \) are obtained, respectively, and the corresponding maximum recoverable scales are \( \sim 4^\prime, \sim 11^\prime, \) and \( \sim 28^\prime \), respectively. Larger-scale emission from the galaxy has been filtered out in the extended configuration observations (panel \( b \)), leaving only the compact structures of its arms. On the other hand, these compact structures are not very discernible in the compact configuration observations (panel \( d \)). A reasonable compromise is found in the somewhat extended configuration observations (panel \( c \)), yet some small-scale detail and larger-scale emission remain missing. Note that though these images are missing angular scales, good science can still be obtained with them, as long as their limitations are properly understood. Combining data obtained from multiple configurations, or having more antennas, would increase the fidelity of recovered images.

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\(^2\)For more information on CASA, see http://casa.nrao.edu
3.7 Multi-configuration Observations

As previously discussed, a given configuration with baselines ranging from $L_{\text{min}}$ to $L_{\text{max}}$ is sensitive to angular scales from $\sim \theta_{\text{MRS}}$ to $\theta_{\text{res}}$. Sensitivity to a broader range of angular scales is possible by combining data obtained in multiple configurations, where more extensive coverage of the $(u, v)$ plane is attained. For example, the same source can be observed in different configurations of the 12-m Array, with more extended configurations providing higher angular resolutions. For sensitivity to extended structures, the more compact 12-m Array configurations or the more tightly clustered 7-m Array can be used. Finally, the individual Total Power (TP) Array antennas of the ACA can be used to map the largest angular scales and address the zero-spacing problem. In this section, we briefly describe so-called multi-configuration observations. Discussion of how ALMA implements
such observations in practice can be found in Sections 7.8 and 7.9.

During proposal preparation, ALMA users should take note of the maximum recoverable scale needed to ensure that the proposed observations will be able to recover the scales needed to address the science in question. The ALMA Observing Tool will determine which combination of configurations will yield the desired angular resolution and maximum recoverable scale.

Data combination appears to work best when the signal-to-noise ratios (SNR) of the datasets are similar. Otherwise, information on scales covered by lower SNR data is relatively less reliable, making interpretation of the images difficult. In addition, the accuracies of the astrometry and calibration of the different datasets are crucial. Assuming the SNR of the individual datasets is high, combination is best done in the visibility domain rather than the image domain to minimize the effect of artifacts produced by aliasing, i.e., incomplete \((u, v)\) coverage, in either dataset. For example, interferometer data obtained from different configurations should be combined in the visibility domain and then the new ensemble should be Fourier transformed to produce a new dirty image.

Combining single-dish data and interferometer data also works best in the visibility domain, as long as both datasets have high SNRs. In this case, the single-dish image can be Fourier transformed into the visibility domain and the resulting visibilities added to the ensemble of those obtained by the interferometer. The new ensemble can be then Fourier transformed en masse to produce a new image. Such data combination works best if the single-dish and interferometric datasets have significant \((u, v)\) coverages in common. For example, a reasonable overlap in \((u, v)\) coverage can provide enough data to reveal amplitude calibration differences that can be minimized by re-scaling the single-dish visibilities relative to the interferometric ones. In general, a reasonable overlap of \((u, v)\) coverage will occur if the single-dish data are obtained by an antenna that has a diameter twice the minimum baseline of the interferometer, e.g., approximately the interferometer antenna diameter. Multiple interferometer pointings, i.e., mosaics, can also partially recover missing low spatial frequency information. In Chapter 7, a specific technique for combining ALMA total power data and 7-m / 12-m Array data called “feathering” is described in some detail.

Single-dish and interferometer data can be done in the image domain as well, and such addition is likely the best route for combining relatively low SNR datasets.

### 3.8 Units and Conversions

Finally, this Chapter ends with a discussion of various units used in millimeter and radio astronomy and describe some useful conversions. Returning to the concept of specific intensity, this quantity can be described alternatively in terms of a temperature:

\[
I_\nu(\theta, \varphi) = \frac{2k\nu^2}{c^2} T_B(\theta, \varphi).
\]

In this equation, \(T_B\) is the brightness temperature, the temperature of a blackbody with the same specific intensity at a given frequency in the Rayleigh-Jeans limit, i.e., \(h\nu/kT \ll 1\). Brightness temperature serves as an equivalent way of expressing the specific intensity of an astronomical source. The unit of brightness temperature is Kelvin (K).

In turn, brightness temperature can be included into the definition of flux density, \(S_\nu\) (Equation 3.2), where

\[
S_\nu = \frac{2k\nu^2}{c^2} \int T_B d\Omega.
\]

Assuming the beam is Gaussian, one can then connect brightness temperature to flux density following:
\[
\left( \frac{T}{\text{K}} \right) = \left( \frac{S_\nu}{\text{Jy}} \right) \left[ 13.6 \left( \frac{300 \text{ GHz}}{\nu} \right)^2 \left( \frac{1''}{\theta_{\max}} \right) \left( \frac{1''}{\theta_{\min}} \right) \right],
\] (3.30)

where \(\nu\) is the observing frequency, and \(\theta_{\max}\) and \(\theta_{\min}\) are the major and minor axis of the synthesized beam, respectively. Note again that flux densities observed by ALMA are typically in units of Janskys, where 1 Jy = \(10^{-26}\) W m\(^{-2}\) Hz\(^{-1}\) = \(10^{-23}\) erg s\(^{-1}\) cm\(^{-2}\) Hz\(^{-1}\). The ALMA Observing Tool converts between temperatures and flux densities using these formulae.

An important point ALMA users must consider when proposing their projects is the dependence of brightness temperature sensitivity on synthesized beam size. Using Equation 3.30, one sees that an rms value in flux density \((S)\) can translate to an rms value in brightness temperature \((\Delta T)\), assuming a given synthesized beam size. Larger beam sizes correspond to lower \(\Delta T\), i.e., the surface brightness sensitivity increases. In turn, extended low surface brightness objects may be harder to detect at higher angular resolutions as the corresponding sensitivities may be too low. Typically, a compromise must be obtained between angular resolution and brightness sensitivity when planning interferometric observations.

### 3.9 Further Reading

If more information about interferometry is desired, the topic is covered in more detail in the following seminal texts:

Chapter 4
Receivers

The ALMA front end can accommodate up to 10 receiver bands covering most of the wavelength range from 8.5 to 0.3 mm (35–950 GHz). Each band is designed to cover a tuning range which is approximately tailored to the atmospheric transmission windows. These windows and the tuning ranges are outlined in Figure 4.1. This illustrates the broad, deep absorption features, mostly due to H$_2$O in the lower few km of the atmosphere, as well as some O$_2$ transitions. The many narrow features seen in this plot are mostly from stratospheric O$_3$, along with some transitions of CO and other trace species. In Cycle 6, Bands 3, 4, 5, 6, 7, 8, 9, and 10 are available, and the basic characteristics of the bands are outlined in Table 4.1. Each of the ALMA receiver bands is described in more detail in the following sections as well as in the references listed in Table 4.2.

![Transmission in All ALMA Bands at Zenith](image)

Figure 4.1: The ten ALMA receiver bands along with atmospheric transmission. The receiver coverage is shown shaded, superimposed on a zenith atmospheric transparency plot at the Array Operation Site (AOS) for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of Precipitable Water Vapor (PWV).

The ALMA receivers in each antenna are situated in a single front-end assembly (see Appendix A, Section A.4). The front-end assembly consists of a large cryostat containing the receiver cold cartridge assembly for
Table 4.1: ALMA Receiver Specifications. Notes to Table: 1. Frequency ranges are the maximum available, at the extreme upper and lower limits of the IF passband. In reality, because of filter roll-off, the coverage is \( \approx 60 \text{MHz} \) less (See Section 6.4). 2. Sideband modes: SSB means single sideband receiver, 2SB means dual sideband receiver where the two sidebands are available simultaneously, DSB means double sideband receiver. See text for details. 3. Usable IF range is extended to allow simultaneous observations of multiple lines. However, the autocorrelation noise performance is degraded by a factor of up to about 1.2 below 5.5 GHz (Section 4.2.4). 4. Maximum instantaneous IF bandwidth, limited by the back-end filters and spectrometers. As both upper and lower sidebands both pass through the same IF bandwidth but are subsequently separated, the effective signal bandwidth given in this column for 2SB receivers is twice the actual IF filter bandwidth. In addition, this is per polarization, so the total effective bandwidth for each receiver is then another factor of 2 higher. Note that the effects of the anti-aliasing filters have been included (see Section 6.4). 5. In Cycle 6, the maximum bandwidth is approximately double in cross-correlation mode, because both sidebands can be separated and correlated using 90-degree phase switching (see Section 6.3.4 and Section B.4.4). 6. The maximum specified SSB receiver temperatures \( T_{\text{rx}} \) are given, unless otherwise noted. These values are the average over the whole IF band. In many cases, the average realised results are better than specifications; the sections on individual receiver bands describe the typical values measured. The numbers adopted in the observing tool (OT) and ALMA sensitivity calculator (ASC) are conservative values of the average value over the band. 7. The specification for Band 3 receivers are described slightly differently to other bands: at LO1=104GHz, \( T_{\text{rx}} < 39 \text{K (including optics)} \), and \( T_{\text{rx}} < 43 \text{K} \) for any other valid LO setting. 8. The specification for Band 10 receivers is \( T_{\text{rx}} < 230 \text{K} \) within a selected 80% portion of that band (787–950 GHz).

<table>
<thead>
<tr>
<th>Band</th>
<th>Frequency / Wavelength range (GHz)(^1)/mm</th>
<th>LO range (GHz)</th>
<th>Sideband mode(^2)</th>
<th>IF range (GHz)(^3)</th>
<th>Inst. IF bandw. (GHz)(^4)</th>
<th>( T_{\text{rx}} ) over 80% of band (K)(^5)</th>
<th>( T_{\text{rx}} ) at any freq. (K)(^6)</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>84.0 - 116.0/2.59 - 3.57</td>
<td>92 - 108</td>
<td>2SB</td>
<td>4-8</td>
<td>7.5</td>
<td>(&lt; 39^7)</td>
<td>(&lt; 43^7)</td>
</tr>
<tr>
<td>4</td>
<td>125.0 - 163.0/1.84 - 2.40</td>
<td>133 - 155</td>
<td>2SB</td>
<td>4-8</td>
<td>7.5</td>
<td>(&lt; 51)</td>
<td>(&lt; 82)</td>
</tr>
<tr>
<td>5</td>
<td>158.0 - 211.0/1.42 - 1.90</td>
<td>166 - 203</td>
<td>2SB</td>
<td>4-8</td>
<td>7.5</td>
<td>(&lt; 55)</td>
<td>(&lt; 75)</td>
</tr>
<tr>
<td>6</td>
<td>211.0 - 275.0/1.09 - 1.42</td>
<td>221 - 265</td>
<td>2SB</td>
<td>4.5-10(^4)</td>
<td>7.5</td>
<td>(&lt; 83)</td>
<td>(&lt; 136)</td>
</tr>
<tr>
<td>7</td>
<td>275.0 - 373.0/0.80 - 1.09</td>
<td>283 - 365</td>
<td>2SB</td>
<td>4-8</td>
<td>7.5</td>
<td>(&lt; 147)</td>
<td>(&lt; 219)</td>
</tr>
<tr>
<td>8</td>
<td>385.0 - 500.0/0.60 - 0.78</td>
<td>393 - 492</td>
<td>2SB</td>
<td>4-8</td>
<td>7.5</td>
<td>(&lt; 196)</td>
<td>(&lt; 292)</td>
</tr>
<tr>
<td>9</td>
<td>602.0 - 720.0/0.42 - 0.50</td>
<td>610 - 712</td>
<td>DSB</td>
<td>4-12</td>
<td>7.5(15)(^5)</td>
<td>(&lt; 175) (DSB)</td>
<td>(&lt; 261) (DSB)</td>
</tr>
<tr>
<td>10</td>
<td>757.0 - 950.0/0.32 - 0.38</td>
<td>795 - 942</td>
<td>DSB</td>
<td>4-12</td>
<td>7.5(15)(^5)</td>
<td>(&lt; 230) (DSB)</td>
<td>(&lt; 344) (DSB)</td>
</tr>
</tbody>
</table>

Table 4.2: Technical papers describing the receiver bands, optics and the water vapor radiometer.
4.1 LOCAL OSCILLATORS AND IF RANGES

Each band (the CCA, including Superconductor-Insulator-Superconductor (SIS) mixers and Local Oscillator (LO) injection) and the Intermediate Frequency (IF) and LO room-temperature electronics of each band (the warm cartridge assembly, WCA). The cryostat is kept at a temperature of 4 K through a closed cycle cooling system. The Amplitude Calibration Device (ACD) used for ambient/hot load calibration is mounted above the front end (see Appendix A.5). Each receiver cartridge contains two complete receiving systems sensitive to orthogonal linear polarizations. The designs of the mixers, optics, LO injection scheme, and polarization splitting vary from band to band, depending on the optimum technology available at the different frequencies; each receiver is described in more detail in the sections below.

Up to three bands can be switched on at a time (more would risk overloading the cryostat cooler). From a hardware point of view, it takes only about 1-2 seconds to switch between two bands which are powered up (limited by the LO locking and antenna movement). For bands that are not switched on, the time to fully thermally stabilize a receiver from an off state is up to ~15 minutes - this is mainly to ensure the optimum flat bandpass shape. All of the receivers are mounted off-axis in order to avoid extra rotating band selection mirrors, which necessitates an offset of the antenna pointing to change band. It means that only one receiver can be used at a given time.

4.1 Local Oscillators and IF Ranges

The observed sky frequencies need to be mixed down to frequencies in the range 2 to 4 GHz in order to send the digitized signals to the correlator (see Section 6.1). This frequency downconversion involves a set of LOs, and the LO and IF systems are described in detail in Appendix B.

The front-end mixer uses LO1 to down-convert the sky frequencies into an IF band with a range of 4–12 GHz. This covers the needs of all the ALMA bands, since the mixers for Bands 3, 4, 5, 7, and 8 have an output range of 4–8 GHz, Band 6 a range of 4.5–10 GHz and Bands 9 and 10 a range of 4–12 GHz (Table 4.1). The possible sky frequency ranges covered by each receiver with the first LO, LO1, set to a frequency $F_{LO1}$ are:

- For the lower sideband (LSB): $(F_{LO1} - IF_{lo})$ to $(F_{LO1} - IF_{hi})$
- For the upper sideband (USB): $(F_{LO1} + IF_{lo})$ to $(F_{LO1} + IF_{hi})$

where $IF_{lo}$ and $IF_{hi}$ are the lower and upper IF ranges in the “IF Range” column of Table 4.1, and the IF bandwidth (per sideband) is $IF_{hi} - IF_{lo}$. The ranges of LO1 are given in column 3 of Table 4.1. This is illustrated in Figure 4.2. Note that the maximum IF bandwidth in Table 4.1 may be a few percent less than the IF range in Table 4.1 (see Section 6.4).

![Figure 4.2: IF ranges for the two sidebands in a heterodyne receiver.](image-url)
4.2 The Cycle 6 Receivers

The receivers for Bands 3 to 8 are two-sideband (2SB) receivers, where both the upper and lower sidebands are provided separately and simultaneously. There are 4 outputs from each of the receivers, comprising the upper and lower sidebands in each of the two polarizations. Each output has a bandwidth of 4 GHz (slightly more for Band 6) which is reduced to an effective total bandwidth of 3.75 GHz due to the anti-aliasing filters, etc., see Appendix B.3.5. The mixers give 10 dB or more unwanted sideband rejection, which is adequate for reducing the degradation of S/N from noise in the unwanted sideband, but not adequate for suppressing astronomical signals in the unwanted sideband. Further suppression is performed by offsetting LO1 and LO2 (and eventually the tunable filter LO, TFB LO) by small and opposite amounts, as well as 180-degree phase-switching, the offsets of which depend on the antenna such that the signals from two antennas in the image sideband do not correlate (see Chapter 6).

Bands 9 and 10 use double sideband (DSB) receivers, where the IF contains noise and signals from both sidebands. They only have two outputs, one per polarization. The IF effective bandwidth is 7.5 GHz per polarization (after passing through the IF processing units), so the total instantaneous receiver bandwidth is the same as Bands 3-8. However, both sidebands can be simultaneously correlated using 90-degree phase switching (see Chapter 6), although this does not remove the noise from the opposite sideband. The effective system bandwidth available in Bands 9 and 10 is therefore double that of the 2SB receivers.

Each of the ALMA receiver bands is different in several aspects, and the following sections describe the individual bands in more detail.

Figure 4.3 shows the layout of the ALMA receiver bands in the ALMA cryostat.

**Definition of ALMA Receiver Polarization**

![Diagram of ALMA receiver bands and polarization](image)

Figure 4.3: The layout of the ALMA receiver bands within the ALMA cryostat. Also shown is the orientation of the receiver polarization position angles.
4.2. THE CYCLE 6 RECEIVERS

4.2.1 Band 3 Receiver

Band 3 is the lowest frequency band available in Cycle 6, covering a frequency range of 84 to 116 GHz (in the 3 mm atmospheric window). The cartridge is fed by a “periscope” pair of ellipsoidal pickoﬀ mirrors located outside the cryostat, which refores the beam through the cryostat window, allowing for a smaller window diameter (Figure 4.4). A single feedhorn feeds an ortho-mode-transducer (OMT) which splits the incoming signal into two linear orthogonal polarizations and feeds the SIS mixers.

![Figure 4.4: Input optics for Band 3, showing the warm pickoﬀ mirrors. The location of the antenna beam from the secondary mirror is shown by the solid line, and the Cassegrain focus is shown by the small circle to the upper right.]

A block diagram of the Band 3 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.5. The Cold Cartridge Assembly (CCA) contains the corrugated feedhorn, OMT, SIS mixers and the low-noise HEMT first IF amplifiers. At room temperature, the Warm Cartridge Assembly (WCA) includes further IF amplification and the Local Oscillator covering 92–108 GHz.

The acceptance criteria for the Band 3 receiver noise performance ($T_{rx}$) is $< 39$ K at LO1=104 GHz, and $< 43$ K for any other valid LO setting. This includes a 2 K component due to the losses in the external warm optics. The achieved median performance is typically somewhat better than this, with $T_{rx} \approx 37 \pm 6$ K. The OT is a little conservative, and assumes $T_{rx}=45$ K. The atmospheric transmission over most of Band 3 is very high, even with a large PWV (Figure 4.6) which means observations in Band 3 can (and do) take place with up to 10 mm or more of PWV. The system temperature ($T_{sys}$) shows the expected rise at the higher end due to the edge of an atmospheric oxygen line (Figure 4.7), so this increase is mostly independent of the PWV.
Figure 4.5: Block diagram of the Band 3 receiver (left) including CCA (upper) and WCA (lower). Right image shows a Band 3 CCA. Note the single feedhorn which feeds the OMT, splitting the two polarization signals for the 2SB mixers. The Band 3 cartridges were constructed in Canada at NRC-HIA, Victoria.
4.2. THE CYCLE 6 RECEIVERS

Figure 4.6: Band 3 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. The decrease at the high frequency end is due to the edge of an \( \mathrm{O}_2 \) absorption line at 118.75 GHz.

Figure 4.7: Typical system temperature \( (T_{\text{sys}}) \) at zenith for Band 3 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. \( T_{\text{sys}} \) was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.2 Band 4 Receiver

The Band 4 receiver covers the 125 to 163 GHz spectral window (in the 2 mm atmospheric window). The signal collected by the telescope is focused to the Band 4 cartridge using a set of warm mirrors (Figure 4.8). A single feedhorn feeds an OMT which splits the two linear polarizations and feeds the 2SB SIS mixers.

![Figure 4.8: Optical layout of the Band 4. Red indicates Band 4 Gaussian beam of 5 times the beam width.](Image)

A block diagram of the Band 4 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.9. The Band 4 CCA contains a feed horn, an OMT as a polarization splitter, 2SB SIS mixer assemblies, cold IF amplifiers, isolators, and LO frequency doublers. The RF signal is down converted to 4–8 GHz using a 2SB mixer unit.

The atmospheric transmission in Band 4 is shown in Figure 4.10 for different PWV values. Most observations in Band 4 will be done with PWV < 5 mm. The specification for Band 4 receiver noise performance ($T_{rx}$) is $<51$ K over 80% of the band, and $<82$ K over the whole band (SSB). The OT assumes 51 K. The resulting system temperatures ($T_{sys}$) for a range of PWV values are shown in Figure 4.11. The on-array performance of the receivers is often more typically $\sim40$ K over most of the band, so the achieved $T_{sys}$ may be slightly lower.
Figure 4.9: Block diagram of the Band 4 receiver (left) including CCA (upper) and WCA (lower). Right image shows a Band 4 CCA. A single feedhorn feeds the OMT, splitting the two polarization signals for the 2SB SIS mixers. The LO is generated in the WCA, with the final x2 multiplier on the 110 K stage. The Band 4 cartridges were constructed in Japan at the NAOJ Advanced Technology Center (ATC) in Mitaka.
Figure 4.10: Band 4 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV.

Figure 4.11: System temperature ($T_{sys}$) at zenith for Band 4 with 0.3 to 5.0 mm of PWV. $T_{sys}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.3 Band 5 Receiver

The Band 5 receiver covers the 158 to 211 GHz spectral range (in the 1.9–1.42 mm atmospheric window), slightly extended from the original 163–211 GHz coverage. The signal collected by the telescope is focused to the Band 5 cartridge using a set of cold mirrors (Figure 4.12). A single feedhorn feeds an OMT which splits the two linear polarizations and feeds the 2SB SIS mixers.

Figure 4.12: Optical layout of the Band 5. Placing the mixer assembly along the side of mirror M1 and the streamlined layout of the IF system are the key to fulfil the requirements on $5w_o$ optical beam clearance at the entrance to the cryostat.

A block diagram of the Band 5 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.13. The Band 5 CCA contains a feed horn, an OMT as a polarization splitter, 2SB SIS mixer assemblies, cold IF amplifiers, isolators, and LO frequency doubler. The RF signal is down converted to 4–8 GHz using a 2SB mixer unit.

The atmospheric transmission in Band 5 is shown in Figure 4.14 for different PWV values. The specification for Band 5 receiver noise performance ($T_{rx}$) is <55 K over 80% of the 163-211 GHz range, and <75 K over the whole 158 - 211 GHz extended range (SSB $T_{rx}$). However, the realised performance of the receivers are often better than 50 K over the band. The OT assumes 55 K. The resulting system temperatures ($T_{sys}$) are shown in Figure 4.15.
Figure 4.13: Block diagram of the Band 5 receiver (left) including CCA (upper) and WCA (lower). Right image shows a Band 5 CCA. Note the single feedhorn which feeds the OMT, splitting the two polarization signals for the 2SB SIS mixers. The production design has a room-temperature 6x multiplier in the WCA, then a 2x cold multiplier in the CCA (at 110K).
4.2. THE CYCLE 6 RECEIVERS

Figure 4.14: Band 5 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. The strong absorption due to H$_2$O at 183.31 GHz allows for observations at this frequency only when the PWV is < 1 mm. Continuum observations in Band 5 should avoid this part of the band if possible.

Figure 4.15: Typical system temperature ($T_{sys}$) at zenith for Band 5 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. $T_{sys}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.4 Band 6 Receiver

The Band 6 receiver covers a frequency range of 211–275 GHz (the 1.3 mm atmospheric window). This receiver has a window with a pair of off-axis ellipsoidal mirrors inside the cryostat (Figure 4.16). A single feedhorn feeds an OMT which splits the two linear polarizations and feeds the SIS mixers. A block diagram of the Band 6 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.17.

Figure 4.16: Band 6 cold off-axis ellipsoidal mirrors feeding the single feedhorn. The off-axis beam from the telescope secondary mirror (shown by the dashed line) feeds directly through the cryostat window, and the Cassegrain focus is just inside the inner infrared blocking filter. Note the slightly inclined inner window, designed to minimize standing waves.

The Band 6 allowed IF frequency was 5-10 GHz, but has been extended further for Cycle 6 to allow for multiple simultaneous line observations with more bandwidth per spw\(^1\); it now covers the range 4.5–10.0 GHz. Although there is up to \(\sim 10\% - 20\%\) excess receiver noise below 5.5 GHz due to LO1 (but the increase of \(T_{\text{sys}}\) is less), the CO multi-transition setup is still considerably more efficient than observing each line separately. It is recommended that for continuum observations, the IF range 6–10 GHz is still used. Also, it should be noted that the full range 4.5–10 GHz cannot be completely sampled because of the limited 4 GHz width of the two basebands per polarization.

The atmospheric transmission in Band 6 is shown in Figure 4.18 for three typical PWV values. Most of the narrow absorption lines are from ozone.

The specification for Band 6 receiver noise performance \((T_{\text{rx}})\) is \(< 83\text{ K over 80\% of the band, and }< 136\text{ K over the whole band (SSB } T_{\text{rx}})\). The achieved on-array results are considerably better, typically \(\sim 40–50\text{ K over most of the band. The OT assumes 55 K. The resulting system temperatures (}T_{\text{sys}}\text{) for different PWV values are shown in Figure 4.19.}

\(^1\) Specifically, the \(^{12}\text{CO}/^{13}\text{CO}/^{18}\text{O } J=2-1\) combination at 230.538/220.398/219.560 GHz, which has a minimum separation of 10.14 GHz. In Cycle 6 the lower end of the allowed IF range has been extended to 4.5 GHz in order to cover all three lines with broader spws.
Figure 4.17: Band 6 receiver block diagram, and (right) image of cartridge. Note the OMT used to split the polarizations feeding the two 2SB mixers. The LO around 80 GHz requires an extra ×3 multiplier inside the cryostat. The Band 6 cartridges were built at NRAO, Charlottesville. Note that the IF output range has been extended to cover 4.5-10 GHz; the range shown is the one recommended for continuum observations (see text).
Figure 4.18: Band 6 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. Most of the narrow absorption lines are from ozone.

Figure 4.19: Typical $T_{sys}$ at zenith for Band 6 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV, based on measured values of the receiver temperatures. $T_{sys}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2. **THE CYCLE 6 RECEIVERS**

4.2.5 **Band 7 Receiver**

The Band 7 receiver covers the frequency range 275 to 373 GHz (the 0.85 mm atmospheric window). It has a similar cold optics design as Band 6, but uses a wire-grid polarization splitter instead of an OMT (Figure 4.20). A block diagram of the Band 7 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.21.

![Figure 4.20: Band 7 cold optics arrangement, showing the off-axis ellipsoidal mirrors and the polarization splitter wire grid.](image)

The atmospheric transmission in Band 7 is shown in Figure 4.22 for different PWV values. The original specification of the Band 7 receiver noise temperature was $T_{\text{rx}} < 147$ K over 80% of the range and $< 219$ K over the whole tuning range, except at the upper end of the band (370-373 GHz), where the specifications were $< 300$ K SSB. However, the performance of the receiver as measured in the lab and on the array is considerably better than this, with typically 65 K achieved, in mid-band. The OT assumes 75 K over the whole band. The resulting system temperatures ($T_{\text{sys}}$) for different PWV values are shown in Figure 4.23. Note that the atmospheric transmission (and hence $T_{\text{sys}}$) at frequencies below 300 GHz is considerably better than that of the top half of Band 7; in that respect the performance is closer to that of Band 6. Also to maximise continuum sensitivity, the deep water absorption lines around 325 GHz should be avoided if possible.
Figure 4.21: Band 7 front-end receiver block diagram, and (right) annotated image of the Band 7 cartridge. Note the polarization-splitting grid and LO injection in the cold optics above the mixers. The Band 7 cartridges were built at IRAM in France.
4.2. THE CYCLE 6 RECEIVERS

Figure 4.22: Band 7 atmospheric zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. The deep atmospheric absorption at 325 GHz is due to water, and the narrower feature at 369 GHz is due to oxygen.

Figure 4.23: Typical $T_{\text{sys}}$ at zenith for Band 7 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. $T_{\text{sys}}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.6 Band 8 Receiver

Band 8 covers the frequency range 385 to 500 GHz (the 650 µm atmospheric window). The cryogenic optics of this receiver adopts a single mirror to couple a feed horn in front of an SIS mixer block to the sub-reflector. A single feedhorn feeds an OMT which splits the two linear polarizations and feeds the 2SB SIS mixers (Figure 4.24).

A block diagram of the Band 8 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.25. The Band 8 CCA consists of a cold optics, a feed horn, an OMT, 2SB SIS mixers assemblies, cold IF amplifiers, isolators, and LO frequency sextuplers.

The atmospheric transmission in Band 8 is shown in Figure 4.26 for different PWV values. The specification of the Band 8 receiver noise temperature is $T_{rx} < 196$ K over 80% of the range and <292 K over the whole tuning range. However, the performance of the receiver as measured in the lab and on-array is considerably better than this, with typical values of 70-120 K. The OT assumes 150 K. The resulting system temperatures ($T_{sys}$) for 0.472 mm PWV are shown in Figure 4.27. Note that transmission in the lower part of the band is almost as high as Band 7, whereas at the high-frequency end of the band (e.g. around the CI line at 492 GHz) the transmission is as low as the centre of Band 9.
Figure 4.25: Block diagram of the Band 8 receiver (left) including CCA (upper) and WCA (lower). Right image shows a Band 8 CCA. Note the single feedhorn which feeds the OMT, splitting the two polarization signals for the 2SB SIS mixers. The Band 8 cartridges were constructed in Japan at the NAOJ Advanced Technology Center (ATC) in Mitaka.
CHAPTER 4. RECEIVERS

Figure 4.26: Band 8 atmospheric zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. The atmosphere in the Band 8 frequency range has some deep absorption features due to water and oxygen.

Figure 4.27: Typical $T_{\text{sys}}$ at zenith for Band 8 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. $T_{\text{sys}}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.7 Band 9 Receiver

Band 9 covers the frequency range 602 to 720 GHz (the 450 µm atmospheric window). It uses a wire grid in order to separate the two orthogonal polarizations, as well as to provide the LO injection scheme (Figure 4.28).

![Figure 4.28: Basic Band 9 optics layout. The signal path is symmetrical for the vertical polarization (P0) and the horizontal polarization (P1). For P1 the path is as follows: the telescope focal point (FP) is followed by mirror M3, wire grid polarization splitter, mirror M4, a beam splitter for LO insertion, and finally the mixer horn U1. For P0 the signal follows from FP to the same mirror M3 and then, reflected by the grid, comes to mirror M4, a beam splitter, and mixer horn U1.]

The mixers are double sideband (DSB), and therefore additional techniques must be employed during the observations to either separate the sidebands or reject the unwanted sideband. LO offsetting can be used to reject one of the two sidebands, which can be chosen independently for each spectral window. Note that LO offsetting does not reject the noise from the unwanted sideband, it simply moves any correlated signal to a high fringe rate so that the signal is smeared over a larger bandwidth increasing noise incoherently. In Cycle 6, 90-degree phase switching can also be used to correlate both USB and LSB simultaneously (see Chapter 6).

The IF bandwidth in this receiver is 8 GHz per polarization (7.5 GHz effective bandwidth after the IF Processor units, see Section 6.4), covering 4-12 GHz. A block diagram of the Band 9 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.29.

The Band 9 atmospheric transmission is strongly dependent on the PWV, as illustrated in Figure 4.30, where only the lower values of PWV have significant transmission. The specifications for the receiver are $T_{\text{rx}} < 175$ K over 80% of the band and $< 261$ K over all the band. However, the performance is considerably better than this, with engineering and typical on-array values for $T_{\text{rx}}$ of 65–120 K. The OT assumes 110 K. Figure 4.31 shows the expected $T_{\text{sys}}$ for 0.472 mm of PWV, over most of the band given the OT default expected receiver noise, however, the achieved $T_{\text{sys}}$ is extremely dependent on the line-of-sight PWV. Phase stability also limits when observations can be made, therefore, most observations in Band 9 are done at night and early morning (and more commonly during austral winter). As well as having a lower atmospheric transmission and a less stable
Figure 4.29: Block diagram of Band 9 cartridge (left) and a schematic image (right). Note that there are only two IF outputs, one from each polarization in this DSB receiver. The extra Faraday rotation mirror in the LO system is part of the ALMA fibre-optic Line Length Corrector system (see Appendix B.3.3), and means that the Band 9 receiver must be available on all the antennas for this to work. The Band 9 receiver was built at SRON in the Netherlands.
atmosphere, Band 9 observing provides additional challenges, including finding sufficiently bright calibrators (most QSOs are relatively faint at this frequency), requiring accurate pointing for the relatively small primary beam (although the actual pointing is done in Band 6), and the need for the highest level of stability in the rest of the system.

Figure 4.30: Band 9 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. Most of the band is affected by the wings of pressure-broadened H$_2$O.

Figure 4.31: Typical $T_{sys}$ at zenith for Band 9 with 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. $T_{sys}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
4.2.8 Band 10 Receiver

Band 10 covers the frequency range 787 to 950 GHz (the 350 µm atmospheric window). Band 10 is the highest frequency receiver of the ten bands envisioned for the ALMA front-end system. The development of the Band 10 receiver was extremely difficult and faced many technical challenges from its material selection. Niobium (Nb) superconducting tuning circuits, which are used in other ALMA receiver bands, cannot be used for Band 10 SIS mixers due to large losses from pair-breaking above a superconducting gap frequency of about 700 GHz. Therefore, niobium-titanium-nitride (NbTiN) with a critical temperature of about 15 K, has been utilized in the tuning circuit of Band 10 mixers. The Band 10 Nb/AlOx/Nb tunnel junctions with NbTiN-based tuning circuitry achieved ALMA requirements and the best DSB receiver noise temperature was 125 K, corresponding to about 3 times the quantum limit for 4 K operation.

It uses a wire grid in order to separate the two orthogonal polarizations, as well as to provide the LO injection scheme (Figure 4.32). The mixers are double sideband (DSB), and therefore LO offsetting can be used to reject one of the two sidebands. As in the case of Band 9, Cycle 6 allows 90-degree phase switching to be used to correlate both USB and LSB simultaneously (see Chapter 6). The IF bandwidth in this receiver is 8 GHz per polarization (7.5 GHz effective bandwidth after the IF Processor units, see Section 6.4), covering 4–12 GHz. A block diagram of the Band 10 receiver, including the cold cartridge and warm cartridge assembly, is shown in Figure 4.33.

Figure 4.32: Schematic of ALMA Band 10 optics. ALMA Band 10 optics are composed of two elliptical mirrors, M1 and M2, a wire grid and two corrugated horns. The wire grid is used to separate the two linear polarizations, P0 and P1, and it is located after the two elliptical mirrors to minimize the number of optical components required.

The Band 10 atmospheric transmission is significantly dependent on the PWV, as illustrated in Figure 4.34 for different values of PWV. The specifications for this receiver are $T_{\text{rx}} < 230$ K over 80% of the band and <344 K over all the band. Typical on-array $T_{\text{rx}}$ values are 175–275 K. The OT assumes 230 K. Figure 4.35 shows the expected $T_{\text{sys}}$ for 0.472 mm of PWV, over most of the band given the OT-assumed receiver noise. Phase stability also limits when observations can be made, therefore, most observations in Band 10 will generally only be performed at night and early morning, when the sky is most stable. Also the antenna surface accuracy has a significant effect on the aperture efficiency at this high-frequency band; improvements to the surface accuracy of the 12m antennas during 2017 has doubled the aperture efficiencies in Band 10. As well as having a lower atmospheric transmission and a less stable atmosphere, Band 10 observing provides the most challenges for observing, including finding sufficiently bright calibrators (most QSOs are relatively faint at this frequency), requiring accurate pointing and focus for the relatively small primary beam (although pointing and focus is normally done in lower Bands), and the need for the highest level of stability in the rest of the system.
Figure 4.33: Block diagram of Band 10 cartridge (left) and a schematic image (right). Note that there are only two IF outputs, one from each polarization in this DSB receiver. The Band 10 cartridges were constructed in Japan at the NAOJ Advanced Technology Center (ATC) in Mitaka.
Figure 4.34: Band 10 zenith transmission for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. Transmission of most of the band is affected by the wings of pressure-broadened lines of H$_2$O.

Figure 4.35: Typical $T_{sys}$ at zenith with Band 10 for 0.3, 0.5, 1.0, 2.0 and 5.0 mm of PWV. $T_{sys}$ was computed using only the receiver temperature values adopted in the OT and the atmospheric contribution. No spill-over or background terms have been included. Temperature is given in Kelvin.
Chapter 5

The Correlators

A correlator can be viewed as the virtual focal plane of an interferometer array where all voltage-based signals from all individual antennas are processed to derive both the cross-correlation products from all independent antenna pairs and the autocorrelation function of each antenna. It delivers complex fringe visibilities which, once calibrated in amplitude and phase, allow the ALMA users to synthesize astronomical images. The correlator also provides pre-correlation delay and phase tracking functions to adjust the response to the wavefronts of received signals in order to maintain the coherence of the complex visibilities. Walsh-switching modulations for sideband separation (90° phase switching) and spurious-signal suppression (180° phase switching) are demodulated in the correlator.1

All signals received by ALMA antennas are processed in one of two correlators: the 64-input Correlator (also known as the Baseline Correlator or BLC) and the ACA Correlator. The 64-input Correlator is used primarily for the main 12-m Array, while the ACA Correlator is used for the Morita Array comprised of the ACA 7-m Array and the Total Power Array (i.e. single-dish), respectively2. Both correlators run simultaneously and independently. Thus, while the 12-m Array observes an object using the 64-input Correlator, the ACA Correlator can be used with the 7-m Array and/or the Total Power Array observing either the same or a different object.

Celestial signals received by the antennas are down converted to lower frequency bands using a set of Local Oscillators (LOs) and mixers as described in Appendix B. The outputs from the IF system form four Base Bands (BBs), each covering a bandwidth of 2 GHz in two orthogonal linear polarizations. These analog BB signals are sampled at the sampling frequency of 4 GHz and quantized with eight quantization levels (3 bits per sample) in digitizers, and then transferred via fiberoptic cable to one of the two correlators.

Both correlators generate auto-correlation and cross-correlation products at the same time. The auto-correlation is used not only for TP-Array observations but also for normalization of cross power spectra and measurements of system noise temperatures. The cross-correlations are used for interferometry with the 12-m Array and the ACA 7-m Array, and also for pointing and focus calibrations for all Arrays.

This chapter addresses capabilities of the correlators to be offered for Cycle 6 observations. The specifications of the correlators determine a lot of observational performance such as bandwidth, spectral resolution, time resolution, and polarimetry. The phase tracking performance, including online WVR correction, provides coherence and improves phase stability in correlated data delivered to users. Imperfect correction for the non-linear response in the correlators invokes systematic errors in complex visibilities. The response of digital signal processing is also presented in this chapter.

1Some of these features can be employed outside the correlator. In the case of the ALMA correlators, the phase tracking is taken in the LO1 and LO2. The 180° phase switching is modulated in the LO1 and demodulated in the DTS transmitter after digitization. See also Sections 5.5.4 and B.4.4, and also Emerson 2005, ALMA memo No. 537.

2 Crossbar switching allows for some flexibility in this arrangement.
5.1 The 64-input Correlator

The 64-input Correlator employs a hybrid design, also known as the FXF\textsuperscript{3} system (Escoffier et al. 2007, A&A 462, 801), that increases by a factor of 32 the spectral resolution of its traditional lag (XF) part. It operates in two basic modes, Time Division Mode (TDM) – equivalent to an XF correlator with a wide bandwidth and a coarse spectral resolution for mainly continuum observations, and Frequency Division Mode (FDM) with fine spectral resolutions for spectral-line observations. A simplified overview diagram of the 64-input Correlator is shown in Figure 5.2. It consists of four quadrants, all of which are available for Cycle 6. Each quadrant can handle a 2 GHz dual-polarization BB for up to 64 antennas\textsuperscript{4}. The full set of four quadrants is capable of accepting four BBs to cover a total 8 GHz bandwidth per polarization, that is 16 GHz of instantaneous bandwidth.

5.1.1 TDM mode

TDM is mostly used for continuum observations. Its simplicity, compared with FDM, offers advantages of a lower data rate and better linearity. Therefore, it is used for standard setups such as pointing, focus, delay calibration, system temperature measurements, sideband ratio measurements, etc. TDM provides also a higher time resolution capability as described in Section 5.5.3.

The full 2 GHz BB is directly sent to the correlator bypassing the Tunable Filter Banks (TFBs). The correlator cuts off the least significant bit to reduce quantization levels from 3- to 2-bits per sample (see Section 5.3).

The TDM mode provides a Spectral Window (spw)\textsuperscript{5} that consists of up to 256/\(N_{\text{pol}}\) channels per BB, where

\textsuperscript{3}F, X and F stand for filtering, correlation and Fourier transform, respectively.

\textsuperscript{4}\(N_{\text{ant}}(N_{\text{ant}} - 1) = 2016\) baselines and 64 auto-correlations for \(N_{\text{ant}} = 64\).

\textsuperscript{5}An spw is a continuous spectrum composed of uniformly spaced frequency channels. See also Chapter 6 about the relation between BB and spw.
5.1. THE 64-INPUT CORRELATOR

\[ N_{pol} \] is the number of polarization products per BB\(^6\). As the full 2000 MHz BB is covered, this requires some truncation of the band-edge channels in offline data processing yielding 1875 MHz of usable bandwidth – see Section 6.4.

5.1.2 FDM mode

FDM is adequate for spectral-line observations that require a spectral resolution higher than that of TDM. Each 2-GHz BB in an antenna is processed in a TFB card, where digital filtering, a digital mixer and LO (LO4) are implemented in an FPGA. (FPGA). In a TFB there are 32 independent 62.5-MHz bandpass filters or sub-bands. To avoid edge artifacts, sub-band center frequencies are set such that their edges overlap and the effective bandwidth is thus reduced by a factor of 15/16, that is, 58.59375 MHz of effective bandwidth per sub-band. Down stream, the correlator software stitches together contiguous sub-bands to output a seamless cross power spectrum to form an spw with as many as \( \frac{15}{16} \times 8192/N_{pol} \) channels.

The number of channels in each spw can be reduced by averaging 2, 4, 8, or 16 channels into one to accommodate the maximum data rate. Table 5.1 lists the spectral setups for Stokes I \((N_{pol} = 2)\) observations. The number of valid sub-bands to form an spw is selectable from 32, 16, 8, 4, 2, or 1. That is, spw bandwidths

\(^6\) \(N_{pol} = 2\) for standard Stokes-I observations that employ XX and YY products, and \(N_{pol} = 4\) for full-Stokes observations.
of 1875, 937.5, 468.75, 234.375, 117.1875, and 58.59375 MHz, respectively. It is possible to have multiple (up to 4) spws in the same BB. While different spectral setups can be set for spws in different BBs, all of spws in the same BB must have the same channel spacing.

The center frequency of each spw can be tuned over the 2 GHz-wide BB using the digital LO (LO4) in the TFB card. However, the edges of the full bandwidth of the sub-bands cannot fall outside the 2 GHz BB range, and the frequency tuning is made in steps of 30.5 kHz set by the frequency resolution of the TFB digital mixer. See also Section 5.5.2 for details about spectral setting.

<table>
<thead>
<tr>
<th>Correlator / Mode</th>
<th>Bandwidth (MHz)</th>
<th>Num. of ch. per pol. (in dual-pol.)</th>
<th>Ch. spacing (MHz)</th>
</tr>
</thead>
<tbody>
<tr>
<td>BLC/TDM</td>
<td>2000*</td>
<td>128</td>
<td>15.625</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>1875</td>
<td>3840</td>
<td>0.488</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>938</td>
<td>3840</td>
<td>0.244</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>469</td>
<td>3840</td>
<td>0.122</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>234</td>
<td>3840</td>
<td>0.061</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>117</td>
<td>3840</td>
<td>0.0305</td>
</tr>
<tr>
<td>BLC/FDM</td>
<td>58.6</td>
<td>3840</td>
<td>0.0153</td>
</tr>
<tr>
<td>ACA</td>
<td>2000*</td>
<td>4096</td>
<td>0.488</td>
</tr>
<tr>
<td>ACA</td>
<td>1000</td>
<td>4096</td>
<td>0.244</td>
</tr>
<tr>
<td>ACA</td>
<td>500</td>
<td>4096</td>
<td>0.122</td>
</tr>
<tr>
<td>ACA</td>
<td>250</td>
<td>4096</td>
<td>0.061</td>
</tr>
<tr>
<td>ACA</td>
<td>125</td>
<td>4096</td>
<td>0.0305</td>
</tr>
<tr>
<td>ACA</td>
<td>62.5</td>
<td>4096</td>
<td>0.0153</td>
</tr>
</tbody>
</table>

Table 5.1: Correlator modes and spectral performances per BB for dual parallel polarization (XX and YY). *Usable bandwidth excluding band edges is \(
\sim 1875 \text{ MHz}\).

5.1.3 Correlation and Realtime Processing

The correlation cards perform the multiply-and-add operations to produce the correlation functions at a clock rate of 125 MHz (4 GHz samples demultiplexed by 32). Four quadrants handle four BBs. A quadrant of the correlator consists of 32 planes of \(64 \times 64\) 256-lag correlator circuits, and it yields auto-correlations and cross-correlations for 64 antennas. In FDM there are always 8192 spectral points per BB for both X and Y polarizations, and higher spectral resolutions can be achieved by reducing the number of sub-bands (i.e. bandwidth) used to form an spw. It is possible to set different modes in different quadrants (i.e. BBs); for example, while one BB is set in TDM, other BBs can be set in FDM. For spectral setup details, see Chapter 6. The Long Term Accumulator (LTA; see Figure 5.2) takes short 1 ms or 16 ms integrations from the correlator circuits and provides longer term integration. Further time averaging through multiple dump intervals is performed in the Correlator Data Processor (CDP) computers. See also Section 5.5.3 about time resolution. Section 5.5 describes in detail the correlator data processing.

5.2 The ACA Correlator

The ACA Correlator is dedicated to observations with the Atacama Compact Array that consists of the 7-m Array with twelve 7 m antennas and the Total Power Array (TP Array) with four 12 m antennas. The ACA Correlator is based on the FX design in which the incoming time-domain data stream is converted into frequency-domain spectrum via an FFT (Fast Fourier Transform) module before cross multiplication to form the power spectrum. The FFT part always accepts a full 2 GHz bandwidth and outputs 524288 channel spectra with a frequency resolution of 3.815 kHz. This design enables flexible spectral handling as shown in Figure 5.3

\(^7\) F and X stand for Fourier transform and multiplication, respectively.
producing multiple spws with different channel spacings by averaging multiple channels (spectral binning). Since the spectral resolution function is different from that of the 64-input Correlator, a frequency profile synthesis (FPS) is performed in the Correlator Data Processor computer so that the outputs of both correlators are matched (Kamazaki et al. 2008, ALMA Memo 580).

Figure 5.3: Block diagram of the ACA Correlator. One of four quadrants is shown. Time-series data from each antenna are divided in the time domain and processed in an 8-way parallel stream. FFT is performed in an FPGA delivering signals in the spectral domain which are then multiplied. The auto and cross-correlation data are then accumulated in time, and the parallel streams averaged together in the Correlation and Integration Processor (CIP) module. The correlated spectra are then fed to the ACA-CDP for further accumulation and processing.

A detailed diagram of the signal processing in the ACA Correlator itself is shown in Figure 5.3. Similar to
CHAPTER 5. THE CORRELATORS

Figure 5.4: Two quadrants of the ALMA Compact Array (ACA) correlator installed in the ACA correlator room. Credit: ALMA (ESO/NAOJ/NRAO), S. Okumura

the 64-input Correlator, each quadrant of the ACA Correlator processes a BB pair; so the complete system can process up to sixteen antennas independently from the 64-input Correlator.

After receiving the digitized BB signals in the Digital Transmission System Receiver (DTS-Rx), the DTS-Rx and FFT Processor (DFP) modules compensate for geometrical delays between antennas and perform the $2^{20}$-point FFT that produces a $2^{19}$-point complex spectrum\(^8\) (hereafter, voltage spectrum) for every BB per antenna per polarization, with a channel separation of 3.815 kHz ($= 2$ GHz $\div 524288$ channels) in a 16-bit complex integer form. The 16-bit complex voltage spectra are re-quantized into a 4-bit complex integer and sent to the Correlation and Integration Processor (CIP) modules.

The CIP module trims the required frequency range and multiplies the antenna-based voltage spectra to generate baseline-based cross power spectra that correspond to cross-correlations. Antenna-based power spectra, corresponding to auto-correlation, are also generated in the same way. Cross-polarization power spectra can be optionally produced. The (cross) power spectra are channel-averaged and time-integrated as designated before they are sent to the ACA-CDP computers in the Computing subsystem through optical fibers.

The ACA-CDP performs further spectral processing such as non-linearity correction (see Section 5.6.2), FPS and temporal integration, before data are sent to the archive.

The overall hardware design of the ACA Correlator is shown in Figure 5.5. The ACA Correlator is equipped with high-speed FPGA chips rather than Application Specific Integrated Circuit (ASIC) chips used in the 64-input Correlator. For technical detail, see Kamazaki et al. 2012, PASJ 64, 29.

5.3 Digitizers

A digitizer is a device converting continuous voltage waveforms into quantized voltage levels, sampled at discrete timings, and encoded in a digital format. The ALMA digitizers are located in the antenna back ends (BEs) where they convert BB signals into the ALMA digital format (Recoquillon et al. 2005, ALMA Memo No. 532). The Digital Transmitter (DTX) on the BE transfers the digital signals through optical fibers to the data

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\(^8\)2^{20}\text{-point FFT produces a }2^{20}\text{-point complex voltage spectrum of a double sideband including the image sideband that will be discarded.}
5.4 Online WVR Correction

The Water Vapor Radiometer (WVR; see Appendix A.6 for technical description) is a device mounted on each of the 12 m antennas to measure the amount of Precipitable Water Vapor (PWV) in order to correct for pathlength fluctuations in the troposphere. The WVR data are recorded together with visibilities. The later can be chosen...
to be WVR uncorrected, corrected, or both at the same time. If online correction is selected, then, the correction happens in real-time every dump duration, based on path length fluctuations in WVR data samples from each antenna. However, starting from Cycle 6 the online software will only produce uncorrected visibilities, and the correction will be applied offline in the CASA pipeline. Recording both WVR corrected and uncorrected results impacts data rates and storage, producing only one data stream allows for shorter integration times.

The WVR correction will not be applied in the ACA Correlator because the array is so compact that PWV fluctuation is not significant.

5.5 Capabilities of the Correlators

5.5.1 Polarization

Both correlators offer polarimetry capability by delivering all correlation products of \( \langle XX^* \rangle \), \( \langle XY^* \rangle \), \( \langle YX^* \rangle \), and \( \langle YY^* \rangle \) where \( X \) and \( Y \) stand for two linear orthogonal polarization components. These correlations relate to the Stokes visibilities of \( I \), \( Q \), \( U \), and \( V \) as

\[
\begin{pmatrix}
  I \\
  Q \\
  U \\
  V
\end{pmatrix}
= \frac{1}{2}
\begin{pmatrix}
  1 & 0 & 0 & 1 \\
  \cos 2\psi & -\sin 2\psi & -\sin 2\psi & \cos 2\psi \\
  \sin 2\psi & \cos 2\psi & \cos 2\psi & -\sin 2\psi \\
  0 & -i & i & 0
\end{pmatrix}
\begin{pmatrix}
  \langle XX^* \rangle / (G_X G_X^*) \\
  \langle XY^* \rangle / (G_X G_Y^*) \\
  \langle YX^* \rangle / (G_Y G_X^*) \\
  \langle YY^* \rangle / (G_Y G_Y^*)
\end{pmatrix},
\]

where \( G_X \) and \( G_Y \) stand for complex antenna-based gains in \( X \) and \( Y \) polarizations and \( \psi \) is the parallactic angle. This formulation doesn’t include cross talk, also known as ‘D-terms’, whose calibration will be discussed in Chapter 10.

The 64-input Correlator forms cross-polarization products of \( \langle XY^* \rangle \) and \( \langle YX^* \rangle \) only when polarimetry is required. This is obtained at the expense of less spectral channels. Although the ACA Correlator is capable of polarimetry, the cross-polarization products of the 7-m Array or the TP Array observations will not be offered for Cycle 6.

5.5.2 Spectral Setup

Both correlators are capable of accepting eight 2 GHz bandwidth signal streams consisting of four BBs and two polarizations, and to be configured to have multiple spws within a BB.

Each spw yields a continuous spectrum composed of uniformly spaced spectral channels. Table 5.1 summarizes the maximum number of spectral channels and the channel spacing of the correlators in dual linear polarization mode (\( XX \) and \( YY \)). The maximum number of channels per spw would be doubled (and the channel spacing halved) for single polarization (\( XX \)) observations. Conversely, for full polarization observations (\( XX \), \( XY \), \( YX \), and \( YY \)), the maximum number of channels is halved and the channel spacing correspondingly doubled as compared to Table 5.1. The bandwidth and the number of channels available from the ACA correlator are identical through the application of FPS (Frequency Profile Synthesis). Note that the channel spacing is not the same as the spectral resolution because of the effects of the applied weighting function, as listed in Figure 5.5.2 and shown in Figure 5.5.2. See also Chapter 6 about multiple-spw setup.

The spectral profile which the correlators output is a convolution of the true spectrum with the spectral resolution function shown in Figure 5.5.2. The spectral resolution function is given by Fourier transform of the weighting function applied to the correlation function in the lag domain. The spectral resolution is often characterized by the FWHM (full width at half maximum) of the spectral resolution function.

Uniform weighting in the lag domain results in the sinc function with the FWHM of 1.21× the channel spacing and yields spectral sidelobes. This ‘ringing’ phenomenon affects the spectra when a narrow line, interference spike, or strong edge channels are present. Alternative weighting functions\(^9\) are applicable to suppress

\(^9\)For a full description, see http://mathworld.wolfram.com/ApodizationFunction.html.
the spectral sidelobes as listed in Table 5.2. Note that the sidelobes do not matter in most of astronomical observations where a line profile spreads over several spectral channels. The choice of a weighting function is a trade off between resolutions and sidelobes. The default weighting function is the Hanning function, which gives the spectral resolution of $2 \times$ the channel spacing and the maximum sidelobe level of $-2.6\%$. A proposer may request another weighting function through the OT if necessary for the science goal.

<table>
<thead>
<tr>
<th>Weighting</th>
<th>FWHM (ch)</th>
<th>Max sidelobe</th>
</tr>
</thead>
<tbody>
<tr>
<td>Uniform</td>
<td>1.21</td>
<td>-0.217</td>
</tr>
<tr>
<td>Hanning</td>
<td>2.00</td>
<td>-0.026</td>
</tr>
<tr>
<td>Hamming</td>
<td>1.82</td>
<td>+0.007</td>
</tr>
<tr>
<td>Bartlett</td>
<td>1.77</td>
<td>+0.047</td>
</tr>
<tr>
<td>Blackmann</td>
<td>2.30</td>
<td>+0.001</td>
</tr>
<tr>
<td>Welch</td>
<td>1.59</td>
<td>-0.086</td>
</tr>
</tbody>
</table>

Figure 5.7: (left table) Spectral resolutions (FWHM) and the maximum sidelobe levels for weighting functions. ALMA employs the Hanning window in default. (right figure) Spectral resolution functions corresponding to various weighting functions

Spectral channel averaging is available to bin or average spectral channels in the CDP. Channels can be averaged together; factors of $N = 2, 4, 8,$ or 16 are available. The main purpose is to reduce the data rate to the archive and the total data volume. It provides a broader spread of correlator functionality between the current TDM (which has only 128 channels in dual polarization) and full FDM (with 3840 channels in dual polarization mode). It might be quite acceptable for those who need something with more resolution than TDM, but where the FDM channels at the full resolution are unnecessary. Table 5.2 shows the resolutions (in kHz) for different values of $N$, using Hanning weighting, in the different bandwidth modes. The channel spacings are in parentheses. In Cycle 6, the default value is $N = 2$.

Note that the default Hanning window function gives a resolution 2 times the channel spacing, so using $N = 2$ (cutting the number of spectral channels from 3840 to 1920) results in negligible loss of final resolution. It is recommended, unless the maximum spectral resolution is required by the observations, to reduce the number of channels when feasible. This is selected in Phase 2 of the SB creation. However, note that this is a non-reversible operation. For Cycle 6 the default value for spectral averaging will be 2. The proposer must specifically request and justify the use of $N = 1$ for spectral averaging.

The ACA Correlator is an FX correlator where the uniform weighting in the FFT segment is equivalent to the Bartlett weighting in the lag domain. FPS in the CDP can be applied to match the spectral resolution function to that of the 64-input Correlator. The bandwidth and the number of channels in the ACA correlator remain unchanged through FPS since Cycle 4.
CHAPTER 5. THE CORRELATORS

<table>
<thead>
<tr>
<th>Usable bandwidth (MHz)</th>
<th>$N = 1$</th>
<th>$2^*$</th>
<th>4</th>
<th>8</th>
<th>16</th>
</tr>
</thead>
<tbody>
<tr>
<td>Channels = 3840</td>
<td>977 (488)</td>
<td>1129 (977)</td>
<td>1938 (1953)</td>
<td>3904 (3096)</td>
<td>7813 (7812)</td>
</tr>
<tr>
<td>937.5</td>
<td>488 (244)</td>
<td>564 (488)</td>
<td>969 (977)</td>
<td>1952 (1953)</td>
<td>3906 (3096)</td>
</tr>
<tr>
<td>468.8</td>
<td>244 (122)</td>
<td>282 (244)</td>
<td>485 (488)</td>
<td>976 (977)</td>
<td>1953 (1953)</td>
</tr>
<tr>
<td>234.4</td>
<td>122 (61)</td>
<td>141 (122)</td>
<td>242 (244)</td>
<td>488 (488)</td>
<td>977 (977)</td>
</tr>
<tr>
<td>117.2</td>
<td>61 (31)</td>
<td>71 (61)</td>
<td>121 (122)</td>
<td>244 (244)</td>
<td>488 (488)</td>
</tr>
<tr>
<td>58.6</td>
<td>31 (15)</td>
<td>35 (31)</td>
<td>61 (61)</td>
<td>122 (122)</td>
<td>244 (244)</td>
</tr>
</tbody>
</table>

Table 5.2: Spectral resolution and channel spacing (in parentheses) in kHz for different correlator bandwidth modes (left column) and for different spectral channel averaging factors (columns, $N = 1$ to 16), using Hanning smoothing. The number of channels can be reduced from 3840 (for the unaveraged case, $N = 1$) down to 240 (for $N = 16$). As $N$ increases, the spectral resolution functions of adjoining channels are combined and then the spectral resolution approaches the channel spacing. Values are given for the dual polarization case. For Cycle 6 the default value for spectral averaging will be 2.

Figure 5.8: Basic data processing and accumulation/averaging steps between the correlator and archive. The CDP applies channel truncation, spectral averaging, and time averaging for the channel-averaged data. It also applies online WVR correction to produce the WVR-corrected spectral data. They are packed in Binary Data Format (BDF) and are archived together with the WVR data.
5.5. CAPABILITIES OF THE CORRELATORS

5.5.3 Time Resolution

As shown in Figure 5.8 there are three stages where correlated data are accumulated or averaged within a correlator sub-system:

1. Accumulator #1: implemented in the correlator itself. Data are accumulated without applying a proper averaging factor. The accumulation period is called dump duration.

2. Accumulator #2: implemented in the CDP, this stage first truncates ~5% band edges for each spw and then averages the whole spectral channels into one. Those channel averages are then further averaged in time during a number of dump duration periods. This average period is called channel average duration, which must be an integer multiple of the dump duration.

3. Accumulator #3: implemented in the CDP, this stage averages full resolution spectra during a number of of dump duration periods. This average period is called integration duration, which must be an integer multiple of the channel average duration. This also outputs WVR-corrected spectral data.

Most of the data produced by the correlator corresponds to full resolution spectra, and its data rate can be approximated by the following expression:

\[ R_{\text{int}} = N_{\text{ant}}^2 \times N_{\text{BB}} \times N_{\text{ch}} \times N_{\text{pol}} \times N_{\text{byte}} / T_{\text{int}}, \quad \text{(bytes/sec)} \]  

(5.2)

where \( N_{\text{ant}}, N_{\text{BB}}, N_{\text{ch}} \) are the number of antennas, BBs, and total number of spectral channels, respectively. \( N_{\text{pol}} = 1, 2, \) and 4 for single, dual and full polarization products, respectively. \( N_{\text{byte}} \) (bytes/visibility) is the byte size of the real or imaginary part of a visibility, and \( T_{\text{int}} \) is the integration duration in the Accumulator #3. For example, an array configuration with \( N_{\text{ant}} = 42 \), a spectral setup with \( N_{\text{BB}} = 4, N_{\text{ch}} = 3840 \times 2 \) polarization products (\( XX \) and \( YY \)), \( T_{\text{int}} = 6 \) sec, and assuming \( N_{\text{byte}} = 2 \) yields a data rate of 18.1 MB/sec.

In Cycle 6, ALMA regulates \( R_{\text{int}} \) at 70 MB/sec. The ALMA Observing Tool (OT) estimates the data rate for each science goal and warns if it exceeds the limit. In this case, it is required to consider the optimal trade between time and spectral resolutions, number of polarization products, number of BBs, and number of antennas in the array.

Some specific considerations about different time duration follows:

- **Integration duration of spectral data with the 64-input Correlator**
  
  The basic integration time of the 64-antenna Correlator is 16 and 1 ms for cross- and auto-correlation, respectively. The correlator integration duration can be, in TDM and FDM, any value that would keep the integration data rate \( R_{\text{int}} \) under the regulated limit for a given science cycle. Note that the 64-input Correlator sets \( N_{\text{byte}} \) in real-time, based on the actual magnitude of the data produced (correlation regime). For validation purposes, and only in FDM, a favorable condition of \( N_{\text{byte}} = 2 \) is assumed.

- **Integration duration of spectral data with the ACA Correlator**
  
  The ACA Correlator uses \( N_{\text{byte}} = 4 \) and its dump time is a multiple of 1 ms and 16 ms for autocorrelation and cross-correlation, respectively. The maximum data rate from the ACA Correlator is 3.6 MB/s, independently of the number of antennas. However, the average data rate to the archive will be lower during typical observing modes, because of overheads and the use of TDM in calibration scans.

- **Channel-averaged duration**
  
  The average power spectra across the whole bandwidth, in each spw, can be recorded in an independent spw consisting of only one spectral channel. This single channel spw, without bandpass calibration, is not used for astronomy, but for recording phase and amplitude variations through the observing time. Its duration should normally be equal to or shorter than the WVR integration duration.

- **Integration duration of online WVR Correction**
If online WVR corrections are requested by a given spectral configuration, then the integration duration \( T_{int} \) must be equal to or longer than that of the WVR system (1.152 sec). Given that for Cycle 6 neither the 7-m Array nor the TP Array uses online WVR, the above restriction applies to the 64-input correlator only.

A pulsar gating function is not implemented in the ALMA correlators. Timing sensitive analysis must be taken offline under the time resolution limited by the actual integration duration, \( T_{int} \), used during an observation.

![Figure 5.9: Schematic diagram of phase switching. The LO1 signal phase is modulated with the the Walsh series pattern and demodulated in the DTS before cross correlation. The Walsh function series are employed to generate switching patterns that are orthogonal to each other.](image)

**5.5.4 Phase Switching**

The 180° phase switching capability is implemented to suppress spurious signals and IF cross talks between the LO1 and the digitizers\(^\text{10}\). The 90° phase switching capability is offered to separate, in real-time, upper and lower sidebands of DSB receivers (Bands 9 and 10). The expedient is based on 90° phase switching at the LO1, and accumulation in different memory areas (bins) in the correlator\(^\text{11}\). Walsh function series are employed to make antenna-based switching patterns orthogonal to each other. The Walsh sequences permutes every 16 ms, and the full cycle completes after 2048 ms to achieve orthogonality for all antennas in the array. The dump duration, during a side band-separation observation, must be an integer number of the complete Walsh cycle. See also Appendix B.4.3.

**The 180° phase switching**

The phase switching capability (see principles in Figure 5.9) is realized by phase modulation in the LO1 and demodulation in the DTS transmitter module after digitization.

**The 90° phase switching**

This function aims to separate USB and LSB for the DSB receivers of Bands 9 and 10. When the LO phase of antenna 2 is offset by 90°, the phases of USB and LSB signals in IF will be \(-\phi_{LO} - \frac{\pi}{2}\) and \(\phi_{LO} + \frac{\pi}{2}\), respectively. The 90°-demodulation is taken by swapping real and imaginary parts of the cross correlations in the CDP. After demodulation, the sign of LSB correlation will be inverted with respect to USB. Thus, the average and the difference of cross correlations in two states gives the USB and the LSB visibilities, respectively.

\(^{10}\)Note that this scheme doesn’t work for autocorrelations.

5.6. PRACTICAL PERFORMANCE

<table>
<thead>
<tr>
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<tbody>
<tr>
<td>0</td>
<td>1</td>
<td>$\phi_{LO}$</td>
<td>$S_1 e^{-i\phi_{LO}} + N$</td>
<td>0</td>
<td>$\langle S_1 \cdot S_2^* \rangle + \langle N \cdot N^* \rangle$</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>$\phi_{LO}$</td>
<td>$S_2 e^{-i\phi_{LO}} + N$</td>
<td>0</td>
<td>$\langle S_1 \cdot S_2^* \rangle + \langle N \cdot N^* \rangle$</td>
</tr>
<tr>
<td>1</td>
<td>1</td>
<td>$\phi_{LO}$</td>
<td>$S_1 e^{-i\phi_{LO}} + N$</td>
<td>0</td>
<td>$\langle S_1 \cdot S_2^* \rangle - \langle N \cdot N^* \rangle$</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>$\phi_{LO} + \pi$</td>
<td>$S_2 e^{-i(\phi_{LO} + \pi)} + N$</td>
<td>$\pi$</td>
<td>$\langle S_1 \cdot S_2^* \rangle - \langle N \cdot N^* \rangle$</td>
</tr>
</tbody>
</table>

average $\langle S_1 \cdot S_2^* \rangle$

Table 5.3: Cross correlation of a baseline under $180^\circ$ phase switching. The received signal of antenna 2, $S_2$, is modulated by $180^\circ$ by the LO1 and demodulated in the DTS transmitter module at each antenna. Spurious signals, $N$, in the IF or cross talks are not modulated via the LO1 and thus are canceled via demodulation and accumulation of data in states 0 and 1.

<table>
<thead>
<tr>
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<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>1</td>
<td>$\phi_{LO}$</td>
<td>$S_1 e^{-i\phi_{LO}} + S_1 L e^{i\phi_{LO}}$</td>
<td>0</td>
<td>$\langle S_1.U \cdot S_2^* U \rangle + \langle S_1.L \cdot S_2^* L \rangle$</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>$\phi_{LO}$</td>
<td>$S_2 e^{-i\phi_{LO}} + S_2 L e^{i\phi_{LO}}$</td>
<td>0</td>
<td>$\langle S_1.U \cdot S_2^* U \rangle + \langle S_1.L \cdot S_2^* L \rangle$</td>
</tr>
<tr>
<td>1</td>
<td>1</td>
<td>$\phi_{LO}$</td>
<td>$S_1 e^{-i\phi_{LO}} + S_1 L e^{i\phi_{LO}}$</td>
<td>0</td>
<td>$\langle S_1.U \cdot S_2^* U \rangle - \langle S_1.L \cdot S_2^* L \rangle$</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>$\phi_{LO} + \pi$</td>
<td>$S_2 e^{-i(\phi_{LO} + \pi)} + S_2 L e^{i(\phi_{LO} + \pi)}$</td>
<td>$\pi$</td>
<td>$\langle S_1.U \cdot S_2^* U \rangle - \langle S_1.L \cdot S_2^* L \rangle$</td>
</tr>
</tbody>
</table>

(State 0 + State 1)/2 $\langle S_1.U \cdot S_2^* U \rangle$

(State 0 − State 1)/2 $\langle S_1.L \cdot S_2^* L \rangle$

Table 5.4: Cross correlation of a baseline under $90^\circ$ phase switching. The suffixes of 1, 2, U, L stand for antenna 1 and 2, and USB and LSB, respectively. The LO phase of antenna 2 is modulated by $90^\circ$. After $90^\circ$-demodulation and correlation, the cross correlations of USB and LSB have opposite signs. Thus, the average and the difference of cross correlations in two states gives the USB and the LSB visibilities, respectively.

5.6. PRACTICAL PERFORMANCE

5.6.1 Sensitivity

A digitizer adds quantization noise to its input analog signal, with a consequent signal-to-noise reduction or sensitivity loss. The ALMA digitizer employs 3-bit (8-level) quantization\(^\text{12}\), and additional re-quantization processes are applied in the correlators.

In Cycle 6, the 64-input Correlator TDM mode makes use of 2-bit quantization that yields a quantization efficiency of $\eta_Q = 0.88$. In the FDM case, digital filtering is applied to 3-bit quantized signals followed by 2-bit re-quantization. The two stages of quantization would cause the net quantization efficiency of $\eta_Q = 0.90 \times 0.88 = 0.85$ if they were independent (see mode chart tables of Escoffier et al., ALMA Memo 556), nevertheless, the real efficiency can be slightly better than the multiplication (Iguchi et al. 2005, PASJ 57, 259)\(^\text{14}\).

The ACA Correlator feeds 3-bit quantized signals into the FFT processors where butterfly arithmetic is performed with 16-bit precision. The FFT output spectrum is re-quantized into 4-bit (16 levels) before cross multiplication. The combination of these quantizations results in a loss of 4.8% (i.e. $\eta_Q = 0.952$; Kamazaki et al. 2012, PASJ 64, 29).

The quantization efficiencies above are theoretical values under optimal conditions of the input signal level and threshold voltages. Although the signal level is adjusted for each scan, significant variation of power level (e.g. observation of strong sources, unstable weather, or a frequency band near atmospheric absorption lines) can violate the conditions required to optimize the quantization efficiency.

\(^{12}\)The quantization efficiency would be $\eta_Q = 0.96$ if optimal signal levels, threshold voltages, and perfect signal processing were performed (Thompson 1998, MMA Memo 220).

\(^{14}\)It is reported that the combination of 2-bit quantization, digital filtering, and 2-bit re-quantization yields $\eta_Q = 0.81$, better than $0.88 \times 0.88 = 0.77$. 
5.6.2 Linearity

Digital quantization also causes non-linearity in the visibility measurements. The relationship between the correlation coefficient of quantized signals with that obtained for analog signals (i.e. infinite quantization levels) is known as the Van Vleck relationship, which is used for the non-linearity correction.

![Graph showing Van Vleck relationship](image)

Figure 5.10: Van Vleck relationship of 2-, 3-, and 4-bit quantizations at the optimal signal power level and threshold voltages. (Left): Correlation coefficient of digitized signals as a function of analog correlation coefficients, \( \rho \). (Right): Departures of the Van Vleck relationship from linear relation using the single factors of \( \eta_4 = 0.881 \), \( \eta_8 = 0.963 \), and \( \eta_{16} = 0.988 \) for 2-, 3-, and 4-bit quantizations, respectively.

Figure 5.10 shows the Van Vleck relationships in 2-, 3-, and 4-bit quantizations. The correlation coefficients of quantized signals keep adequate linearity for small correlation coefficients\(^{14}\). This indicates that for most sources the cross-correlation response is effectively linear.

However, auto-correlations are affected by the non-linearity because the correlation coefficient at zero lag must be unity by definition. The non-linearity in auto-correlation power spectra influences system temperature measurements and the TP Array observations. In both correlators, the CDPs measure total power of digitized signals and determine the Van Vleck relationship to correct the (cross) power spectra.

5.7 Final data product - the ASDM

The final product from each observation in the archive is known as the ASDM (the ALMA Science Data Model), each of which has an unique hexadecimal name (e.g. uid://A002/X2fed6/X3f). The ASDM contains the metadata (headers, descriptions of the observation setup, ancillary data, etc), and the binary data (the raw data itself), and is described in more detail in Section 12.2. The following describes the spectral data in the ASDM and how it related to the correlator output.

In the ASDM, the binary data are saved in a structure of spws. All of the data in a single spw must share the same frequency setup, including the number and width of spectral channels, and the integration time. The observed spws will be a combination of the science spectral windows set up by the proposers in Phase 1 of the OT, and additional spws from observations needed for calibration (pointing, and sometimes system temperature) set up during Phase 2. Additionally, the WVR data are stored in a spectral window with 4 channels around the water line at 183 GHz. In the ASDM, except for the WVR spectral window, each requested spectral window maps into two output spws in the data: one with the requested dimensions of

\[^{14}\text{The departure from linear relation is } \delta \rho / \rho < 10^{-3} \text{ for } \rho < 0.2\]
channels per polarization product (for example 128 or 3840), and a second channel averaged version with one channel per polarization product (this averaging is done in the correlator). The channel averaged data are used by the online telescope calibration system (TelCal) and for realtime diagnostic purposes (QuickLook), and are typically not used downstream in the data reduction. Overall, this can lead to ASDMs with a large number of spectral windows. For example, a typical science observation in FDM mode can have more than 25 spectral windows. Every scan/spw combination has an “intent” associated with it that indicates its purpose (pointing, system temperature, science, etc). The intent can be used in CASA to decode how to utilize each spectral window in the data reduction process.
Chapter 6

Spectral Setups

This chapter describes the frequency setup of ALMA. In particular it shows how spectral setups are defined from the user viewpoint and how these then set up the ALMA hardware. It summarises the functions and operation of the local oscillator (LO) and Intermediate Frequency (IF) systems, as well as outlining some of the other points pertaining to the frequency setup. For those interested in more details of the LO and IF components and how the hardware works, please jump to Appendix B. The main change for Cycle 6 is the extension of the Band 6 IF coverage to 4.5–10 GHz, allowing for broader spws covering the J=2-1 $^{12}$CO transition and its $^{13}$C and $^{18}$O isotopologues.

To the ALMA system, a spectral setup includes the hardware LO, IF and correlator settings, such that each Spectral Window (spws) covers the desired lines and/or continuum frequencies. To the end-user, however, the spectral setup is normally defined in the Observing Tool (OT) just in terms of the desired lines or observing frequencies, spectral window bandwidths and spectral resolutions: there is normally no need for the user to worry about the details of each hardware setting. For full details of the OT and how to use it, see the user manual and reference manuals, available from the ALMA website (and also in the OT itself). In this chapter, we try to reconcile these two viewpoints.

6.1 Introduction

ALMA uses two stages of heterodyne conversion to shift the signals from the sky or observing frequency down to a range where electronics can be used to perform the digital sampling and cross-correlations. In general, the signals of observing frequency $f_{\text{sig}}$ are mixed to IF signals of frequency $f_{\text{IF}}$ using a narrow Local Oscillator signal (LO, at $f_{\text{LO}}$). $f_{\text{IF}}$ is the difference frequency between $f_{\text{LO}}$ and $f_{\text{sig}}$, where $f_{\text{IF}}$ is always, for ALMA, much lower than $f_{\text{sig}}$. The output signal is split into two separate sidebands at frequencies offset from $f_{\text{LO}}$, known as the upper and lower sideband (USB or LSB).

Figure 6.1 illustrates the ALMA signal path and basic principle of operation. The first heterodyne stage uses the SIS mixer and LO1 to mix (or downconvert) the astronomical signal to 4 GHz- to 8 GHz-wide IF bands, giving both USB and LSB. Up to four 2 GHz wide Basebands (or BBs) can be placed in the available IF range in either or both the USB or LSB IFs using the second stage of downconversion. Within each of these basebands it is possible to place up to four Spectral Windows (spws). Each spws forms a final contiguous spectrum, with bandwidths from 58.594 MHz up to 1.875 GHz wide (see Section 6.3).

Using the OT during both the proposal preparation (Phase 1) and scheduling block (SB) creation (Phase 2), the user adds spws to each baseband by choosing the frequencies (or spectral lines) to be observed in the BB, along with the required bandwidth or spectral resolution. During the subsequent spectral tuning, there are effectively four different LOs set up by the system (Figure 6.1):

\[^1\text{http://almascience.org/documents-and-tools/}\]
• LO1 which sets the frontend tuning frequency. This has continuous coverage.
• LO2 which positions the basebands within the frontend receiver IF (each baseband uses a different LO2, hence there are four LO2s). This does not have continuous frequency coverage, as it is generated by a 125 MHz frequency comb plus a 20–42.5 MHz continuous frequency synthesizer.
• LO3 which is the clock frequency of the digitizers (fixed at 4 GHz).
• LO4 (also known as the tunable filterbank LO, or TFB LO), which is a digital LO synthesized in the correlators allowing positioning of the spectral windows within each baseband. Each spws has effectively a separate LO4.

Figure 6.1: Summary block diagram of ALMA signal path and LO system for one spectral window, one baseband, one polarisation. The Backend IF system is located in each antenna, and the digitised signal for each baseband/polarisation is fed to the correlator in the central technical building through buried optical fibers. See Appendix B for a more detailed description of the components and the other tasks that the LO systems perform.

The OT and the realtime system have a tuning algorithm that attempts to find the best tuning solution for LO1, LO2 and LO4 based on the requested observing frequencies. Because the LOs themselves are generated by combinations of frequency synthesizers, several different tuning solutions may be possible, and the algorithm picks the best one. If no solution is possible, the OT will notify the user. A more detailed description of the LO operation is given in Appendix B.2.

2Note that with tuning setups having multiple basebands, only the hardware LO1 and LO2, the requested frequencies may not appear in exactly the center channel in all the spws, because of the non-continuous coverage of LO2. Normally the tuning algorithm will use LO4 to compensate for this and re-center the lines. However in the widest bandwidth modes (FDM 1875 MHz and TDM) the spws cannot be shifted using LO4. Then the sum of the absolute frequency offsets in all basebands can be as much as 26.38 MHz for 4-baseband setups, giving a worst-case average baseband offset of 6.6 MHz; this is considered small compared with the bandwidth (and the frequency labelling will still be correct). See section B.2 and "ALMA LO System Setup Algorithms" (Scott, 2009) for more details.
6.2. FREQUENCY DEFINITIONS

Cycle 6 has similar restrictions on the spectral setups and tuning as Cycle 5 (see Table 6.2 in Section 6.9), except for Band 6, where the lower end of the IF tuning range has been extended down to 4.5 GHz. The main limits are that the edges of the 2 GHz-wide basebands cannot lie outside the receiver tuning range listed in Table 4.1, and the edges of the (untruncated) spectral windows cannot lie outside the individual 2 GHz-wide basebands. The settings of each baseband in the correlator are independent, so the bandwidth, resolution and correlator mode as well as the value of LO4 can be different for each baseband. For example, it is possible to have a 58 MHz-wide spws centered on a particular line in baseband 1, and simultaneously use a broadband time division multiplexing mode (or TDM mode - see correlator chapter) for all the other basebands. It is possible to have multiple FDM spectral windows within each baseband, described in Section 6.3.2 below; however, currently the underlying resolution of all spws within the same baseband has to be equal. The spectral setups of the ACA Correlator and the 64-input Correlator are - to the end user - effectively the same, with the ACA Correlator having the same allowable spectral functions as the 64-input Correlator.

6.2 Frequency Definitions

There are several frequency definitions used in the ALMA system, in the OT Phase 1 and Phase 2. Most important to the user are:

**Center Frequency (Rest)** \( f_{\text{spws}} \) This is set by the user in the OT Phase 1, and is the central frequency of the spws in the requested rest frame of the source. Note that the source rest frame is selectable, but is commonly set to the local kinematic standard of rest (LSRK), which is the conventional local standard of rest based on the average motion of the Sun with respect to the solar neighborhood (see Section 6.8).

**Center Frequency (Sky)** This is seen in the OT Phase 1, and is the actual central frequency of the spws in the local ALMA (i.e. sky) velocity frame, after including the velocity of the source. However, it does not include the velocity shift of the chosen coordinate frame with respect to ALMA (i.e. the extra velocity from Earth rotation and orbit) as this is not known until actual runtime.

**Baseband desired Center Frequency** \( f_{BB} \) This is the frequency center of the baseband, shown in OT Phase 2. Both the Sky and Rest Baseband center frequencies are seen. If the spws are centered on the basebands, these will be the same as the Phase 1 Center Frequencies (above). Otherwise they can be different because of the TFB LOs (see below).

**Center Offset Frequency** \( f_{\text{offset}} \) Used in the OT Phase 2, this is the offset due to the TFB LO. The center frequency of the spws, \( f_{\text{spws}} \), is given in terms of this and Baseband Center Frequency by:

\[
    f_{\text{spws}} = f_{BB} \pm (f_{\text{offset}} - 3.0\text{GHz})
\]

where the sign depends on the observing sideband.

6.3 Spectral Setups

The wide IF bandwidth and tuning ability allows for routine simultaneous imaging of multiple lines at different spectral resolutions in different basebands. Some examples (with the approximate line frequencies in GHz) are shown in Table 6.1 (for a source of redshift zero). Note that in many cases the lines will not necessarily appear in the center of the spws (for example, in the Band 6 combination 6b). When <4 spectral windows are required for the primary lines, the other BBs can be set up to cover fainter lines or to observe the continuum, potentially in TDM mode. If lines are the main goal in one spws, the other continuum spws frequencies need to be setup manually; one method of doing this is described in the next section. The selection of secondary lines to be observed can be done using the OT spectral interface. In the case of continuum spws, to maximize the sensitivity, the widest bandwidth mode should be chosen (i.e. in TDM, or 1.875 GHz in FDM - preferably with some channel-averaging, see Section 6.4). Also the continuum spws should cover as much of the IF band as possible. Not only will this maximize the continuum SNR on the science target, but these continuum data
can be used to improve phase and amplitude calibration. For this reason it is recommended that all four available basebands be used in all observations. Note that there is no gain in sensitivity by overlapping spws in frequency, since the input signal is the same in the overlap region; the OT uses the unique aggregate bandwidth in calculating the sensitivity, taking any spws overlap into account. However, for sources with known high line density (1 spectral feature per 10 MHz), PIs are encouraged to set up all the spectral windows in FDM mode. This will allow a more robust determination of the line-free channels used to form the aggregate continuum during data processing and imaging.

### 6.3.1 Observing Frequencies for Continuum

For optimum sensitivity, continuum spws should be set to the frequencies with the lowest system temperature. Because the mixer frequency responses are fairly flat, normally this corresponds to the best atmospheric transmission. For full continuum observations (without any required lines, or with only low resolution) the OT has standard optimized frequencies for each band which are used in the OT Phase 1 when continuum is selected; these are noted in Table 6.1 and illustrated by the red in Figure 6.2.

![Figure 6.2: ALMA bands available in Cycle 6 (Bands 3 - 10), showing the frequency of the standard continuum settings as red shading. This gives the coverage of both USB and LSB, except for Bands 9 and 10, which have 8 GHz of bandwidth using the USB only. The available frequency coverage of each band is shown by the grey shading, and atmospheric transmission for 0.6 mm of PWV is given by the blue line.](image)

If mixed line+continuum operation is desired, for example with a single line in BB1 observed at frequency
### Table 6.1: Examples of spectral setups possible in Cycle 6. This includes the standard continuum-only setups for each band, and some multiple line/continuum configurations. Notes:

1. Frequency for standard continuum setups are the mean observing frequency of all spws, $\approx$LO1 in dual-sideband (2SB) receivers, so this frequency is not actually covered if both sidebands are used. See Figure 6.2.

2. * Based on a template spectral setup released with Cycle 6 version of the OT.

<table>
<thead>
<tr>
<th>Band/setup</th>
<th>Species/transition</th>
<th>Freq. (GHz)</th>
<th>Sideb.</th>
<th>bandwidth</th>
<th>BB spws</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td>3 a Standard cont.</td>
<td>HCO$^+$ (1-0)</td>
<td>97.5</td>
<td>dual</td>
<td>TDM</td>
<td>LO1=97.5</td>
<td></td>
</tr>
<tr>
<td>3 b *</td>
<td>HCO$^+$ (1-0)</td>
<td>89.188</td>
<td>LSB</td>
<td>58 MHz</td>
<td>1 1</td>
<td>HCO$^+$/HCN/H$_2$CO</td>
</tr>
<tr>
<td>-</td>
<td>HCN (1-0)</td>
<td>88.632</td>
<td>LSB</td>
<td>58 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>CH$_3$OH (7-2(6)-7(1,6))</td>
<td>101.293</td>
<td>USB</td>
<td>125 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>H$_2$CO (6(1,5)-6(1,6))</td>
<td>101.333</td>
<td>USB</td>
<td>125 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>101.3</td>
<td>USB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>3 c *</td>
<td>CO (1-0)</td>
<td>115.271</td>
<td>USB</td>
<td>58 MHz</td>
<td>1 1</td>
<td>CO/CN/C$^{18}$O</td>
</tr>
<tr>
<td>-</td>
<td>CN (N=1-0)</td>
<td>113.499</td>
<td>USB</td>
<td>58 MHz</td>
<td>1 2</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>C$^{18}$O (1-0)</td>
<td>112.359</td>
<td>USB</td>
<td>117 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>102.5</td>
<td>LSB</td>
<td>TDM</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>100.5</td>
<td>LSB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>4 a Standard cont.</td>
<td>145.0</td>
<td>dual</td>
<td>TDM</td>
<td>TDM</td>
<td>LO1=145.0</td>
<td></td>
</tr>
<tr>
<td>4 b</td>
<td>CS (3-2)</td>
<td>146.969</td>
<td>LSB</td>
<td>117 MHz</td>
<td>1 1</td>
<td>CS/DCO$^+$</td>
</tr>
<tr>
<td>-</td>
<td>DCO$^+$ (2-1)</td>
<td>144.077</td>
<td>LSB</td>
<td>117 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>SO2 (4(2,2)-4(1,3))</td>
<td>146.605</td>
<td>LSB</td>
<td>117 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>H$_2$CO (2(0,2)-1(0,1))</td>
<td>145.603</td>
<td>LSB</td>
<td>117 MHz</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>5 a Standard cont.</td>
<td>203.0</td>
<td>dual</td>
<td>TDM</td>
<td>TDM</td>
<td>LO1=203.0</td>
<td></td>
</tr>
<tr>
<td>5 b</td>
<td>H$_2$O (3(1,3)-2(0,0))</td>
<td>183.31</td>
<td>USB</td>
<td>250 MHz</td>
<td>1 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>181.4</td>
<td>USB</td>
<td>TDM</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>H$_2$S (1(1,0)-1(0,1))</td>
<td>168.763</td>
<td>USB</td>
<td>250 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>170.8</td>
<td>LSB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>6 a Standard cont.</td>
<td>233.0</td>
<td>dual</td>
<td>TDM</td>
<td>TDM</td>
<td>LO1=233.0</td>
<td></td>
</tr>
<tr>
<td>6 b *</td>
<td>$^{13}$CO (2-1)</td>
<td>230.538</td>
<td>USB</td>
<td>937 MHz</td>
<td>1 1</td>
<td>J=2-1 CO isotopes</td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>234.0</td>
<td>USB</td>
<td>TDM</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>C$^{18}$O (2-1)</td>
<td>219.560</td>
<td>LSB</td>
<td>937 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>$^{13}$CO (2-1)</td>
<td>220.399</td>
<td>LSB</td>
<td>937 MHz</td>
<td>3 2</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>216.5</td>
<td>LSB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>7 a Standard cont.</td>
<td>343.5</td>
<td>dual</td>
<td>TDM</td>
<td>TDM</td>
<td>LO1=343.5</td>
<td></td>
</tr>
<tr>
<td>7 b *</td>
<td>continuum spws</td>
<td>343.0</td>
<td>LSB</td>
<td>TDM</td>
<td>1 1</td>
<td>CO/HCO$^+$/HCN</td>
</tr>
<tr>
<td>-</td>
<td>$^{13}$CO (3-2)</td>
<td>345.796</td>
<td>LSB</td>
<td>469 MHz</td>
<td>2 1</td>
<td>1/2-corr</td>
</tr>
<tr>
<td>-</td>
<td>HC$^{15}$N (4-3)</td>
<td>344.200</td>
<td>LSB</td>
<td>234 MHz</td>
<td>2 2</td>
<td>1/4-corr</td>
</tr>
<tr>
<td>-</td>
<td>H$^{13}$CN (4-3)</td>
<td>345.340</td>
<td>LSB</td>
<td>234 MHz</td>
<td>2 3</td>
<td>1/4-corr</td>
</tr>
<tr>
<td>-</td>
<td>HCN (4-3)</td>
<td>354.505</td>
<td>USB</td>
<td>469 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>HCO$^+$ (4-3)</td>
<td>356.734</td>
<td>USB</td>
<td>117 MHz</td>
<td>4 1</td>
<td>1/4-corr</td>
</tr>
<tr>
<td>-</td>
<td>NH$_2$D (2(2,0)-2(1,2))</td>
<td>356.230</td>
<td>USB</td>
<td>234 MHz</td>
<td>4 2</td>
<td>1/2-corr</td>
</tr>
<tr>
<td>7 c</td>
<td>$^{13}$CO (3-2)</td>
<td>345.796</td>
<td>USB</td>
<td>58 MHz</td>
<td>1 1</td>
<td>J=3-2 CO/$^{13}$CO</td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>344.8</td>
<td>USB</td>
<td>TDM</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>$^{13}$CO (3-2)</td>
<td>330.588</td>
<td>LSB</td>
<td>58 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>331.6</td>
<td>LSB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>8 a Standard cont.</td>
<td>492.160</td>
<td>USB</td>
<td>117 MHz</td>
<td>1 1</td>
<td>CI/$^{13}$CI</td>
<td></td>
</tr>
<tr>
<td>8 b</td>
<td>CI ($^3$P$_1$ - $^3$P$_0$)</td>
<td>492.160</td>
<td>USB</td>
<td>117 MHz</td>
<td>1 1</td>
<td>CI/$^{13}$CI</td>
</tr>
<tr>
<td>-</td>
<td>CS (10-9)</td>
<td>489.751</td>
<td>USB</td>
<td>117 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>477.5</td>
<td>LSB</td>
<td>TDM</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>479.5</td>
<td>LSB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>9 a Standard cont.</td>
<td>679.0</td>
<td>USB</td>
<td>TDM</td>
<td>TDM</td>
<td>LO1=671.0</td>
<td></td>
</tr>
<tr>
<td>9 b</td>
<td>$^{13}$CO (6-5)</td>
<td>691.472</td>
<td>USB</td>
<td>469 MHz</td>
<td>1 1</td>
<td>CO/CS</td>
</tr>
<tr>
<td>-</td>
<td>CS (14-13)</td>
<td>685.436</td>
<td>USB</td>
<td>117 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>H$_2$S (2(0,2)-1(1,1))</td>
<td>687.303</td>
<td>USB</td>
<td>234 MHz</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>C$^{18}$O (6-5)</td>
<td>674.009</td>
<td>LSB</td>
<td>469 MHz</td>
<td>4 1</td>
<td></td>
</tr>
<tr>
<td>10 a Standard cont.</td>
<td>806.652</td>
<td>USB</td>
<td>469 MHz</td>
<td>1 1</td>
<td>CO/HCO$^+$</td>
<td></td>
</tr>
<tr>
<td>10 b</td>
<td>CO (7-6)</td>
<td>806.652</td>
<td>USB</td>
<td>469 MHz</td>
<td>1 1</td>
<td>CO/HCO$^+$</td>
</tr>
<tr>
<td>-</td>
<td>HCO$^+$ (9-8)</td>
<td>802.458</td>
<td>USB</td>
<td>117 MHz</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>805.5</td>
<td>USB</td>
<td>TDM</td>
<td>3 1</td>
<td></td>
</tr>
<tr>
<td>-</td>
<td>continuum spws</td>
<td>803.65</td>
<td>USB</td>
<td>TDM</td>
<td>4 1</td>
<td></td>
</tr>
</tbody>
</table>
For 2SB receivers, typically this would have BB2 in the same sideband (offset by 2 GHz from \(f\)) and BB3 and BB4 in the opposite sideband (offset from BB1 and BB2 by the front end center IF frequency, \(f_{IF}\)). To maximize sensitivity the continuum BBs should cover the maximum aggregate unique bandwidth, i.e. they should not overlap in frequency. If \(sb\) is the sign of the sideband of BB1 (\(sb = +1\) for USB, \(sb = -1\) for LSB), then the four BBs could be set up as follows:

- **BB1**: \(f\) (BB covering the primary line)
- **BB2**: \(f - 2.0\) (continuum BB in the same sideband, 2 GHz below the primary line)
- **BB3**: \(f - (2.sbf_{IF})\) (continuum BB in the opposite sideband)
- **BB4**: \(f - (2.sbf_{IF}) - 2.0\) (continuum BB in the opposite sideband)

where \(f_{IF}\) is the front-end IF frequency (6.0 GHz for Bands 3, 4, 5, 7 and 8, and normally 8.0 GHz for Bands 6, 9 and 10). For DSB receivers, it is common to keep all the BBs in the same sideband, so BB3 and BB4 would be at \(f - 4.0\) GHz and \(f - 6.0\) GHz. Note that this is an approximate rule - if possible, adjustment of the continuum BBs (for example, choosing whether the opposite sideband is LSB or USB) should be done using the OT spectral display, to avoid deep atmospheric absorption features. This is particularly important at Bands 8, 9 and 10, and near the water lines around 183 GHz in Band 5 and 325 GHz in Band 7 (see Figure 6.2). Also, if contiguous spectral coverage of the continuum is desired, the BBs should be offset by 1.875 GHz rather than 2.0 GHz as suggested above. Adding TDM spws in otherwise unused BBs (at non-overlapping frequencies) will improve the overall sensitivity for calibration and it is recommended to use all four BBs whenever possible.

### 6.3.2 Multiple Spectral Windows in the Same Baseband

The system allows the capability to have up to four spws in each baseband, and each baseband can have completely different setups. This is useful for projects where several lines need to be observed simultaneously at high spectral resolution. However, in cases where four lines or fewer are observed in total, if the full bandwidth/resolution is needed for each line and they are well-separated, it may be preferable to place each in a different baseband. Alternatively, in cases where the lines are close to one another, or less spectral resolution/bandwidth is acceptable, it may be better to include multiple lines in the same baseband or even in the same spws, and use the other basebands in TDM mode to maximize the total (aggregate) bandwidth for phase calibration. An example of this is shown in Table 6.1 for spectral setup 6 b, where both C\(^{18}\)O and C\(^{13}\)CO are observed in the same baseband (3.1 and 3.2). Channel-averaging can be applied to each spws, which simply bins spectral channels together in the correlator data processing (see Section 6.5).

Multiple spws per baseband does have some restrictions: each spws within the same baseband must have the *same* correlator channel width (i.e. before channel-averaging), and the sum of the correlator resources in each baseband should add up to 1.0 (or less). Given these restrictions, it is possible to set up different bandwidths and/or different channel-average values for each spws. Multiple spws are set up in the Phase 1 OT simply by adding more spectral lines in the same baseband, and setting the spws to have a correlator fraction less than 1.0 (so for 4 spws, each has a correlator fraction of 1/4). There may be up to 4 spws per baseband, however, the total number of spectral channels in each baseband is limited by the correlator resources. So doubling the number of spws in one baseband will result in half the number of channels per spws, i.e. a lower resolution.

Figure 6.3, which is adapted from the Phase 2 Spectral Editor of the Observing Tool (OT), illustrates a more complex spectral setup with multiple spws and basebands. The lines observed are given in Table 6.1, example 7 b. In this case, the frequency of LO1 is 350.0 GHz, and the upper and lower sidebands of the Band 7 receiver are shown as green shaded areas. The four basebands, illustrated in this case by the red horizontal bars, can be moved around, but only within the two sidebands. The spectral windows are also shown, labelled by the name of the primary spectral line or continuum. In this example, baseband 1 is set for a TDM spws covering the whole 2

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3Note - this case assumes one spws per BB

4although, as before, spws in different basebands can have different channel widths
6.3. **SPECTRAL SETUPS**

GHz baseband for a continuum measurement, whereas in other basebands, FDM is used. Baseband 2 has three spws with 0.56 MHz resolution and channel-average of 2; two of 234 MHz bandwidth and 1/4 of the correlator resources per spws, and one with 469 MHz and 1/2 of the correlator resources. At the same time, in BB3 one line is observed with 469 MHz bandwidth and 0.49 MHz resolution (with a factor of 4 channel-averaging). In BB4, one spws has 234 MHz bandwidth and 0.28 MHz resolution (with 2 channel-average) and one spws has 117 MHz bandwidth and 0.48 MHz resolution (with 4 channel-average). Basebands (and spws) are allowed to partially overlap, although the overlapping region only counts once in the aggregate bandwidth and estimate of continuum sensitivity. So it may be preferable to observe closely-spaced lines with a single baseband (e.g. C$^{18}$O and $^{13}$CO in example 6 b above) or better, with a single spws, as this can free an additional TDM spws for improved continuum sensitivity. Many combinations of spws and basebands can be set up in this way.

Figure 6.3: Illustration of a frequency setup, based on the OT spectral display. Green areas are the IF ranges, horizontal red bars are the 2 GHz-wide basebands (BB1–BB4), and smaller horizontal lines represent the spectral windows (up to 3 per baseband in this example). The frequency of LO1 is shown by the central vertical line. The blue hashed area represents the tuning range of the front end (in this case a section of Band 7), and the curved line represents the nominal atmospheric transmission for the chosen PWV. The bandwidths of the spws are illustrated by the widths of the horizontal lines. The lines observed are given in example 7 b in Table 6.1.

### 6.3.3 Spectral Setups for Lines near the Edge of the Bands

The restriction that the baseband coverage cannot fall outside the maximum or minimum tuning range of the receiver is an issue for certain lines at the edge of the tuning range. One obvious example is the $^{12}$CO($J=1-0$) line at a redshift of zero (rest frequency of 115.271 GHz). This is close to the maximum tuning range of Band 3 (116.000 GHz, see Table 6.1). A setup using the wide bandwidth modes (TDM or FDM/1875 MHz) centered at the line frequency with zero redshift will not validate in the OT, because some of the basebands will fall outside the maximum Band 3 frequency (116.0 GHz). Narrower modes will not go over the edge and are allowed.

There are two possible solutions: If the full bandwidth is absolutely required, the center rest frequency can be set to the closest valid frequency, resulting in an offset of the line from the spws center (in this case, set to 115.0 GHz and the line will be offset by 0.271 GHz). Another solution is to choose a narrow spws bandwidth (e.g. 1 GHz or less); in this case the spws will be offset from the center of the baseband, and the line will be at the center of the spws.
### 6.3.4 Spectral Setup in Bands 9 and 10: DSB Considerations and 90-degree Phase Switching

Bands 9 and 10 use double-sideband (DSB) receivers. If only one sideband is required (the so-called 'signal' sideband) the other sideband (known as the 'image' sideband) can either be suppressed using LO offsetting (currently using LO1 and LO2). Alternatively, both sidebands can simultaneously be recorded using 90-degree phase switching, although this can only be used in interferometric mode - see Section B.4.2.

#### Observing a single sideband

The choice of which sideband a particular baseband is configured to observe (and which sideband is suppressed) depends only on the relative sign of the LO1 and LO2 offset. This can be arbitrarily different for different basebands. So it is possible to set up two basebands at approximately the same LO1/LO2 frequency, but observe one line in the upper sideband in one baseband, and another in the lower sideband using the next baseband, just by having different signs in the LO2 offsetting. However, in this case it might be better to use 90-degree switching (see next section) to record both using the same baseband. Note that different spws in the same baseband must have the same sideband, as image suppression uses LO2, which is common to the whole baseband.

#### Observing both sidebands - 90-degree switching

In Cycle 6 it is possible to double the continuum bandwidth by simultaneously recording both sidebands from the DSB receivers using 90-degree phase switching. This is enabled by phase modulation of LO1, with ±90 degree demodulation in the correlator, controlled with a Walsh sequence (see Chapter 5; also ALMA Memo 287). It has the advantage of doubling the continuum bandwidth to a nominal 16 GHz per polarisation (2 x 2 GHz per baseband) - particularly helpful in these high-frequency bands as it can improve the continuum sensitivity by \( \sim \sqrt{2} \), as well as making available the lines in the opposite sideband. The actual improvement however, depends on the atmospheric transmission in the opposite sideband, which in Bands 9-10 can result in image system temperatures 20% (or more) higher or lower than the main signal sideband. Even though this doubles the number of spectral windows there is no loss of correlator spectral resolution, so the correlator data rate is doubled. This may cause correlator data rate limitations in FDM mode, which could be mitigated by longer integration times, lowering resolution by increasing the channel-averaging factor, and/or using TDM in the continuum-only basebands. For Cycle 6, there is a restriction that this mode can only be used with TDM or the maximum-bandwidth (1875 MHz) FDM configuration, and the spws can only be centered on each baseband. It also requires that the integration duration should be a multiple of 2048 ms.

### 6.4 Usable Bandwidth

The IF system contains an anti-aliasing filter which limits the bandwidth of the basebands. Nominally this filter has -1 dB points at 2.10 and 3.90 GHz, giving a maximum bandwidth of 1.8 GHz. However, the IF response is such that the usable bandwidth is slightly wider – i.e. closer to 1.9 GHz. In FDM mode, the correlator outputs a maximum bandwidth of 15/16 of the nominal bandwidth, and reduces the number of channels by the same factors, from 4096 to 3840 (or 2048 to 1920 after the default channel-average factor - see Chapter 5). So for 2 GHz nominal, the correlator outputs a bandwidth of 1.875 GHz. Thus in FDM wideband, the filters do not truncate the spectrum, and the full available correlator bandwidth in wideband mode can be used. In TDM the correlator outputs a bandwidth of 2.000 GHz, but typically the edges of the spectra are affected by low power due to this filter and some ringing effects (see upper panel in Figure 6.4). It is recommended that 4 (in double-polarization) or 8 (single-polarization) channels are removed or flagged manually offline. This results in approximately the same usable bandwidth in both TDM and FDM modes and is illustrated in Figure 6.4. Note that if the centers of two basebands are separated by 1.875 GHz it is possible to set up spws that can be joined in data reduction to yield a contiguous 3.75 GHz spectrum.
6.5 Spectral Resolution and Channel-averaging

The spectral resolution of each spws is set by a combination of the inherent resolution/bandwidth setup in the correlator, the correlator weighting function (also known as the smoothing function) and the channel-averaging (also known as spectral-averaging) factor (which is carried out at a late stage of processing in the correlator). This is described in the correlator Section 5.5.2. The default FDM setup from the OT in Cycle 6 uses Hanning smoothing with a channel-averaging factor of 2 (see Section 5.5.2). This is changed from previous Cycles, which had a channel-average factor of 1, and has been introduced to reduce the size of the FDM datasets with negligible loss of spectral resolution. Hanning smoothing plus a channel-average of 2 results in a spectral resolution equal to the final channel spacing x 1.15 - only ~15% worse than Hanning without channel-averaging but has the advantage of halving the data rate. If the line is well resolved, with ≥3 channels covering the linewidth, then this will not cause a problem. However, marginally resolved lines should keep the maximum resolution. If possible it is recommended to use this default in order to reduce the data rate and size of the final product; this is more critical where fast dump times are needed, e.g. for long baseline observing. If the extra 15% resolution is required or the line is expected to be marginally resolved, a justification for using channel-average=1 should be made. Note that for TDM mode, Hanning smoothing without channel-averaging (1) is the default.

6.6 Spectral scans

For coverage of frequency ranges greater than the IF bandwidth, a spectral scan observing mode is available. If spectral scan is selected in the OT, this will observe up to 5 different LO1 frequency tunings to cover the requested range, using all four BBs in USB and LSB. Tunings will be programmed to include an overlap between each setting to avoid issues with edge channels. Band 9 and 10 will also use the 90-degree phase switching to double to effective bandwidth per tuning. If more than 5 tunings are required to cover the range, additional science goals need to be created.
6.7 Spurious Signals

Most spurious signals in the cross-correlation data are suppressed using a combination of 180-degree phase modulation of LO1 and using different LO offsets for different antennas, both indexed with a Walsh pattern (see Section B.4.4). Not only does this effectively suppress signals generated after the front ends, but it also improves the image sideband rejection for 2SB receivers by an additional factor of ~20 dB. However, it does not reduce spurious signals coming in at the observing frequency. Results from tests using these Walsh functions show very few remaining spurious signals, and these are further reduced by the fringe tracking. However, harmonics of the LO in the WVRs (ie 91.66, 183.32, 274.98, 366.64, 458.3 ... GHz) are very bright and cannot be removed through these methods. These are very narrow and may require a few spectral channels to be flagged out during data reduction. Observing lines at these frequencies should be avoided if possible.

6.8 Doppler Setting and Velocity Reference Frames

In most cases the system will be set up to provide on-line correction for the science target velocity in a particular reference frame and the Earth motion in that frame. The primary velocity reference frames recommended for use in ALMA are:

Topocentric In this case no correction for the source or Earth motion is made. The Center Frequency (both Rest and Sky) will be identical.

Barycentric This is with respect to the center of mass of the Earth-Sun system, and is very close to the heliocentric frame.

LSRK Velocity with respect to the Kinematic local standard of rest, at 20.0 km/s in the direction 18h, +30° [B1900.0]. This frame is based on the mean velocity of the stars in the Solar neighbourhood.

If a target velocity and reference frame other than topocentric is selected in the OT, at the start of the observation, the velocity of the science target in the chosen reference frame and the velocity of the observatory relative to the chosen reference frame are used to set the center frequencies of the science spws. For example, for a source of radial velocity 100 km/s in the LSRK frame, the difference between rest and sky frequency in the OT will be the equivalent of 100 km/s. At runtime, the extra velocity from the Earth orbit and rotation (up to ±30.3 km/s and ±0.46 km/s), plus the motion of the LSRK (20.0 km/s toward RA=18 hrs, δ =30° for epoch 1900.0) will be included. This means that the LO solutions are slightly different at runtime compared with those from the OT, although this is normally transparent to the user. It is also possible to just make the correction for the motion of the observatory relative to the chosen reference frame, ignoring the source velocity in that frame. The frequency setting for the science target will normally also be used in the same execution for both the bandpass and amplitude calibration.

For sources with an external ephemeris file, the rate of change of distance between the target and the observatory taken from the ephemeris is used to compute the source velocity at the start of the execution. This velocity is used throughout the observation (like non-ephemeris observations, the velocity from the ephemeris is not updated). Combining executions with different velocities must be done in CASA offline.

The user can specify either the velocity or the redshift of the target. Note that velocities can be defined in three different ways, resulting in different conversions to frequency. These differences only become significant at high velocities.

---

5 Although the WVR LOs have small offsets which are different for each antenna, which helps to reduce their cross-correlated signal

6 Note that ALMA does not do Doppler tracking, where the frequency would be continuously updated for Earth motion during the observation. Doppler corrections are only set once at the start of each execution. Compensation for the small changes during an observation (<0.1 km/s for an execution lasting 2 hours) are made during CASA offline reduction.

7 The system allows for a small (< 1 MHz at Band 6) overhang of the spectral windows outside the 2 GHz-wide basebands at runtime, to allow for the very slightly different Doppler corrections in different spws, which will shift the spws with respect to one another.
6.8. DOPPLER SETTING AND VELOCITY REFERENCE FRAMES

The velocity formulae are:

Radio \( v = c(f_o - f)/f_o \). This is the default.

Optical \( v = c(f_o - f)/f \)

Relativistic \( v = c(f_o^2 - f^2)/(f_o^2 + f^2) \).

In these equations, \( v \) is the source velocity, and \( f \) and \( f_o \) are the observed and the rest frequencies of the line. If a redshift \( z \) is specified, the conversion to frequency uses \( z = (f_o - f)/f \). For further details, see the Knowledgebase article “What are the frequency reference frames in CASA?”

<table>
<thead>
<tr>
<th>Rule</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.</td>
<td>LO1 must lie within the LO tuning ranges given in Table 6.1.</td>
</tr>
<tr>
<td>2.</td>
<td>No part of the 2.0 GHz-wide basebands can extend over the edge of the IF passband. So, the baseband centers cannot be closer than 1.0 GHz to the IF passband edge. For example, with a 4.0–8.0 GHz IF range, the baseband center frequency must lie between 5.0–7.0 GHz. The system actually does allow a small extension of the edges of the basebands over the IF edges, to cope with the differential Doppler shifts in the different basebands, but this is small (&lt; 1 MHz at Band 6) and is transparent to the OT user (see Section 6.8).</td>
</tr>
<tr>
<td>3.</td>
<td>For 2SB receivers (Bands 3–8), the number of basebands in one sideband can only be 0, 1, 2, 3, and 4. However, note that three basebands in one sideband and one baseband in the other sideband is not allowed whereas the combination three-zero is. For DSB receivers (Bands 9 and 10), there is no such restriction (the number can be 0, 1, 2, 3 or 4).</td>
</tr>
<tr>
<td>4.</td>
<td>For &gt; 2 basebands, BB1 and BB2 should be in the same sideband, as should BB3 and BB4.</td>
</tr>
<tr>
<td>5.</td>
<td>No part of the full nominal bandwidth of the spws can extend over the edge of the 2 GHz-wide baseband. For a mode with nominal bandwidth B (e.g. 62.5 MHz), that means the spws center IF frequency (a.k.a. Center Offset Frequency in the OT, Phase 2) must be &gt;(2000+B/2) and &lt;(4000-B/2) MHz. The current version of the OT forces this restriction. For 2 GHz FDM and TDM modes, this means that the spws must be at the center of the baseband. However there is a further restriction on this, as noted in the next rule.</td>
</tr>
<tr>
<td>6.</td>
<td>The spws usable bandwidth (i.e. 15/16 of the nominal spws bandwidth of a multiple of 62.5 MHz) should be in an allowed region of the baseband. This is in addition to (5). In practice this means that the required range of the spws should normally be inside the range 2050 - 3950 MHz (i.e. &gt;50 MHz from the edges of the 2–4 GHz second IF) to ensure that the edge of the anti-aliasing filter does not significantly affect the IF power. In practice it is possible to extend some fraction of the spws to &lt;50 MHz of the IF edges (although this is no longer so critical for the Band 6 setup 6b in Table 6.1, as the allowable IF range has been extended down to 4.5 GHz in Cycle 6).</td>
</tr>
<tr>
<td>7.</td>
<td>The line frequency does not need to be in the center of the spws in Phase 1 of the OT; if a line like 12CO 1-0 is requested, it will generate an SB with the correct TFB LO offset, as long as rule 5 is obeyed (which may require using a narrower spws). This is mentioned in more detail in Section 6.2</td>
</tr>
<tr>
<td>8.</td>
<td>Only 2-bit, Nyquist sampling is allowed in the correlator.</td>
</tr>
<tr>
<td>9.</td>
<td>It is possible to have multiple targets with different redshifts within the same Science Goal in the OT. For SGs including sources with more than one redshift, all the observations must be achievable using five or fewer tunings within the same receiver band, considering the source redshifts and, in the case of spectral lines, the line widths and configuration of spectral windows.</td>
</tr>
<tr>
<td>10.</td>
<td>The number of spws per baseband can be 1, 2, 3 or 4. For 3 spws, the correlator resources per spws should be set to 1/4, so only 3/4 of the available resources are used in this case.</td>
</tr>
<tr>
<td>11.</td>
<td>All the spws within the same baseband are required to have the same correlator channel width (before channel-averaging). An individual spws within a baseband may occupy 1, 1/2 or 1/4 of the resources available in the baseband and the sum of the fractional resources within one baseband must be ( \leq 1 ). The correlator resources are proportional to the number of correlator spectral channels.</td>
</tr>
<tr>
<td>12.</td>
<td>With DSB receivers (Band 9 and 10), 90-degree phase switching is allowed, giving both sidebands simultaneously. This is only available for TDM and FDM 1875 MHz bandwidth modes.</td>
</tr>
</tbody>
</table>

Table 6.2: Rules for spectral setups.
6.9 Limitations and Rules for Spectral Setups in Cycle 6

The correlators already allow for a broad flexibility of spws in a single observation, however, the full capabilities are gradually being introduced and tested by the observatory before release. The current rules for spectral setups are given in Table 6.2.
Chapter 7

Imaging with ALMA

7.1 Introduction

Imaging ALMA data consists of Fourier-transforming calibrated interferometric visibilities from one or more configurations or arrays; de-convolving the interferometer point spread function (PSF); and, if Single-Dish data have been collected, merging them with the interferometric image or cube. We describe in this chapter some of the key considerations of this process that should be taken into account in order to obtain the best results from ALMA observations.

As described in Chapter 3, the van Cittert-Zernike theorem is the mathematical foundation of synthesis imaging. It describes a fundamental relationship between the sky brightness distribution ($I$), the beam pattern ($A$) and the visibility distribution ($V$):

\[ A(l, m)I(l, m) = \int \int V(u, v)e^{2\pi i(u\lambda + v\gamma)}dudv \]  \hfill (7.1)

A more detailed discussion of the van Cittert-Zernike theorem can be found, e.g., in Rohlfs & Wilson (2004). An interferometric observation can be represented by $N$ discrete points in the $(u, v)$ plane $V_k(u_k, v_k)$ where $k = 1, \ldots, N$. The variables $u_k$ and $v_k$ represent the $x$ and $y$ components of a a specific $k^{th}$ baseline between a pair of antennas, measured in wavelengths $\lambda$. There are in general $N_{\text{ant}}(N_{\text{ant}} - 1) \times 0.5$ distinct baselines; factoring in distinct frequency channels and temporal integrations, $N$ can reach a few million or more for ALMA. One implication of the van Cittert-Zernike theorem, which will be discussed in more detail elsewhere in this chapter, is that the interferometer point-spread function (PSF) is the Fourier Transform of the set of $(u_k, v_k)$ points measured. In the parlance of radio astronomy this set of $(u, v)$ points is often called the “$(u, v)$ coverage” of an observation, and the PSF is called the “synthesized” or “dirty” beam. In practice it is often useful to define a weighted visibility function $V^W$ (see Briggs, Schwab & Sramek, 1999) that enables the properties of the synthesized beam to be controlled to some extent, often at the expense of sensitivity. The weighted visibility function can be represented as:

\[ V^W(u, v) = \sum_{k=1}^{N} R_k T_k D_k \delta(u - u_k, v - v_k)V_k(u_k, v_k). \]  \hfill (7.2)

In Equation 7.2, $R_k$, $T_k$ and $D_k$ are weights which control the relative weighting of individual visibilities in imaging. $R_k$ is a noise-variance weight derived from the sensitivity of the $k^{th}$ visibility measurement. It accounts for variables such as the integration time, the system temperature, and the bandwidth. These factors are determined by the instrument, technical set-up, and observing conditions, and are not controlled by the imaging process. $T_k$ is a “taper” weight— usually taken to be a multiplicative, Gaussian factor in the $(u, v)$ plane— that can be used to down-weight the longest baselines, suppress small-scale sidelobes, and increase the synthesized beam width and surface brightness sensitivity. Finally, $D_k$ is a $(u, v)$ “density” weight that can be
used to offset the high concentration of measurements near the center of the \((u, v)\) plane, typically increasing angular resolution and reducing the large-scale sidelobes caused by abrupt gaps in the \((u, v)\) coverage density.

Figure 7.1: This Science Verification image is one of the sharpest images ever taken by ALMA (ALMA Partnership 2015). It shows the protoplanetary disc surrounding the young star HL Tauri. These observations reveal substructures within the disc that have never been seen before and even show the possible positions of planets forming in the dark patches within the system. The data, calibration and imaging scripts used to create this image are available from http://almascience.org/alma-data/science-verification and http://casaguides.nrao.edu/index.php?title=ALMA2014_LBC_SVDATA. Credit: ALMA (ESO/NAOJ/NRAO)

The following sections describe: the ALMA configurations (Section 7.2); the effects of antenna shadowing (Section 7.3); the factors that effect angular resolution and beam shape (Section 7.4); the special case of short “snapshot” observations (Section 7.5); the large spatial scale response of ALMA (Section 7.6); mosaicing (Section 7.7); multi-array and multi-configuration ALMA observations (Section 7.8); and, finally, multi-array and multi-configuration imaging (Section 7.9).

### 7.2 Cycle 6 Configurations

In Cycle 6, depending on the range of angular scales required, an ALMA Science Goal will obtain data from either: a single 12-m Array configuration; two 12-m Array configurations; one or two 12-m Array configurations plus the 7-m Array; or one or two 12-m Array configurations plus the 7-m and Total Power (TP) Arrays\(^1\). The detailed logic by which specific combinations are chosen based on the PI Science Goals is described in Section 7.8. Generally, the PI’s desired angular resolution determines the most extended array, while the largest angular scale of interest determines the most compact array—whether that is a second 12-m Array, the 7-m Array, or the Total Power in order of progressively larger spatial scales. Intermediate Configurations or Arrays are added to connect these spatial scales as required.

During Cycle 6 the 12-m Array will be arranged in ten different configurations. As discussed below, operational factors impact the actual configuration achieved in a given observation, so these ten configurations are called “representative configurations”. The representative configurations for Cycle 6 are identical to those used in Cycle 5. Each 12-m Array configuration will have at least 43 antennas. The maximum baseline will

\(^1\)The combination of the 7-m Array and the Total Power Array constitutes the ALMA Compact Array (ACA). The ACA is also known as the Morita Array in honor of Prof. Koh-ichiro Morita.
be 16.2 km for Bands 3-6, 8.5 km for Band 7 and 3.6 km for Bands 8-10. The 7-m Array will be available in only one configuration, with at least ten antennas. Four extended configurations will be available in Cycle 6: C43-7, C43-8, C43-9, and C43-10. The most extended configuration - C43-10 - will be offered alone; it cannot be combined with other configurations. All other configurations are offered in combination with other configurations.

Tables 7.1 and 7.2 give the basic properties of the 12-m Array and 7-m Array configurations. Each configuration is characterized by a maximum recoverable (spatial) scale $\theta_{MRS}$ and an angular resolution $\theta_{res}$. Generally these are taken to be determined by the shortest and longest baselines in the configuration, respectively. As discussed in Section 7.4 and Section 7.6, the ALMA project uses the somewhat more robust 5th and 80th percentile baseline lengths instead. Figures 7.2, 7.3, and 7.4 show antenna locations for the configurations of the 12-m and 7-m Arrays. The three TP Array single-dish 12-m antennas will be in fixed positions, although their relative spatial locations do not substantially affect their imaging characteristics.

<table>
<thead>
<tr>
<th>Configuration</th>
<th>Frequency (GHz)</th>
<th>Band</th>
<th>3</th>
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<th>6</th>
<th>7</th>
<th>8</th>
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<tr>
<td>$\theta_{res}$ (arcsec)</td>
<td>12.5</td>
<td>8.35</td>
<td>6.77</td>
<td>5.45</td>
<td>3.63</td>
<td>2.72</td>
<td>1.93</td>
<td>1.44</td>
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<td>$\theta_{MRS}$ (arcsec)</td>
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<tr>
<td>$\theta_{res}$ (arcsec)</td>
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<td>$\theta_{res}$ (arcsec)</td>
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<td>1.53</td>
<td>1.24</td>
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<td>4.68</td>
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<td>0.2</td>
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<td>0.295</td>
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<td>1.94</td>
<td>1.46</td>
<td>1.03</td>
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<td>$\theta_{MRS}$ (arcsec)</td>
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<td>2.22</td>
<td>1.78</td>
<td>1.19</td>
<td>0.892</td>
<td>0.632</td>
<td>0.472</td>
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<td>C43-7</td>
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<tr>
<td>$\theta_{res}$ (arcsec)</td>
<td>0.231</td>
<td>0.141</td>
<td>0.114</td>
<td>0.0917</td>
<td>0.0612</td>
<td>0.0459</td>
<td>0.0325</td>
<td>0.0243</td>
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</tr>
<tr>
<td>$\theta_{MRS}$ (arcsec)</td>
<td>2.58</td>
<td>1.72</td>
<td>1.4</td>
<td>1.12</td>
<td>0.749</td>
<td>0.562</td>
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<tr>
<td>$\theta_{res}$ (arcsec)</td>
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<td>0.064</td>
<td>0.0519</td>
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<tr>
<td>$\theta_{MRS}$ (arcsec)</td>
<td>1.42</td>
<td>0.947</td>
<td>0.768</td>
<td>0.618</td>
<td>0.412</td>
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<td>0.038</td>
<td>0.0308</td>
<td>0.0248</td>
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<td>$\theta_{MRS}$ (arcsec)</td>
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<td>0.543</td>
<td>0.44</td>
<td>0.354</td>
<td>-</td>
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<tr>
<td>C43-10</td>
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<tr>
<td>$\theta_{res}$ (arcsec)</td>
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<td>0.028</td>
<td>0.0227</td>
<td>0.0183</td>
<td>-</td>
<td>-</td>
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</tr>
<tr>
<td>$\theta_{MRS}$ (arcsec)</td>
<td>0.496</td>
<td>0.331</td>
<td>0.268</td>
<td>0.216</td>
<td>-</td>
<td>-</td>
<td>-</td>
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</tr>
</tbody>
</table>

Table 7.1: Resolution ($\theta_{res}$) and maximum recoverable scale ($\theta_{MRS}$) for the 7-m Array and 12-m Array configurations available during Cycle 6 as a function of a representative frequency in a band. The value of $\theta_{MRS}$ is computed using the 5th percentile baseline (L50) from Table 7.2 and Equation 7.7. The value of $\theta_{res}$ is the mean size of the interferometric beam obtained through simulation with CASA, using Briggs $(u,v)$ plane weighting with $\text{robust}=0.5$. The computations were done for a source at zenith; for sources transiting at lower elevations, the North-South angular measures will increase proportional to $1/\sin$(ELEVATION).

12-m Array configurations in particular are impacted by a variety of “real-world” factors including, but not limited to, the availability of specific antenna pads\(^2\), and necessary, ongoing antenna relocations. To help accommodate these factors, “hybrid” configurations are often used. These operational considerations will often change the resolution and maximum angular scale that are achieved by executing a given scheduling block. The scheduling subsystem takes these changes into account alongside the performance margins that the ALMA QA2 process allows. Thus from the user’s point of view configurations are not directly selected: instead the

\(^2\)ALMA’s 54 12-m antennas can be distributed among 192 “antenna pads” which provide power, IF signal & network connections, and a structurally stable foundation.
CHAPTER 7. IMAGING WITH ALMA

<table>
<thead>
<tr>
<th>Configuration</th>
<th>7-m</th>
<th>C43-1</th>
<th>C43-2</th>
<th>C43-3</th>
<th>C43-4</th>
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<td>5th percentile or L05 (m)</td>
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<td>80th percentile or L80 (m)</td>
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<td>369.2</td>
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<td>500.2</td>
<td>783.5</td>
<td>1397.9</td>
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</table>

<table>
<thead>
<tr>
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<th>C43-7</th>
<th>C43-8</th>
<th>C43-9</th>
<th>C43-10</th>
</tr>
</thead>
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<td>Minimum baseline (m)</td>
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<td>244.0</td>
</tr>
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<td>5th percentile or L05 (m)</td>
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<td>235.2</td>
<td>427.3</td>
<td>746.9</td>
<td>1228.1</td>
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<td>3527.3</td>
<td>6482.6</td>
<td>8685.9</td>
</tr>
<tr>
<td>Maximum baseline (m)</td>
<td>2516.9</td>
<td>3637.8</td>
<td>8547.7</td>
<td>13894.2</td>
<td>16194.0</td>
</tr>
</tbody>
</table>

Table 7.2: Basic parameters of the 7-m Array configuration and the ten 12-m Array configurations offered during Cycle 6. The baselines are projected for a transiting source (HA = ±1h) at a declination of −23°. Note that C43-9 and C43-10 will not be available for Bands 7-10, and C43-8 will not be available for Bands 8-10.

required angular resolution (θ_{res}) and largest angular scale (θ_{LAS}) are specified in the OT, and projects are observed using array configurations that will achieve the specified science goals. Since Cycle 5 it has also been possible for PI’s to specify an acceptable range of θ_{res}. More details about array combination are provided in Section 7.8.

### 7.3 Shadowing

During interferometric array observations at low elevations—particularly in compact configurations—one antenna’s view of the source can be partially blocked by another antenna, corrupting data from the first antenna. This phenomenon is known as shadowing. Given ALMA’s location (latitude=−23.02917°), targets as far North as declination +47° can in principle be observed, corresponding to a maximum source elevation at Chajnantor of ⇐20°. However shadowing by adjacent antennas can be a significant problem at such low elevations. Figure 7.5 shows the fraction of data that will be shadowed (i.e., the “shadowing fraction”) when sources of various declinations are observed in the most compact ALMA configurations. As can be seen, the shadowing fraction can be as large as 55% for sources observed with the most compact 12-m Array configuration (C43-1). Image quality and time on source will necessarily be limited for such northern sources. Given the short baselines in the ACA, sources with declinations less than −70° or greater than +25° are subject to significant shadowing. For the 12-m Array, shadowing becomes significant (> 5%) in the most compact configuration for sources with declination lower than −65° or higher than +20°.

Note that independent of shadowing, observations at low elevation often result in elliptical (u, v) coverage and, therefore, elliptical synthesized beams.

### 7.4 Angular Resolution and Beam Shape

The angular resolution θ_{res} of an array can be estimated with the following equation:

\[
θ_{res} = \frac{kλ}{L_{max}} \text{ [radians]} \tag{7.3}
\]

where θ_{res} is the FWHM of the main lobe of the synthesized beam; k is a factor that depends on the (u, v) plane weighting function; λ is the observing wavelength in meters; and L_{max} is the longest baseline in meters. For commonly used weighting schemes, k is typically in the range 0.7 < k < 1.2. Since circumstances can sometimes result in unrepresentatively long baselines\(^3\) for a given configuration, the single longest baseline

\(^3\)For example, there could be one or two baselines much longer than the next longest baselines; these would provide very poor
Figure 7.2: Representative 12-m Array compact configurations for Cycle 6.
$L_{\text{max}}$ is a somewhat fragile indicator of angular resolution. Consequently ALMA uses the 80th percentile of the $(u, v)$ distance as a more robust proxy to angular resolution. Using the ALMA representative configurations, $(u, v)$ coverage and commensurately useless imaging.
7.4. ANGULAR RESOLUTION AND BEAM SHAPE

The following equation was determined:

\[ \theta_{res} \approx \frac{0.574 \lambda}{L_{80}} \, \text{[radians]} \]  

(7.4)

where again \( \lambda \) is the observing wavelength in meters, and \( L_{80} \) is the 80\(^{th} \) percentile of the \((u, v)\) distance in meters (see also Table 7.2).

One important consideration in imaging is the choice of \((u, v)\) plane weighting scheme, in particular the \((u, v)\) density weighting scheme. Two commonly used, limiting forms of density weighting are natural weighting and uniform weighting. These are controlled by the density weight \( D_k \) (see Eq. 7.2):

- \( D_k = 1 \) is called natural weighting, and results in maximum sensitivity but a relatively large synthesized beam due to the typically high density of points in the inner region of the \((u, v)\) plane.

- \( D_k = \frac{1}{N_s(k)} \) is called uniform weighting, where \( N_s(k) \) is the number of visibilities in a region centered on the \( k^{th} \) visibility. It removes the dependence of spatial-scale sensitivity on the density of visibilities (samples) in the \((u, v)\) plane. Uniform weighting increases angular resolution at the expense of sensitivity.

To bridge the extremes of natural and uniform weighting, Briggs (1995) defined a continuous scheme that uses a “robustness” parameter \( R \) (called \( robust \) in CASA). In CASA, uniform weighting is close to \( robust = -2 \) and natural weighting is close to \( robust = 2 \). Figure 7.6 shows the angular resolution achieved for an observation with the C43-4 configuration at 100 GHz using \( robust \) values between \(-2\) and \( 2 \). As can be seen, the angular resolution varies from 0.7\(^{\prime\prime} \) (\( robust = -2 \)) to 1.2\(^{\prime\prime} \) (\( robust = 2 \)). In general, the angular resolutions presented in this chapter were computed using CASA simulations with Briggs weighting and \( robust = 0.5 \).

The synthesized beam shape, which is the Fourier transform of the \((u, v)\) plane sampling during the observation(s), is a function of the source declination. In addition to shadowing, sources that must be observed at low elevations also have shorter projected North-South baselines and thus the shape of the \((u, v)\) plane sampling distribution becomes more elongated and the beam shape more elliptical. For example, Figure 7.7 shows the different beam shapes for sources observed at declinations of \(-70^\circ\) (similar to the SMC and LMC) and \(-30^\circ\). Also, Figure 7.8 shows the \((u, v)\) coverage for these sources, revealing the elongation of the \((u, v)\) plane sampling.
distribution, together with the large fraction of shadowing, for the more southern sources. To mitigate this effect, one can convolve the resulting images to a more circular (albeit larger) beam after deconvolution.

Figure 7.9 shows the minor and major axes of the synthesized beam widths ($\theta_{\text{res}}$) for each array configuration as a function of source declination for a 2-hour observation at 100 GHz. Figure 7.10 shows the geometrical mean of the major and minor axes of $\theta_{\text{res}}$ at the same frequency. The beam width scales with $\lambda$, but bear in mind that not all configurations can be used with the higher-frequency bands, e.g., C43-8 and more extended configurations are not available at Bands 8-10.

Note that in the OT, the angular resolution and sensitivity are now computed consistently, assuming Briggs weighting with $\text{robust} = 0.5$. If the imaging requirements of a project are particularly scientifically critical, then for purposes of proposal preparation it is advisable to carefully evaluate the expected performance using CASA simulations that employ the representative configurations, target declination, and expected $(u, v)$ coverage of the proposed observations. In this case it is also advisable to include realistic noise and to CLEAN the images carefully in order to realistically assess the results that can be expected. For a given real dataset, and/or for planning special imaging observations, it may be desirable to vary the value of $\text{robust}$ and other imaging parameters.

As will be discussed in Section 7.9, combining data from compact configurations with those from more extended configurations primarily has the effect of compensating for negative sidelobes in the synthesized beam and thereby improving the interferometer’s response to larger angular structures on the sky. Such combination, however, does also slightly degrade the resolution. This degradation can be mitigated by moving towards uniform density weighting using “robust”.

---

4Most interferometric imaging in CASA is now done using TCLEAN, which is an improved implementation of the CLEAN algorithm. The CLEAN task in CASA has been deprecated; in a near-future release, CLEAN and TCLEAN will be renamed and eventually the “old” CLEAN will be removed from the code.
7.4. ANGULAR RESOLUTION AND BEAM SHAPE

Figure 7.7: Beam shape for configuration C43-1 with a 2-hour observation of a transiting source at a declination of either $-70^\circ$ (left) or $-30^\circ$ (right).

Figure 7.8: $(u, v)$ plane coverage for configuration C43-1 with a 1-hour observation of a transiting source at a declination of either $-70^\circ$ (left) or $-23^\circ$ (right). For the source with a declination of $-70^\circ$, 25.7% of visibilities are expected to be flagged due to shadowing.
Figure 7.9: Cycle 6 angular resolutions as a function of source declination. Each stripe shows, for a particular configuration, the range of the major and minor axes of the synthesized beams expected from a 2-hour observation at 100 GHz and at transit. 12-m Array configurations are arranged with C43-1 at the top and C43-10 at the bottom.

Figure 7.10: Geometrical mean of the major and minor axes of the synthesized beams as a function of source declination. These correspond to a 2-hour observation at 100 GHz and at transit. 12-m Array configurations are arranged with C43-1 at the top and C43-10 at the bottom. The 7-m Array is also shown.
7.5 Snapshots

The sampling function $S$ of the visibility distribution is defined as:

$$S(u, v) = \sum_{k=1}^{N} \delta(u - u_k, v - v_k) \quad (7.5)$$

If $S$ were a continuous function, e.g., a Gaussian, the synthesized beam would also be a Gaussian, i.e., a central peak with a smoothly decreasing response away from the center. Given the finite number of baselines, however, the sampling function is an ensemble of Dirac functions and its Fourier transform is a central peak surrounded by a complex pattern of sidelobes. This “dirty beam” response is a consequence of the gaps in the $(u, v)$ plane, the sidelobes becoming increasingly prominent as the gaps increase (following the Gibbs phenomenon). Conversely, as there are fewer gaps in the $(u, v)$ plane—i.e., better $(u, v)$ coverage—the PSF sidelobes become less prominent.

Because the baseline lengths and orientations as seen from the source change as a function of time due to the Earth’s rotation, the $(u, v)$ coverage increases and the PSF sidelobes are commensurately reduced in proportion to the duration of the observation. Therefore long integrations on a target can give better image quality. Short integrations can still be valuable if the science target has a relatively compact, simple structure and is not too faint. Such observations are sometimes called snapshots, and they produce obvious sidelobes which can be mitigated by applying a $(u, v)$ taper, i.e., down-weighting of the visibilities on longer baselines, during the imaging process. The $(u, v)$ taper is represented as the set of taper weights $T_k$ in Eq. 7.2.

<table>
<thead>
<tr>
<th>Configuration</th>
<th>7-m</th>
<th>C43-1</th>
<th>C43-2</th>
<th>C43-3</th>
<th>C43-4</th>
<th>C43-5</th>
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<td>5.0%</td>
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<td>6.6%</td>
<td>6.7%</td>
<td>7.3%</td>
<td>8.0%</td>
</tr>
<tr>
<td>uniform</td>
<td>39.0%</td>
<td>7.7%</td>
<td>16.7%</td>
<td>12.4%</td>
<td>11.0%</td>
<td>13.1%</td>
</tr>
<tr>
<td>Configuration</td>
<td>C43-6</td>
<td>C43-7</td>
<td>C43-8</td>
<td>C43-9</td>
<td>C43-10</td>
<td></td>
</tr>
<tr>
<td>natural</td>
<td>4.6%</td>
<td>9.6%</td>
<td>7.8%</td>
<td>11.1%</td>
<td>13.0%</td>
<td></td>
</tr>
<tr>
<td>briggs ($R = 0.5$)</td>
<td>6.2%</td>
<td>5.6%</td>
<td>9.7%</td>
<td>9.6%</td>
<td>11.3%</td>
<td></td>
</tr>
<tr>
<td>uniform</td>
<td>11.3%</td>
<td>11.2%</td>
<td>10.1%</td>
<td>8.5%</td>
<td>12.0%</td>
<td></td>
</tr>
</tbody>
</table>

Table 7.3: Sidelobe levels for a 1-hour observation of an unresolved source at a declination of $-23^\circ$ with the different array configurations. The levels are indicated with three different weighting schemes used for the imaging.

Table 7.3 gives the level of sidelobes for each configuration, for a 1-hour observation of a source at Dec=$-23^\circ$. Figures 7.11 and 7.12 show examples of a $(u, v)$ plane sampling distribution and cleaned image for a snapshot of 1-minute duration and a longer integration of 1-hour duration expected with the configuration C43-1. As shown in Figure 7.11, the $(u, v)$ coverage in 1 minute is quite uniformly sampled, but much less dense than that of the 1-hour integration. In addition, the cleaned images of the snapshot and the longer integration (see Figure 7.12) are similar in terms of angular resolution and the apparent differences are quite small. The main difference between the two (besides sensitivity) is in the dirty beam: the sidelobe level for the snapshot will be much higher than for the 1-hour integration. The snapshot image will then require more careful cleaning in order to avoid introducing spurious sources from the strong sidelobes. Condon et al. (1998) gave a very comprehensive description of how this issue impacted the 20-cm NRAO VLA Sky Survey (NVSS).

A useful practice to disentangle sidelobe effects from real point sources, especially with relatively strong point sources, is to perform CASA simulations using a component list of the strongest sources together with the actual array configuration. This test allows one to estimate the sidelobe fingerprint left by the strong point sources after deconvolution (cleaning), although the residuals will likely be higher in practice due to imperfect calibration. It is worthwhile to use the interactive mode during the deconvolution so that residuals can be monitored.
Figure 7.11: \((u, v)\) plane sampling distributions of a model ALMA observation with 1-minute integration \((left)\) and 1-hour integration \((right)\), using the C43-1 configuration to observe a source at a declination of \(-23^\circ\).

Figure 7.12: Images obtained from a model ALMA observation with 1-minute integration \((left)\) and 1-hour integration \((right)\), using the C43-1 configuration to observe a source at declination of \(-23^\circ\). Black contours at 100, 200, 300, 500, 700 and 900 mJy beam\(^{-1}\) are overlaid.
7.6 Large Spatial Scale Response

As described in Chapter 3, an interferometer measures the Fourier components of the sky brightness distribution in an area of the \((u, v)\) plane defined by the array configuration used to observe the target. There is, in particular, some shortest baseline that a given array and configuration measures, ultimately limited by the physical size of the antennas. This shortest baseline corresponds to some largest angular scale which can be usefully imaged by the interferometer.

This spatial filtering is a serious issue that must be considered carefully for each science case. To illustrate the concept, the annularly averaged amplitudes of the Fourier transforms of three uniform disks with sizes of 5\(^\prime\), 10\(^\prime\) and 20\(^\prime\) are shown in Figure 7.13 for an observation at 100 GHz. The smallest uniform disk is closest to a point source and so it has large amplitudes up to a baseline of 180 m. Meanwhile, the most extended disk has large amplitudes only up to \(\sim 40\) m. Therefore, an array with baselines larger than 40 m will not be sensitive to emission on angular scales larger than \(\sim 20\)\(^\prime\), and will filter out most of the emission from such an extended disk. An important consequence of such filtering is that an interferometer only detects a fraction of the total flux density for sources with emission on size scales larger than its shortest baseline. Indeed, if the source only has structures on size scales larger than the shortest observed baselines, one can “resolve out” the source entirely. In order to ameliorate the effects of spatial filtering on extended sources, the 7-m and TP Arrays are available to supplement the 12-m Array configurations if needed.

The maximum recoverable scale \(\theta_{\text{MRS}}\) is the largest angular structure to which a given array is sensitive. It is in principle determined by the length \(L_{\text{min}}\) of the shortest baseline in the array, which measures spatial scales \(\text{i.e.} \) Fourier modes with a period) of \(\lambda/L_{\text{min}}\). In practice the sensitivity of this measurement to such large structures is not very good, so a smaller value of \(\theta_{\text{MRS}}\) is typically adopted. The exact filtering depends on the details of the large-scale structure as well as the short-baseline \((u, v)\) coverage and is best determined by simulations. ALMA has adopted a criterion of measuring 10\% of the total flux density of a uniform disk, which for well-constructed array configurations yields:

\[
\theta_{\text{MRS}} \approx \frac{0.6 \lambda}{L_{\text{min}}} \quad \text{[radians]}
\] (7.6)

where \(\lambda\) is the observing wavelength in meters, and \(L_{\text{min}}\) is the shortest baseline in meters. For reasons similar to those explained in Section 7.4, the ALMA project uses the 5\(^{\text{th}}\) percentile shortest baseline \(L_5\) instead of the very shortest baseline \(L_{\text{min}}\) to calculate \(\theta_{\text{MRS}}\); using simulations of configurations the following equation was determined:

\[
\theta_{\text{MRS}} \approx \frac{0.983 \lambda}{L_5} \quad \text{[radians]}
\] (7.7)

where \(\lambda\) is the observing wavelength in meters, and \(L_5\) is the 5\(^{\text{th}}\) percentile of \((u, v)\) distance in meters. This is the equation used to measure MRS for purposes of choosing configurations during proposal preparation. Table 7.1 lists the \(\theta_{\text{MRS}}\) for the Cycle 6 array configurations. If the scientifically required \(\theta_{\text{res}}\) and \(\theta_{\text{MRS}}\) cannot be achieved by a single array or configuration, multiple 12-m configurations and/or the 7-m and TP arrays, will be called for, as described in Section 7.8. Note that these expressions assume good \((u, v)\) coverage. This is typically the case for ALMA 12-m Array observations, but may not be for very short 7-m Array observations due to the relatively small number of antennas.

Figure 7.14 shows the annularly averaged visibility amplitudes \(\text{vs.} \) \((u, v)\) distance of an example astronomical source, M51. This example is taken from an H-\(\alpha\) image, but it is used to indicate the spatial structure expected at 100 GHz. The amplitudes of the visibilities can be approximated as a power law of \((u, v)\) distance with a negative index. This distribution indicates that most of the power is located in larger scale structures and that power decreases rapidly at smaller scales. This result shows that the flux that would be received by the 7-m Array would be much higher than that received by the 12-m Array with extended configurations \(\text{e.g.}, \) C43-6). The only case where flux is independent of the sampling in the \((u, v)\) plane is for point sources, which have the same amplitude for all visibilities \(\text{i.e.}, \) the Fourier transform of a Dirac function. In our M51 example, the visibility amplitude detected by the 7-m Array \((k_u \approx 5k\lambda)\) is \(\sim 23\) Jy, whereas that detected by the configuration C43-6 \((k_u \approx 200k\lambda)\) is only \(\sim 0.9\) Jy. Ideally, ALMA users should use simulations to estimate the...
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Figure 7.13: Annularly averaged amplitudes of the expected visibilities of three model uniform disks (annular averages) at 100 GHz as a function of \((u, v)\) distance.

distribution of power at different length scales for targets they wish to observe.

7.7 Mosaicing

The field of view of a single interferometer pointing is determined by the antenna primary beam. A uniformly illuminated, circular aperture will have a beam width at half maximum \((FWHM)\) of \(1.02 \frac{D}{\lambda}\). It will also have very high side lobes due to the abrupt truncation of the antenna illumination pattern; indeed, these side lobes would be particularly problematic for single-dish observations. Most radio and millimeter receivers illuminate their antennas with approximately Gaussian illumination patterns that smoothly go down to -10 dB to -15 dB response at the edges of the dish. The ALMA feedhorns were designed to illuminate the dish with a -12 dB edge taper to provide a nearly Gaussian primary beam with low sidelobes while preserving as much of the resolution and sensitivity as possible. This tradeoff is a fundamental aspect of radio telescope design. The resulting \(FWHM\) of the ALMA 12-m antennas is measured to be \(FWHM = 1.13 \frac{D}{\lambda} = 58''\) at 100 GHz, and it varies inversely with frequency. For example, Figure 7.15 shows a Gaussian profile that approximates the shape of the primary beam at 112 GHz with \(FWHM = 52''\). As described in Chapter 3, the primary beam attenuation can be corrected as the last step of imaging. In addition to correcting the sky distribution of the signal to its correct value, this correction, however, also increases the noise with angular distance from the primary beam center. For example at the radius of the FWHM, the noise will be increased by a factor of two compared to that at the primary beam center. Moreover, beyond a certain map size, the antenna is not sensitive and it is necessary to observe several adjacent pointings, i.e., a mosaic, to recover the sky emission.

Observing a mosaic with ALMA is needed if a map size larger than approximately the \(FWHM\) of the primary beam is required. The default mosaic pointing pattern used by ALMA is a fully sampled hexagonal grid with equilateral triangles whose vertices are separated by \(\theta_{hex} = \frac{\lambda}{D\sqrt{3}} = 0.511 \times FWHM\). With this spacing, the mosaic will sample the emission at the Nyquist spatial frequency. Note that a hexagonal mosaic has spacing \(\theta_{hex}\) along a row (e.g., in right ascension) and \(\frac{\sqrt{3}}{2}\theta_{hex}\) between rows (e.g., in declination). Figure 7.16 gives an example of such a mosaic. To estimate the number of pointings \((N_p)\) necessary to cover an area of \(L_X \times L_Y\) using this hexagonal pattern, the following expressions can be used:
Figure 7.14: Image of H\textalpha\ emission from M51 used as a model of emission at 100 GHz (left) and the expected visibility amplitudes with (u, v) distance (right). Note that M51 cannot be observed from ALMA.

\begin{align*}
N_X &= \left(\text{int}\left(\frac{L_X}{0.311 \text{FWHM}} + 1\right)\right) \\
N_Y &= \left(\text{int}\left(\frac{L_Y}{0.311 \text{FWHM}} + 1\right)\right) \\
N_p &\approx (2N_X - 1) \times \frac{N_Y}{2}
\end{align*}

For an observation at 100 GHz with $L_X = L_Y = 4$ arcmin, 85 pointings are defined, similar to that obtained by the above formulae. In the case of an odd number of rows, $N_p$ may differ slightly from the value returned by the OT. Note that mosaicing with the 7-m Array is more efficient because with the smaller diameter of its antennas, the FWHMs are commensurately larger ($\text{FWHM} = 1.13\frac{D}{2}$ also holds for the 7-m antennas).

Like a single pointing, a mosaic has an analogous “mosaic primary beam response pattern” that is the combination of the individual primary beams of the different pointings. Near the mosaic center, a Nyquist sampled hexagonal mosaic pattern has a sensitivity about 1.58 times that of a single pointing, with the sensitivity decreasing with the fall off of the mosaic primary beam response pattern. Another frequently used mosaic pattern has pointings separated by $\text{FWHM}/\sqrt{2}$ (see for example the NVSS survey); this pattern covers area more efficiently but with little gain in sensitivity over a single pointing. Generally the Nyquist sampled hexagonal pattern (the default on ALMA) is a good choice for smaller mosaics and the constant noise hexagonal pattern is most often used for larger mosaics. The ALMA Observing Tool (OT) can be used easily to set up a mosaic of adjacent pointings with a user-defined spacing, though it is not recommended to exceed spacings greater than the constant noise pattern ($\text{FWHM}/\sqrt{2}$) if a well-sampled mosaic image is desired.

## 7.8 Multi-array and Multi-configuration Observations

As shown in Table 7.1, different 12-m Array configurations provide different angular resolutions, $\theta_{\text{res}}$, which are a function of the longest baselines of the configuration. Similarly, the maximum recoverable scales, $\theta_{\text{MRS}}$, of each configuration depend on the shortest baselines present. To achieve the requirements entered by the PI in the OT, $\theta_{\text{res}}$ and $\theta_{\text{LAS}}$, multiple configurations may be needed. Note that $\theta_{\text{LAS}}$ is a property of the science target, while $\theta_{\text{MRS}}$ is a property of an array configuration. In particular, the most compact configuration for a project must ensure that $\theta_{\text{MRS}}$ is larger than or equal to the $\theta_{\text{LAS}}$ of the science target.

Based on the user-entered values of $\theta_{\text{res}}$ and $\theta_{\text{LAS}}$ for a given Science Goal, the OT will automatically attempt to choose configurations that will produce a final image with the requested scientific goals.\footnote{The entire suite of potentially multi-array/configuration observations that satisfies the Science Goal requirements is the \textit{Group Obs Unit Set} or GOUS. The individual array, individual configuration components it comprises are the \textit{Member Obs Unit Sets} or}
Figure 7.15: Approximation of the primary beam profile of an ALMA 12-m antenna at 112 GHz with FWHM = 52°.

Figure 7.16: An example of mosaicing with a field of 2 arcmin extent at 100 GHz using a hexagonal pattern with Nyquist sampling (white crosses).
good image quality, there are some restrictions on which arrays and configurations can be combined (based mostly on the \((u, v)\) coverage each provides). For example, at most two 12-m Array configurations are allowed and the most extended configurations (C43-10) cannot be combined with any more compact configurations. If such restrictions result in the requested \(\theta_{LAS}\) not being achievable, the OT will return a validation error and \(\theta_{LAS}\) must be modified. Since the middle of Cycle 4, a “stand-alone ACA” option has been offered for observations that use the 7-m Array, with the TP Array added if required, but without 12-m Array observations. It is not possible to request observations that only require the TP array.

In detail, the array configuration combinations are selected as follows:

1. An interferometric array configuration (12-m or 7-m) is selected that comes closest to achieving the entered value of \(\theta_{res}\).
2. If the \(\theta_{MRS}\) of this array configuration is less than the requested \(\theta_{LAS}\), the largest structures in the source will not be well imaged. A more compact 12-m array configuration will be added, if allowed by the set of combinations listed in Table 7.4.
3. If the \(\theta_{MRS}\) of the most compact configuration included thus far is still less than the requested \(\theta_{LAS}\), add a 7-m array component.
4. If \(\theta_{MRS}\) is still < \(\theta_{LAS}\), add a TP array component if possible. Single dish continuum, band 9, and band 10 observations are not currently supported so there is a fundamental limit to how large a source can be reliably imaged in these cases.

This procedure is carried out automatically by the OT and is entirely dependent on the values of \(\theta_{res}\) and \(\theta_{LAS}\) that are entered by the user. A detailed list of the allowed combinations is given in Table 7.4. Fig. 28 of the Cycle 6 ALMA Primer gives visual representation of required array combinations as a function of \(\theta_{res}\) and \(\theta_{LAS}\).

Note: if the PI specifies a range of \(\theta_{res}\) there may be multiple 12-m array combinations which could achieve the Science Goal. For instance, C43-5 or C43-4 could serve for the extended 12-m configuration. This would imply C43-2 or C43-1, respectively, for the compact 12-m configurations; but C43-5 plus C43-1 is not an allowed combination. In this situation the initial observations will collapse the space of possibilities such that subsequent data will be taken in the most “compatible” array combinations. In this example, if C43-5 data are collected first, then every effort will be made to collect all of the extended array data in C43-5, and subsequent, compact configuration observations will be taken in a C43-2-like configuration. For purposes of total time calculations in Phase 1 the most expensive combination will be assumed.

Table 7.4 also shows the relative integration times that will be used by each array/configuration. While the overall integration time is set by the specific requirements of the Science Goal a PI defines, the relative integration times between its component configurations/arrays are determined by general considerations as discussed in Mason & Brogan (2013). The relative integration times do not depend on the details of the Science Goal beyond the fact that the Science Goal definition determines which combination of configurations or arrays is needed. Broadly speaking, because of their smaller collecting areas, the TP and 7-m Arrays require more on-source integration time than the 12-m Array. Furthermore it is evident that the higher surface brightness sensitivity of the more compact 12-m Array configurations requires longer 7-m Array integrations. These time ratios were calculated by simulating matched image sensitivities when the more extended of two arrays or configurations is \((u, v)\) tapered to give the same angular resolution as the more compact configuration. This approach is an incremental improvement to the approach used in previous Cycles, described in Mason & Brogan (2013). The actual time ratios adopted for ALMA observations, while guided by these calculations, are subject to further operational constraints. In particular, the ideal time ratios are often impractically large for the 7-m and TP Arrays when the most compact 12-m Arrays are in use. In previous Cycles, the time ratio was “capped” at a maximum value— typically 5 for adjacent arrays or configurations in a given imaging combination (with imaging combinations defined by Table 7.4). In Cycle 6, a less restrictive cap of 7 for adjacent arrays or configurations is adopted, but in practice this choice only affects the 7-m Array time ratio with respect to C43-1 and C43-2. Furthermore, all 7-m:12-m Array time ratios are rescaled by the factor by which the 7-m:C43-1 time ratio was reduced from its ideal value. This approach has the effect of more uniformly applying
Two additional considerations can impact the required integration times. First, the OT enforces a minimum on-source integration time in order to ensure that the telescope is used efficiently. The minimum allowed on-source time is taken to be the greater of 5 minutes and half of the integration time spent on the bandpass calibrator. Finally, the 7-m Array needs integrations of at least 1 hour to provide sufficient \((u, v)\) coverage for good image quality. Snapshot observations with the 7-m Array are therefore strongly discouraged (see also Section 7.5).

The operational constraint across different arrays and configurations with a modest and consistent impact on image characteristics.

Two additional considerations can impact the required integration times. First, the OT enforces a minimum on-source integration time in order to ensure that the telescope is used efficiently. The minimum allowed on-source time is taken to be the greater of 5 minutes and half of the integration time spent on the bandpass calibrator. Finally, the 7-m Array needs integrations of at least 1 hour to provide sufficient \((u, v)\) coverage for good image quality. Snapshot observations with the 7-m Array are therefore strongly discouraged (see also Section 7.5).
As previously noted, the configuration in which a given SB is executed may differ from the ideal "representative" configuration(s) that the OT assigned. One reason for this difference is that, to meet its demanding configuration schedule, ALMA antennas are almost daily being moved from one pad to another. Fundamentally, ALMA covers a very wide range of spatial scales (and wavelengths), and reconfiguring several dozen antennas is both operationally challenging and time consuming. Unanticipated proposal pressure distributions, periods of bad weather, weather, and hardware problems can require further flexible adaptation. In response to these circumstances it is not uncommon to merge two adjacent configurations into a single "hybrid" configuration that provides adequate if not optimal coverage of both of the two configurations. The scheduling subsystem accounts for the actual array configuration and matches projects with it based on their requested ✓LAS, allowing for the performance margins that ALMA’s Quality Assurance (QA2) process dictates. For observations that need a second, more compact 12-m Array configuration, observations of the second configuration are scheduled considering Angular Resolutions and (u, v) overlaps corresponding to the combinations of arrays listed in Table 7.4. Some flexibility in QA2 is provided by the use of Briggs’ weighting for QA2 imaging, a scheme which allows a tradeoff between angular resolution and sensitivity. This flexibility in scheduling and analysis is vital to achieve PI’s science goals efficiently.

### 7.9 Multi-array and Multi-configuration Imaging

If data have been collected in multiple arrays or configurations as described in Section 7.8, the data from each will be processed separately and delivered if they pass quality assurance. At present PI’s are responsible for combining these products together, although they are encouraged to contact their ARC for assistance if desired. PI’s may also refer to the ALMA M100 Science Verification CASA Guide for detailed guidance on data combination.

The recommended combination procedure comprises either one or two steps, depending on whether or not total power data are collected. First, the interferometric data will be imaged together in a single, “joint” deconvolution. This step can be done in CASA by passing all interferometric measurement sets— 7-m and one or more 12-m Array configurations— directly to the TCLEAN task. All 7-m and 12-m data delivered during Cycle 6 will be directly combinable.

If total power data are part of a project, then there is a second step to the combination procedure, where the the deconvolved interferometric (7-m+12-m) image cube is combined with the total power image cube. The recommended procedure is to do this by “feathering” them together. Feathering is a commonly used technique in radio imaging for combining two images together by forming a weighted sum of their Fourier transforms. The procedure is as follows:

1. The total power and interferometer images are Fourier transformed.
2. The beam from the total power image is Fourier transformed \(FTSDB(u, v)\), to be used as a weighting function. Alternatively one can specify some smaller portion of the total power antenna aperture, corresponding to a wider (single-dish) beam.
3. The Fourier transform of the interferometer image is multiplied by \(1 - FTSDB(u, v)\). This step down-weights the large spatial scale components of the interferometer map which are poorly measured, if measured at all.
4. The Fourier transform of the total power image is multiplied by the ratio of the volumes of the interferometer restoring beam to the single-dish beam, thereby putting the maps in the same units.
5. The results from 3 and 4 are added and Fourier transformed back to the image plane.

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6 [https://casaguides.nrao.edu/index.php/M100_Band3](https://casaguides.nrao.edu/index.php/M100_Band3)

7 Manually calibrated data deliveries from early ALMA Cycles that were calibrated using CASA versions earlier than 4.3 require a special procedure to put the 7-m and 12-m Array data weights on an equal basis before the data can be combined. The required steps are described in the “Data Weights and Combination” CASA guide [https://casaguides.nrao.edu/index.php/DataWeightsAndCombination](https://casaguides.nrao.edu/index.php/DataWeightsAndCombination). All ALMA data calibrated by the ALMA pipeline have correct weights, irrespective of Cycle number or CASA version.
The resulting cube provides high angular resolution due to the interferometric data, but will have more accurate large-scale features and total flux densities because the single dish data are present as well. An excellent discussion of feathering and a comparison to other techniques for combining single dish and interferometer data is given in Stanimirovic (2002).
7.9. MULTI-ARRAY AND MULTI-CONFIGURATION IMAGING

As an example, the \((u, v)\) coverages and \((u, v)\) plane sampling distributions for a 1-hour observation with the C43-1 configuration and 5-hour observation with the 7-m Array are shown in Figure 7.17 and Figure 7.18, respectively. The Cycle 6 7-m Array configuration with ten 7-m antennas provides \((u, v)\) measurements in the range 8-32 m and overlaps the 12-m Array configurations in the range \(R_{(u,v)} < 15\) m. Using the \(H\alpha\) emission of M51 as a hypothetical model (see Figure 7.14)

Figure 7.19: Images obtained using C43-3 (top left; 1 hour), C43-3 + C43-6 (top right; 1 + 3.3 hours) and C43-3 + C43-6 + 7-m (bottom left; 1 + 3.3 + 1.3 hours) Array combinations, and the image model itself (bottom right).

Figure 7.19 shows the images resulting from observing with a C43-compact configuration (C43-3) only; a compact plus extended 12-m configuration (C43-3 + C43-6); and finally of also adding the 7-m Array. Natural weighting was used and deconvolution was performed using the CLEAN algorithm. The recovery of larger scales is quite noticeable with the inclusion of the 7-m Array data. Using only the C43-6 configuration at 100 GHz provides an angular resolution of 0.27\(^{\prime}\). Combined with the C43-3 configuration, the angular resolution is lowered to 0.31\(^{\prime}\). Finally, the angular resolution with the full combination (C43-3 + C43-6 + 7-m) is 0.32\(^{\prime}\), very similar to that of only the two 12-m Array configuration data but 17% larger than the angular resolution obtained with the extended 12-m Array data alone. That difference is due to the respective \((u, v)\)-coverage and
weight of the array configurations in the combined dataset.

These simulations illustrate two general features of multi-configuration/array observations. First, they are effective in retrieving a wider range of spatial scales than are accessible to a single configuration or Array. Second, the addition of more compact configurations or arrays does tend to broaden the synthesized beam (PSF) slightly. This broadening can be mitigated by changing the \((u, v)\) weighting, e.g., via the robust parameter, at the expense of sensitivity.

One final effect to be aware of when combining data from multiple arrays or configurations is the flux bias which can result by the mismatch between the clean and dirty beam areas. This effect is discussed by Jorsater & van Moorsel (1995), who present an analytic correction and also note that it can be mitigated by deeper cleaning.

References
[6] Mason, B. & Brogan, C., 2013, Relative Integration Times for the ALMA Cycle 1 12-m, 7-m, and Total Power Arrays, ALMA memo 598
Chapter 8

Observing Modes

The ALMA Observing Modes are the set of capabilities that ALMA offers each cycle to its user community. Before offering a particular Mode to the community, it is implemented in the ALMA online software package and verified using the required hardware. Here, the software implementation is described along with a description of particular observing modes.

An observing proposal submitted to the ALMA archive will have an associated structure called the Observing Project that will accompany it along the whole length of its lifecycle. This structure is defined in the ALMA Project Data Model (APDM), which specifies all the relevant components and their contents needed for successful completion of a project. A summary view of the constituents of an Observing Project is shown in Figure 8.1.

8.1 Observing Project Structure

The organization of each science project is subdivided into a well-defined structure with clear hierarchical levels. At the bottom of this structure are the Scheduling Blocks (hereafter SB). An SB is the minimum set of instructions describing an ALMA observation that can be fully calibrated. Projects are broken down into a set of these fundamental units for flexibility, given the properties of the ALMA site and the continuously-varying status of the Observatory as a whole (including the weather), and to encapsulate the scientific objectives of the proposal into individual entities. After a proposal has been accepted for observation, the SBs to be executed during the observations are produced automatically by the ALMA Observing Tool (hereafter OT) during the Phase 2 stage of the project lifecycle. The SB may be edited, using the OT, in Phase 2 after the proposal has been accepted for observation and also during the project execution if necessary2.

An SB contains a large amount of information about what should be observed, the positions and velocities of the science targets, details of the correlator setup, the integration and cycle times of the different calibrations, etc. However, the SB itself does not control the observations as it is just a single set of XML instructions. Instead, the ALMA online software reads the SB and executes the observing script appropriate to the type of observation required e.g., single-dish, single field interferometry, etc. The SB has relatively little influence over the order in which the various sources are observed and does not describe all of the calibrations that should be performed (a prime example being the measurements of the system temperatures).

Each SB typically consists of set-ups, calibrations, and target observations that can be observed within 1.5 hours. The end of an SB execution may be specified in terms of a maximum amount of time or when certain well-defined science goals have been reached as specified in the Science Parameters section of the SB. An SB cannot be stopped and re-started mid-way. Therefore, an SB either runs to completion, fails mid-way, or is terminated by the Astronomer on Duty (hereafter AoD). Given the limited duration of the SBs, it is often necessary to

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1For a description of Phase 1 and Phase 2 see the Cycle 6 Proposer’s Guide
2SBs can also be created in an ad-hoc manner for commissioning purposes
Figure 8.1: Block diagram of an Observing Project from the point of view of the observation preparation (top) and internal hierarchical structure of the Scheduling Block in actual executions (bottom). All projects have the same ObsUnitSet (OUS) levels, the Science Goal OUS level, the Group OUS level and the Member OUS level. Scheduling blocks are attached to the Member OUS. Each time a Scheduling Block is executed, Control creates a new Execution Block structure (see text).
observe them several times to achieve the sensitivity required by the PI. The current maximum execution time for a particular SB has been chosen based on statistical measurements of the stability of the system. A system failure may prevent the online software from storing the data in the archive, resulting in the loss of the observed project data. Thus, executing SBs for more than two hours becomes unsafe from the operational perspective. To optimize this scheme, the concept of a session was implemented, which is the continuous execution of a single SB until a certain goal has been reached (see detailed description later on). Sessions have the added benefit that some of the calibrations can be shared among the different executions, which increases the overall observing efficiency. For Cycle 6, sessions are primarily used for polarization observations and VLBI observations.

While the SB is the smallest entity used for observing, the Observing Unit Set (hereafter OUS) is the smallest unit for data processing. The SB/Member OUS/Group OUS (hereafter MOUS and GOUS, respectively) are the smallest structures that hold science observations that need to be observed/processed/combined together. These data hierarchies therefore maximize observing flexibility and ensure that data gets processed as early and as fast as possible. In the OT, observations are divided into different Science Goals (see the ALMA Proposer’s Guide). In order to follow this structure, each GOUS is attached to a Science Goal OUS.

All SBs in a Science Goal have to be processed together to produce calibrated science products (e.g. images or data cubes). Usually, there will be one SB in each MOUS, but the latter may hold multiple executions of this SB. For polarization observations, all the SBs that belong to the same session should be grouped together for data processing. Unfortunately, this is not explicitly described in the SB.

If the calibrated science observations of a MOUS have to be combined with science observations of another MOUS, then these are grouped together into one GOUS. Otherwise, they will be placed in different GOUSs. A typical example of MOUSs that belong to the same GOUS are observations of the 12-m Array, the 7-m Array and the TP array. According to the method above, these observations would be in three separate MOUSs, but in the same GOUS.

Pipeline processing happens at the MOUS level, as soon as all observations of a MOUS are complete (see Chapter 13 for more details). It is expected that in future Cycles, processing will also happen at the GOUS level, in the case that the GOUS contains several MOUSs which have already been processed. The only other event that triggers data reduction is the end of an observing Cycle, when all the MOUSs and, potentially, GOUSs are reduced, irrespective of their degree of completion. As polarization (and other non-standard modes) observations do not currently fit this scheme, the data reduction will be done manually by the ALMA staff.

8.2 Execution Structure

Once a given SB has been selected for execution (by the Scheduler subsystem or by the AoD), it is read into the online Control software, and specifically by the Science Software Requirements subsystem (hereafter SSR). The SSR subsystem commands the lower level Control software to create an Execution Block (EB) structure that is attached to the SB. As an SB will be executed as many times as needed to fulfill the requirements, several EBs may exist for a given SB. Each EB contains a record of the parameters and conditions under which the SB was executed along with references to the acquired data. The internal hierarchical structure of the EB is also shown in Figure 8.1. The SSR subsystem constructs a sequential series of scans for each of the required calibrations and commands Control to execute them. Each scan execution is in fact carried out by breaking it down into a series of subscans, each of which is itself broken into a series of integrations (the correlator software only admits sequences of subscans). Although commands are issued at the scan/subscan level, the correlator output corresponds to a particular integration. In general, each calibration observation consists of a scan containing several subscans (e.g. the 5 data points of a pointing calibration constitutes a scan with five
subscans (see Chapter 10. Similarly, the integration time on a single science source between phase calibrations consists of one scan comprised of a number of subscans. To optimize the execution, scans may be organized in scan sequences which are passed to Control for execution. Scans and scan sequences can be of arbitrary length, depending on the characteristics of a given observation. But subscans are recommended to be 30 seconds or less in most cases. Integrations tend to be on the order of 1 to 10 seconds where the final value has to be an integer multiple of the correlator dump time. These values can be specified during the generation of SBs in Phase 2 with the OT; specifically, the subscan duration is specified in the SB target parameters section and the length of the correlator integration time and dump time is specified in the SB Spectral Setup. Calibration results from the Telescope Calibration subsystem (TelCal) and QuickLook (QL) pipeline are usually attached to a scan where one example is an antenna pointing result. The SSR and the Control subsystem are responsible for the creation of all the metadata needed downstream for data processing.

8.3 The Observing Process

For each of the SBs in a given project, a set of targets is specified (for more details, see the Cycle 6 OT User Manual). These targets can correspond to either calibration or science executions. Targets are organized in observing groups where the first group (Group 1) is always the initial calibration group, with subsequent groups detailing the science observations. All Science targets and relevant calibrators within a group are observed before the next group is started. An SB can have multiple groups.

All groups other than Group 1 are considered complete when all Science targets in the group have been observed for the requested time or have set below the elevation limit. After all groups are completed or the SB execution time limit is reached, the primary Phase calibrator for the group which triggered the SB execution time limit as well as any deferred calibrators from group 1 are observed.

Most observing modes, including single field interferometry, grouped source executions, pointed mosaics, and polarization use the same ALMA observing script, called the standard interferometry script. The set of necessary calibration measurements (e.g. flux, bandpass, etc) usually specified in group 1 are performed at the beginning of the observation sequence. Unless otherwise prefixed, the sources are selected at run-time by the SSR query algorithm using the parameters defined in the SB as input. If sources of sufficient quality for calibration are found the SB will execute.

Pointing, Atmospheric, and Sideband-Ratio calibrations are associated to the main (Bandpass, Amplitude, Polarization, and Phase) calibrations on an as-needed basis which is determined by the SSR at run-time. In the particular case of pointing calibration, it is usually verified before the amplitude and bandpass calibrators are observed, and again before the main observations of the science target and phase calibrator cycle. Within a group, the Science targets are each observed in turn until observation of the primary phase calibrator (the calibrator with shortest cycle time in the group) is required. A typical cycle time for the phase calibrator may be 7–10 minutes for standard observations, and 1–2 minutes for long-baseline observations (see Chapter 10). This process is repeated until the observing requirements are met, or the SB reaches its maximum execution time limit. Any additional (secondary) phase calibrators are observed as specified in the scheduling block. For a description of what each calibration entails, see Chapter 10.

The user has several options to select optimal calibrators:

- **System-defined Calibration** in the OT: let the OT set up default queries to the ALMA Calibrator Source Catalogue which will be used to select appropriate calibrators at run time. This is the recommended mode.

- **User-defined Calibration** in the OT: manually enter specific calibration sources or set up the queries using alternative values for the parameters, but this carries some risk (for example, calibrators will not be observed during the execution of a group if they are not visible at the time of the observation), and thus must be fully justified in the Technical Justification of the proposal.

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5In some cases (e.g. polarization), additional calibrations may be added into subsequent groups and they will be interleaved according to their cycle times. One example is the polarization calibration.
8.3.1 The Source Selection Algorithm

The ALMA system calibration for a given Scheduling Block is usually done using astronomical sources selected at runtime (exceptions are Long-Baseline and High-Frequency observations where some calibrators are pre-selected and hard-coded into the SBs, or SBs with User-Defined calibration where calibrators are pre-selected and hard-coded into the SBs). To optimally select appropriate calibration sources, selection criteria are implemented in the SSR using the following steps:

- A list of sources is retrieved from the ALMA Calibrator Source Catalogue based on the query center coordinates and the search radius defined in the SB.
- For each source in the list a search for flux measurements at the SB representative frequency is performed. If no flux measurements are found, the flux is extrapolated from other measurements using a spectral index of -0.7 to the SB representative frequency.
- The flux values are weighted based on proximity to the representative frequency, science target and time since the last measurement present in the source catalog.
- Expected SNR is calculated from flux and SB observing setup.
- The returned list of sources that passed the signal to noise criteria are ranked based on SNR, separation with respect to the query center and flux error.
- The final list is sorted based on ranking.

Using these rules, the criteria for all the ALMA calibrations were implemented adding specific requirements based on the nature of the target, such as:

- Bandpass calibrator: The bandpass calibrator integration time defined in the SB (typically 5-15 minutes, depending on e.g. band and spectral setup) and a number of dynamically determined parameters (e.g. the number of antennas in the array and the estimated system temperature) are used to determine a flux density threshold for the catalog search that would correspond to a minimum SNR per antenna of 50. If no suitable sources are found, the SNR limit is lowered. Other factors that are taken into account are elevation and shadowing. The search radius is typically 45 degrees from the query center, but if no suitable calibrator is found the search radius is increased.

- Phase calibrator: An SNR of 15 is used in the flux estimation. The required integration time on the phase calibrator is derived either using the widest spectral window bandwidth (usually 2 GHz), or the aggregate bandwidth when only narrower spectral set-ups are specified. Unless a value is set up in the SB, a search radius of 15 degrees from the query center is used as default (smaller cycle times and radii are used for the highest frequencies and longest baselines, as described in Chapter 10).

- Flux calibrator: A Solar System Object is searched in the first instance for flux calibration purposes. Otherwise, a search is performed among Grid Sources (see Chapter 10 for further details). The first observable grid source is picked up. A detailed list of the Solar System objects and grid sources used for flux calibration is available in Chapter 10.

To help checking the consistency of calibration results, the SSR has a logic to select at least two different sources for the sets of calibrators listed above.

8.4 Single Field Interferometry

Single field interferometry is the most basic form of observation that ALMA supports. It consists of standard calibration scans\(^6\) associated with the constituent calibration targets, and science observations of a single field.
CHAPTER 8. OBSERVING MODES

(Primary beam, see Chapter 3 for details). A typical observation will start with a bandpass calibration. The bandpass observation is executed to measure the spectral response of the system, and thus should be done on a bright source with simple spectral properties, such as a bright quasar with no emission or absorption lines and a reasonably flat spectrum.

The flux scale calibration will be performed next, which is intended to obtain the observed flux of a well known source, such as a solar system object. The observed flux will be used to compare with the established flux model of this object to obtain the scaling factor to be applied to all other sources in the SB. Ideally, the flux scale calibrator sources should be small in angular size with respect to the synthesized beam so as not to resolve the source structure and add uncertainties in the flux calibration. In practice, many solar system objects may be moderately resolved. In such cases, only a subset of the baselines may be used to estimate the flux based on accurate models of the objects available in the analysis software. From this result, a scaling factor is derived which can then be applied to the data at the data reduction stage. If no solar system object is observable, the SSR will pick the first available grid source.

The interferometer is a device for measuring the spatial coherence function (e.g., Clark 1999), so the phase information of the sources must be preserved and discerned from phase variations at the individual antenna elements. This is achieved with the phase calibration during the course of an observation. Since the phase is expected to change much more rapidly in time than the amplitude, phase calibrators will be observed more frequently than other calibrators. Because the phase varies on small scales on the sky, the calibrator must be as close as possible to the science target. The phase calibration observations are taken right before and after each observation of the science target, and the phase correction will be interpolated in time when applied to the science target. As the atmosphere fluctuates rapidly (especially at higher frequencies), radiometric observations of atmospheric water lines are also done to correct for these additional phase variations. These corrections can be applied online and offline (see Sections 10.4.3 for further details).

8.5 Pointed Mosaic Observations

Pointed mosaic observations are also offered for this Cycle. This mode enables a single Science Goal to cover a field of view larger than the primary beam by making observations of multiple single fields that overlap in spatial coverage by an amount specified in the SB. Up to 150 pointings are possible in a single SB. This limitation is set by the maximum execution time for a single SB and the necessity to finish all fields at least once within an execution. In pointed mosaic observations, each of the fields will be assigned a different field ID in the data, but the same source ID. Thus all fields will share their bandpass, amplitude and phase calibrations which are done as single fields. For larger fields that cannot be covered by 150 pointings, multiple Science Goals may need to be defined. The mosaic is arranged as a single scan composed by a number of subscans corresponding to the individual rows in the mosaic (see Section 5.5). Both the science target specific last mosaic pointing position and index within the science target list are handled by the SSR to ensure the next scan begins on the proper science target and proper offset position. The mosaic observations can be set up by specifying a Rectangular Field in the OT. If the user wishes to execute several small mosaics within an SB, the custom-mosaic under Multiple Pointings should be selected (see details in the OT guide).

8.6 Single Dish (Total Power) Observations

The purpose of adding data from single-dish observations using autocorrelations is to recover large scale emission from the science target that may have been spatially filtered-out by even the shortest baselines of the 7-m Array. For this reason, these observations are referred to as zero spacing. For convenience, these observations are also sometimes referred to as total power (hereafter TP), although in practice they are taken in autocorrelation mode rather than using a total power (square law) detector. Four 12 m antennas connected to the ACA Correlator are available for this purpose. For Cycle 6, only spectral-line observations are offered in this mode.

The observing script that executes this mode is called standard single-dish script. As with standard interferometry (see section 8.3), this script uses the SSR capabilities to perform some of the calibrations observations
(e.g., pointing) interferometrically. The SB of a TP observation consists of a group of calibrations followed by an On-The-Fly (OTF) observation of a rectangular area on the science target for line mapping with periodic offsets to a certain reference position observed for calibration purposes. For the OTF and reference position integrations, only the autocorrelation data are written to the ASDM in order to minimize data rate and size, whereas the rest of calibrations use cross-correlation information for analysis.

Most of the calibration scans are executed in group 2. The execution of group 2 starts up with a pointing calibration scan. Then, the OTF mapping of science targets are made with periodic executions of atmospheric calibration scans inserted. The atmospheric calibration is made at the reference position of the science target.

The OTF map is observed as a series of raster rows, scanning in the coordinate system specified in the SB. In Cycle 6, scans are taken either in longitudinal or latitudinal directions as specified in the OT. In future Cycles, scans may be taken in the two perpendicular directions in turn, to minimize scanning artifacts. The reference positions, assumed to be positions which are free of spectral-line emission, are specified either in absolute coordinates or as offsets from the map center. By default, the OTF map will cover an area half a beamwidth larger than the interferometric observations on all sides of the map. This will ensure that undersampling at the map edges does not affect the data combination process with the interferometric data. A raster row in these observations is defined as a subscan, with a maximum length per scan (consisting of some number of subscans) of 600 seconds. The reference position is observed as specified by its cycle time during the science target scans. Pointing is calibrated on a bright calibrator near the Science target with a frequency indicated in the SB, and atmospheric calibrations are taken every 10 minutes at the reference position to measure the system temperature.

The calibrated TP map of the science target will be in units of Kelvin, on the antenna temperature \( T^\star_a \) scale. Since all TP observations are expected to be combined with 12-m and 7-m Array data which come in calibrated units of Jy/beam, the TP data must also be converted into these units. This conversion from Kelvin to Jy/beam requires knowledge of the main beam efficiency \( \eta_{mb} \) and the beam size \( \theta \). These are measured separately for each project by obtaining a continuum map of a bright quasar or a planet with known flux. These “amplitude calibration” maps are reduced in the same way as the science observations, and the emission is compared with their model or observed flux (taken from the most recent calibrator survey measurements in the case of quasars) to calculate the Kelvin to Jansky/beam conversion factor, which is then applied to the science observations.

### 8.7 Polarization

The ALMA antennas have receivers with linearly polarized feeds followed by a waveguide with a polarization splitter.\(^7\)

In this way, the incoming radiation is separated into two orthogonal components (X and Y) which are down-converted and digitized independently. For each baseline, the digital signals are cross-correlated at the correlator where the outputs are the four cross-correlations \( XX, YY, XY, \) and \( YX \) (or \( V_{xx}, V_{yy}, V_{xy}, \) and \( V_{yx} \)). These four cross-correlations, as a function of the Stokes parameters, are ideally given by

\[
\begin{align*}
V_{xx} &= I + Q \\
V_{xy} &= U + iV \\
V_{yx} &= U - iV \\
V_{yy} &= I - Q
\end{align*}
\]

\(^7\)The minimum time per row in a map is 3 sec, and maximum speed at which data can be taken is usually limited by the correlator.

\(^8\)It is expected that during Cycle 6, these measurements will only be needed for the high-frequency ALMA bands, while for bands up to Band 7, the values will be automatically derived from fits to the aperture efficiency of each antenna as a function of time, elevation, outside temperature, surface accuracy, etc., which will be regularly monitored by the observatory.

\(^9\)For the ALMA receiver bands offered in polarization mode for Cycle 6, Bands 3, 4, 5 and 6 the polarization splitter is an Ortho-Mode-Transducer (OMT), and for Band 7 it is a polarized grid. See Chapter 4 for details.
where $I$, $Q$, $U$, and $V$ are the Stokes parameters. In an ideal world, one would be able to combine the observed cross correlated visibilities and recover the Stokes parameters as,

\[
\begin{align*}
I &= \frac{V_{xx} + V_{yy}}{2} \quad \text{(8.5)} \\
Q &= \frac{V_{xx} - V_{yy}}{2} \quad \text{(8.6)} \\
U &= \frac{V_{xy} + V_{yx}}{2} \quad \text{(8.7)} \\
V &= \frac{V_{xy} - V_{yx}}{2i} \quad \text{(8.8)}
\end{align*}
\]

From here, it is easy to see that the total intensity is in Stokes $I$ (function of the parallel hands XX and YY), and the linear polarization is described by Stokes $Q$ and $U$; while circular polarization is described by Stokes $V$. However, there are a number of issues that prevent directly using the measured cross-correlation in this simple way:

1. The splitting of the incoming radiation into orthogonal components is not perfect and small projections of one component into the other are produced. This is called the instrumental polarization or $D$-terms. The instrumental polarization is an antenna based quantity which also depends on the frequency and the band (cartridge design and external optics). Additionally, this quantity is measured in the frame of the antenna which is an elevation and azimuth frame (Alt/Az) and thus, it rotates with respect to the frame of the sky. This rotation introduces an angular dependence into the visibilities as a function of the parallactic angle ($\psi$). By design, the instrumental polarization is small (a few percent), but not negligible.

2. The signal path followed by each of the individual X and Y polarizations is slightly different, which introduces a small delay in the signal that needs to be accounted for. It has also been observed that the usage of the different local oscillators also introduces an offset in the XY-Phase.

3. The usual off-line calibration procedure arbitrarily sets the phase of a reference antenna to zero; this results in the inverse of its X-Y offset being imprinted on the cross-hand phases of the other antennas.

4. Other effects, such as position within the primary beam (off-axis polarization), are still under commissioning. Thus, only on-axis polarization is offered for Cycle 6 (see below for a detailed explanation.)

By taking into consideration the instrumental polarization and the parallactic angle dependence, the equations 8.1–8.4 for the visibilities can be re-written as

\[
\begin{align*}
V_{XX} &= (I + Q_\psi) + (U_\psi + iV) d_{Xj} + d_{Xj} (U_\psi - iV) + d_{Xj} (I - Q_\psi) d_{Xji} \quad \text{(8.9)} \\
V_{XY} &= (I + Q_\psi) d_{yi} + (U_\psi + iV) + d_{Yj} (U_\psi - iV) d_{Yji} + d_{Yj} (I - Q_\psi) \quad \text{(8.10)} \\
V_{YX} &= d_{Yj} (I + Q_\psi) + d_{Yj} (U_\psi + iV) d_{Xj} + (U_\psi - iV) + (I - Q_\psi) d_{Xji} \quad \text{(8.11)} \\
V_{YY} &= d_{Yj} (I + Q_\psi) d_{Yji} + d_{Yj} (U_\psi + iV) + (U_\psi - iV) d_{Yji} + (I - Q_\psi) \quad \text{(8.12)}
\end{align*}
\]

where $d_{Xj}$ are the $D$-terms as a function of polarization and antenna, the asterisk denotes complex conjugates, and $U_\psi$ and $Q_\psi$ are the Stokes parameters as a function of the parallactic angle. Figure 8.2 shows an example of $D$-terms in Band 3. The $D$-term level is typically a few percent at Bands 3, 4, 5, 6, and 7 on axis with some variations over frequency. Without any $D$-term calibration, an unpolarized source may appear to be polarized at the 1% level. The most straightforward way to calibrate the $D$-terms is to observe an unpolarized source. The cross-hands output will be purely $D$-terms since Stokes Q, U, and V will have no signal. However, bright unpolarized sources are rarely found.

\footnote{A resonance has been documented for Band 6 between the feed and the OMT. This resonance introduces a small ripple in the $D$-term solution which increases the leakage amplitude by a small amount, but still only a few percent.}
Figure 8.2: The $D$-term plots of antennae DA42 (top), DV03 (middle), and PM01 (bottom) for an observation in Band 3. The vertical axis is the fraction of the input signal voltage in one polarization that leaks into the output of the other polarization in voltage units and the horizontal axis is the frequency in MHz. The blue and red symbols represent $D_X$ (the fraction of $Y$ polarization signal that leaks into the $X$ polarization) and $D_Y$ (the fraction of $X$ polarization signal that leaks into the $Y$ polarization), respectively.

ALMA therefore normally uses an unresolved quasar as a polarization calibration source. These are monitored to ensure they are bright enough, and the exact polarization is determined during data reduction provided that the source is observed over a wide enough range ($> 60^\circ$) of parallactic angles to separate the effects of the source polarization, the $D$-terms and the cross-hands phase spectra and delay. Fig 8.3 shows that the rate of change of parallactic angle is fastest near transit and slowest for low elevation sources. The $D$-terms may have a slight elevation dependence, and therefore it is favorable for the polarization calibrator to be close to the science target. The ALMA Observatory will select appropriate calibrators for each science target, but the users must take into account these limitations when planning the observation.

With the current calibration scheme, linear polarization imaging of a compact source on-axis, at the level of 0.1% in fractional linear polarization is feasible in TDM mode (see below). The accuracy of absolute polarization position angle will be nominally 6 degrees (based on initial polarization mode verification results), with a contribution of 2 degrees coming from the specification on the orientation of the receiver feeds, and an additional scaling of the error with the number of antennas contributing to the final image. Of course, these levels of accuracy will only be reached if there is sufficient signal-to-noise in the polarised emission. For circular polarization we have determined that the minimum detectable degree of circular polarization, defined as three times the systematic calibration uncertainty, is currently 1.8% of the peak flux for both TDM and FDM observations. This level of accuracy is currently only on-axis, as off-axis circular polarimetry is still being commissioned. Thus, the users should consider this current level of accuracy when proposing.

As the $D$-term component also arises from the off-axis geometry of feed horn, antenna illumination, and the alignment of optics, it will also vary across the primary beam pattern (see Chapter 4). Generally, the $D$-term level becomes larger when increasing the offset from the beam center. This is the so-called off-axis instrumental polarization. In Cycle 6, the off-axis instrumental polarization calibration will not be employed. The expected minimum detectable degree of linear polarization, defined as three times the systematic calibration uncertainty, is 0.1% (1%) for compact sources (i.e., within the inner 1/3 of the primary beam FWHM) and 0.3% (3%) for extended sources for TDM (FDM) observations, respectively. The minimum detectable degree of circular polarization is 1.8% of the peak flux for both TDM and FDM observations. Note that the systematic calibration uncertainty can degrade by a factor of $\sim 2$ depending on the choice of calibrator, parallactic angle coverage, etc. With the current calibration scheme, linear polarization imaging of a compact source on-axis in TDM mode is feasible at the level of 0.1% (3 sigma) fractional polarization for the very brightest calibrators, and 0.2% (3 sigma) level for a typical observation. Users must justify that their science goals can be met by sources exceeding these lower limits.
Spectral line polarimetry is also offered in Bands 3, 4, 5, 6, and 7 at arbitrary frequencies. In the case of circular polarization, the current minimum detectable fractional polarization of 1.8% also applies to FDM mode. This should be considered for Zeeman-like experiments. Figure 8.4 shows the D-term solutions of a representative antenna (DV22) obtained during commissioning. Figure 8.5 shows a zoomed-in view in frequency of the D-term solutions shown in Figure 8.4. These figures show that the spectral shape of the D-term solutions is smooth and no spur-like structure is seen even in highest frequency resolution mode (30.5 kHz at the right bottom panel of Figure 8.5). Figures 8.4 and 8.5 clearly demonstrate that the instrumental polarization is mostly spectral resolution independent for these datasets. No mosaic, ACA, and TP Array observations are offered. Further, a minimum execution time of 3 hours will be imposed to ensure sufficient parallactic angle coverage for calibration.

8.7.1 Sessions

Though not directly related to polarization itself, sessions are the observing scheme designed to observe polarization projects with ALMA. In order to allow the execution of long programs and to avoid the execution of long scheduling blocks (more than 2 hours), the concept of a session was implemented into the SSR. Given the current stability of the system, it was considered an unnecessary risk to have more than 1.5 hours scheduling blocks, because in the event of a system failure all data may be lost. Thus, a session scheme was designed as an alternative. A session is defined as the continuous execution of the same SB until the scientific criteria are met. A session will manage the cycle time of each of the calibrators from the starting point of the session, i.e. the first execution of the SB. In this way, calibrations are only executed when needed, avoiding unnecessary observations and giving an additional level of optimization for ALMA. By default, calibrations such as flux calibration will be done once per session and bandpass calibration every hour, thus saving observing time. Also, the phase calibration science target loop will be interrupted (preserving phase calibration) when an additional calibration is needed (e.g., polarization). The mode that currently exploits the session scheme in ALMA is the observation of polarization projects. Because for an Alt/Az antenna the frame of the sky rotates, the calibration of the instrumental polarization requires sampling a strong compact source as a function of parallactic angle approximately every 30 minutes. The session will remember the last time the polarization calibrator was observed and interleave the calibration when needed. In general, one can achieve the required parallactic angle coverage by running the polarization SB 2 or 3 times, which results in between 3 to 4 hours of observation, which is a perfect example of the session scheme. Other cases in which the session scheme might be useful are large mosaics, surveys (multi-target), and large single dish raster maps, but those cases are not yet offered in session mode for Cycle 6. The implementation of sessions is done by keeping information about the previous execution of the SB into the runtime memory.

---

11 The cycle time parameters are user controlled and can be explicitly specified in the OT target parameters section.
8.8. MULTIPLE REGION MODES

In the frequency division modes (FDM) of the correlator (see Section 5.1.2), the final spectrum is synthesized using individual filters 62.5 MHz wide. When the total bandwidth is between 125 MHz and 1 GHz, it is possible to move these individual filter positions to create spectral windows covering a number of disjointed spectral regions. This is called the Multiple Region Modes. The constraints are:

- The spectral window width must be a multiple of 62.5 MHz.
- The aggregate width of these spectral windows must be equal to that of the original bandwidth selected.
- The spectral windows must all fit within the 2 GHz baseband used.
- The other parameters (resolution, polarization and sensitivity options) must be the same for all spectral windows.

This mode is useful when a number of line features which require high spectral resolution are spread across the IF bandwidth. Since the filters have a unit width of 62.5 MHz, if the user chooses a mode with 250 MHz total bandwidth, it is possible to place four separate windows, each with 62.5 MHz width, anywhere within the 2 GHz baseband. In Cycle 6, a maximum of four spectral windows per baseband will be offered.

8.9 Observation of Ephemeris Objects

Observation of solar system objects (with the exception of the Sun) is supported as in previous Cycles. Several well known solar system objects including planets, satellites and asteroids can be selected from a pulldown menu in the Observing Tool. For other sources, including non-sidereal objects, an external ephemeris file can be supplied as an input. The ephemeris file must be in JPL Horizons format. A typical ephemeris file may consist of the date (time), Right Ascension, Declination, range and range rate, for example:
Figure 8.5: A close-up view in frequency space of the real part of D-term shown in Figure 8.4. Blue circles and red crosses are showing the D-terms obtained from setup #1 (FDM, 2 GHz bandwidth) and setup #2 (FDM, 500 and 62.5 MHz bandwidth respectively).

Date__ (UT)__ HR:MN R.A._ (ICRF/J2000.0) DEC delta deldot

<table>
<thead>
<tr>
<th>Date</th>
<th>HR:MN</th>
<th>R.A. (ICRF/J2000.0) DEC</th>
<th>delta</th>
<th>deldot</th>
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</tr>
</tbody>
</table>

More information on the format and precision needed is available in the ALMA Observing Tool documentation.

8.10 Solar Observations

Solar observing was part of the original science case for ALMA. It has recently been thoroughly tested and verified and released for PI science observations. In Cycle 6, solar observing with ALMA will be supported, albeit with limitations. In particular, the following conditions apply to solar observing in Cycle 6:

- Only Band 3 and Band 6 continuum observations using the following default frequencies are offered:

```
<table>
<thead>
<tr>
<th>Band</th>
<th>LD Freq.</th>
<th>LSB</th>
<th>BB1</th>
<th>BB2</th>
<th>BB3</th>
<th>BB4</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>100 GHz</td>
<td>92-94 GHz</td>
<td>94-96 GHz</td>
<td>104-106 GHz</td>
<td>105-108 GHz</td>
<td></td>
</tr>
<tr>
<td>6</td>
<td>239 GHz</td>
<td>229-231 GHz</td>
<td>231-233 GHz</td>
<td>245-247 GHz</td>
<td>247-249 GHz</td>
<td></td>
</tr>
</tbody>
</table>
```

- Interferometric observations must be done using the TDM mode (see section 5.1.1); frequencies are fixed to 2 GHz-wide spectral windows centered on the frequencies shown in the table above. The spectral-line FDM observing mode is not offered for solar observations (Section 5.1.2).

- Simultaneous observations with Bands 3 and 6 are not offered: each execution block can only include one band.

- Observations will be carried out while the array is in a compact configuration corresponding to C43-1, C43-2, or C43-3.
8.10. SOLAR OBSERVATIONS

- Because the WVR receivers are saturated when the dishes point at the Sun, on-line WVR phase correction will not be applied and off-line WVR correction for on-source (solar) data is not possible (see Section 8.10.2).
- Interferometric observations will use a combined array comprising both 12 m and 7 m dishes and will be processed with the baseline correlator (Section 5.1).
- To minimize shadowing of 7-m antennas, observations will be carried out between 13:00 UT and 20:00 UT.
- Mosaicing of larger FOVs can be carried out with up to 150 different pointing offsets relative to the phase center specified in the ephemeris.
- As for the rest of ALMA observations, a single observation (an Execution Block) cannot exceed 2 hours, which will include the time overheads for bandpass and flux calibration. These calibration overheads amount to about 25 mins.
- The integration time of interferometric observations is fixed to 1 second.
- Observations may be performed using dual linear polarization (XX, YY) or single polarization (XX) correlations; full polarization measurements are not currently offered for solar observations.
- Total-Power full-Sun fast-scanning single-dish observations are offered as context and to recover the largest angular scales for interferometric observations; proposals requesting only Total-Power single-dish observations will not be accepted in Cycle 6.
- Single-dish Total-Power mapping will be “monomode” (double-circle scan pattern) with a 2400 arcsec-diameter circular FOV centered on the solar disk (see below).
- Single-dish full-Sun will not be executed when the Sun is at elevations above 70° because the required fast-scan azimuth slew speeds are too high.
- The time cadence of full-sun images obtained from total power observations is fixed to about 7 minutes for Band 3 and 10 minutes for Band 6.

8.10.1 Solar Observing Modes

Solar observations with ALMA are possible because the surface of the antennas is designed to scatter the optical and IR radiation to an extent that the subreflector and other elements in the optical path are not damaged or degraded. However, additional steps must be taken to allow useful observations of the Sun to be made. ALMA receivers are designed for a maximum RF signal corresponding to an effective brightness of about 800 K at the receiver input. Since the quiet Sun has a temperature of \( \sim 5000-7000 \) K at ALMA frequencies, the solar signal must be attenuated or the receiver gain must be reduced to ensure that receivers remain linear, or nearly so.

The initial solution adopted by ALMA was the use of a “solar filter” (SF) that is mounted on the Amplitude Calibration Device (ACD) of each antenna (Section A.5). When placed in the optical path, the solar filter is required to attenuate the signal by \( 4+2\lambda_{mm} \) dB with a return loss of -25 dB (-20 dB for \( \nu > 400 \) GHz) and a cross polarization induced by the filter of -15 dB, or less. While the use of solar filters has been demonstrated to work, their use introduces several disadvantages, not the least of which are cumbersome calibration procedures.

Yagoubov (2013b, 2014) pointed out that the ALMA SIS mixers could be de-biased to reduce the mixer gain and effectively increase the saturation level to a degree that allows solar observations without the use of the solar filters, at least for non-flaring conditions on the Sun. These produce lower conversion gain and since the dynamic range scales roughly inversely with gain, these settings can handle larger signal levels before saturating. In addition to the SIS bias voltage, the local oscillator (LO) power can be altered in order to further modify the receiver performance. However, LO power settings have not yet been fully optimized for solar observing.

Two so-called “Mixer De-Tuned” or “Mixer De-Biased” settings have been adopted for solar observations in Band 3 and Band 6 for Cycle 6. These are referred to as solar observing modes MD1 and MD2.
• **Band 3**: MD1 mode uses a bias voltage that sets the SIS mixer to the 2nd photonic step below the voltage gap whereas MD2 mode employs a bias voltage corresponding to the 2nd photonic step above the gap. For MD1 mode, the Band 3 receiver temperature suffers a modest increase, to $\sim 50$ K, and receiver compression is limited to $\sim 10\%$. For MD2 mode, however, the Band 3 receiver temperature increases significantly, to $800$ K, but receiver performance is believed to be essentially linear.

• **Band 6**: MD1 mode settings are nominal for ALMA; i.e., the bias voltage is that used under normal observing conditions. Receiver compression is again on the order of $\sim 10\%$. MD2 mode employs a bias voltage corresponding to the 1st photonic step above the voltage gap. Again, the receiver temperature is significantly higher, $800$ K, but the receiver is essentially linear.

The MD1 mode is considered to be a “quiet Sun” mode (coronal holes, the solar limb, quiescent filaments, prominences, quiet areas outside of active regions) whereas the MD2 mode is recommended for the “active Sun” (active regions, active filaments, science objectives that require accurate photometry).

### 8.10.2 Array Observations

The Sun is an extremely large source compared with the Primary Beam of either the 7 m or 12 m antennas (see Chapter on Antennas). The Primary Beam is filled with complex emission when pointing at the Sun, as are the beam sidelobes. The ALMA array ultimately measures the brightness temperature contrast relative to the background Sun, which is resolved out by the array. As noted earlier, in order to recover the absolute brightness temperature of solar targets, it is necessary to include not only interferometric observations (by the 7 m and 12 m antennas), but also Total Power measurements made with a single dish. Single dish fast-scan mapping of the Sun in Total Power mode is addressed in the next subsection.

An advantage to using MD mode observing is that the water vapor radiometers (WVRs), which are used to correct differential phase errors introduced by precipitable water vapor over the array, are not blocked by the ACD. They can therefore be used, in principle, to make such corrections to solar data. Unfortunately, unless the optical depth of the sky is $\sim 2.5$ or more, which would represent highly non-optimum observing conditions, the WVRs saturate on the Sun. Until the WVRs are modified or replaced to increase their dynamic range to accommodate the Sun, phase corrections based on WVR measurements will not generally be possible when pointed at the Sun. For this reason, solar observations are currently restricted to compact array configurations to minimize such phase errors.

Another consideration, again regardless of whether SFs or the MD modes are used (see Section 8.10.1), is the system IF attenuator settings. The input power changes significantly as the antennas move from the (solar) source to a calibrator and back. The IF chain has two variable attenuators (in steps of 0.5 dB) to ensure that signal levels remain within nominal limits: one in the IF Switch and one in the IF Processor. A concern is whether the variable attenuators themselves introduce unacceptable (differential) phase variation between source and calibrator settings, thereby corrupting phase calibration referenced against suitable sidereal calibrators; and whether there are differences between the spectral window bandpass response between source and calibrator scans as a result of attenuator settings. Careful testing has shown that this should not be a significant concern. While these tests show that phase shifts caused by the attenuation level changes do in practice difference out, verification that this is the case cannot be checked from observing data obtained using the standard solar Scheduling Block. As a check, the observatory will carry out a test observation of a calibrator source using normal and MD attenuation levels before solar observations begin on a given day or at least once before a campaign program.

Bandpass calibration is carried out in the usual manner using MD modes: i.e., a strong calibrator is observed in an MD mode and the bandpass solution is obtained. Bandpass shape and stability were checked for MD modes and attenuator states in Bands 3 and 6. It was found that perturbations to bandpass amplitudes and phases were small. For the IF Switch and IF Processor settings adopted for MD mode observing it was found that the RMS difference between bandpass phases for an MD attenuator state and the nominal attenuator state was generally a fraction of a degree for both Band 3 and Band 6, the maximum being $1.2$ deg. Similarly, the
normalized amplitude difference was typically a fraction of 1%. No explicit correction for differential bandpass is needed.

In the non-solar case, the antenna temperature \((T_{\text{ant}})\) is small compared to the system temperature, and \(T_{\text{ant}}\) can therefore be neglected for amplitude calibration (see Chapter 10). In contrast, unlike most cosmic sources, the antenna temperature of the Sun is large (~7000 K at 100 GHz). It is therefore necessary to measure both the system temperature and the antenna temperature when pointing at the Sun in order to compute the System Equivalent Flux Density (SEFD) to correctly scale visibility amplitudes.

To estimate the antenna temperature \(T_{\alpha}\) on the Sun, “single-dish” measurements must be performed using all antennas of the array. Specifically, the standard observing sequence for solar interferometric observations will include the following measurements:

- a “sky” observation \(P_{\text{sky}}\), offset from (by typically 2°) and at the same elevation as, the target (Sun)
- a “cold” load observation \(P_{\text{cold}}\) (also known as the “ambient” load), in which an absorber at the temperature of the thermally-controlled receiver cabin (nominally 20 °C) fills the beam path
- a “hot” load observation \(P_{\text{hot}}\), in which an absorber heated to ~70 °C fills the beam path
- a “zero” level measurement \(P_{\text{zero}}\), which reports the levels in the detectors when no power is being supplied

Then the telescope moves to the target (Sun) where the IF attenuation levels are set appropriate to the input power. After the target scan, the telescope again moves to the “sky” position and takes another measurement, called the “off” measurement \(P_{\text{off}}\), without changing the IF attenuation.

The antenna temperature of the science target is then given by:

\[
T_{\alpha} = (P_{\text{sun}} - P_{\text{off}}) \frac{(P_{\text{sky}} - P_{\text{zero}})}{(P_{\text{off}} - P_{\text{zero}})(P_{\text{hot}} - P_{\text{cold}})(T_{\text{hot}} - T_{\text{cold}})}
\]

The autocorrelation data output from the baseline correlator cannot be used for this measurement because it has insufficient dynamic range to measure \(P_{\text{zero}}\). Instead, the necessary measurements rely on Total Power data obtained by the baseband detectors.

### 8.10.3 Single-Dish Mapping

Fast scanning observations are ideal for recovering the flux or brightness distribution on angular scales ranging from the ALMA primary beam width to the scale of the target in question (typically a few arcminutes), or up to the full disk of the Sun. Briefly, fast-scan mapping entails making Total Power (and more recently, autocorrelation measurements) as the telescope pointing is driven continuously and smoothly through a sampling pattern on the target that avoids sudden acceleration or deceleration of the antenna drive motors. A major advantage of fast scanning is that it minimizes the impact of atmospheric variation, and the full solar disk can be mapped in as short as 7 minutes.

While various types of scan patterns have been developed and tested for ALMA dishes, Cycle 6 supports the use of a “double circle” pattern, which maps a circular region on the sky. The double-circle pattern is particularly well suited for full-disk mapping because its coverage matches the shape of the solar disk and it repeatedly revisits the region of the center of the disk, allowing atmospheric opacity variations to be corrected for (see Figure 8.6). Standard observing procedures include focus and pointing checks on suitable sources prior to the fast-scan mapping.

### 8.10.4 ALMA Solar Ephemeris Generator Tool

The tool (URL http://celestialscenes.com/alma/coords/CoordTool.html) was developed by the Czech ARC Node to be used together with the OT in preparing solar observations and was tested during the December
2015 solar campaign. It is a javascript based application which runs in any modern browser and on many different operating systems.

User interface consists of several panels. In the input panel, it is possible to select the latest SDO/AIA image in several bands or to upload the users’ own FITS files. Currently, only uncompressed FITS with defined solar WCS keywords using CROTA2 formalism are supported. Visualization and display panels enable the user to pan and zoom in/out the region of interest or to show/hide coordinate grids and tweak the image display by false coloring and level scaling functions typical of other astronomical software packages and the OT.

The actual pointing is done by clicking the desired feature with a green cross marking the position. The coordinates of the pointing are displayed in several coordinate systems inside the pointing panel, where it is also possible to manually define the pointing.

Finally, in the observation panel, the user defines start and end times of the observation and differential rotation profile which will be used for generation of an OT-compatible ephemeris file. There are several rotation profiles to choose from or the user can define his/her own. Clicking Generate ephemeris file for OT will display the generated file which can be downloaded by following the Download data below link and then imported into the OT.

The ephemeris file is generated from the JPL Horizons file which the tool queries directly from Horizons website. However, during Cycle 4 regular observing sessions there was a period when JPL Horizons site went offline and hence the Ephemeris Tool was unable to function properly. Work is currently underway to enable precise ephemeris calculations in the Ephemeris Tool even when JPL site goes down.

The ALMA Solar Ephemeris Generator Tool comes with a user manual and it is served on several sites for backup purposes. Currently, there are no plans to include the tool functionality into the OT.

### 8.11 VLBI Observing Mode

The Very Long Baseline Interferometry (VLBI) Observing Mode (VOM) is a variant of the standard interferometry mode with some additional capabilities to allow ALMA to participate in global VLBI networks operating at millimeter and submillimeter wavelengths. Since Cycle 4, ALMA VLBI mode observing has been offered in
8.11. VLBi Observing Mode

Band 3 in conjunction with the Global Millimeter VLBI Array (GMVA)\textsuperscript{12} and in Band 6 in conjunction with the Event Horizon Telescope (EHT) network\textsuperscript{13}. For details about proposing for these opportunities, see the ALMA Cycle 6 Call for Proposals (http://almascience.org/proposing/call-for-proposals/).

For all VOM observations in Cycle 6, ALMA will be operated as a phased array of 12 m antennas (expected to contain up to 39 phased antennas). With 39 12 m antennas, phased ALMA is equivalent to an antenna of diameter of 75 m, neglecting efficiency losses. From Table 9.3 the aperture efficiencies in Band 3 and Band 6 are 0.71 and 0.68, and $T_{\text{sys}}$ is 70 K and 100 K, respectively. Since $Gain \equiv A_{\text{eff}}/(2k_B)$ where $k_B$ is the Boltzmann constant and $SEFD \equiv T_{\text{sys}}/Gain$. Corresponding estimates for these quantities are:

\[
\begin{align*}
Gain &= 0.71\pi(37^2)/(2k_B) \times 10^{-26} = 1.13 \text{ K/Jy} \\
SEFD &= 62 \text{ Jy} \quad \text{(Band 3)} \\
Gain &= 0.68\pi(37^2)/(2k_B) \times 10^{-26} = 1.06 \text{ K/Jy} \\
SEFD &= 94 \text{ Jy} \quad \text{(Band 6)}
\end{align*}
\]

Note that the VOM is not compatible with the use of subarrays.

When ALMA is operated as a VLBI station, all standard interferometry data products are also output by the ALMA correlator. Therefore, during any VOM execution an observer simultaneously obtains data equivalent to a standard ALMA interferometric observation of the science target, in addition to the VLBI data products. The latter are recorded independently on Mark 6 VLBI recording systems and will be shipped to common sites (MIT Haystack Observatory and MPIfR Bonn) for correlation with the other VLBI site data before delivery to the PI (see below). Some important details of the VOM are elaborated in the next subsections.

8.11.1 General VLBI Considerations

The supporting VLBI networks (i.e., the GMVA for Band 3 and the EHT for Band 6), together with ALMA, place certain restrictions on the observing process; these are largely driven by the complexity of orchestrating the simultaneous, reliable execution of a common VLBI schedule, as well as the calibration requirements of the individual telescopes that comprise the global array. In practice, for the ALMA VLBI observer there will be very little difference between planning an observation in the two observing bands, aside from details such as differing bandwidths and frequencies.

Successful proposals will be passed to a network Scheduler who builds the common VLBI schedules (encoded in a so-called VLBI EXperiment, or VEX file built with SCHED\textsuperscript{14}) for a particular observing campaign. In these schedules, some number of contiguous hours will be devoted to each set of science targets and necessary VLBI calibrators. The schedule will be worked out between the Scheduler and the PI well in advance of the observation. Most sites participating in the VLBI network also have a “friend” of VLBI who assists the Scheduler in working out details specific to that site. Sample VLBI schedules from previous campaigns are available from the aforementioned GMVA website.

In general, the required time on the VLBI science target(s) and VLBI calibrators will be broken into a number of VLBI scans of several minutes duration, separated by gaps of the order of several minutes to allow time for local calibrations (pointing, system temperature measurements, etc.) and antenna slew time sufficient to meet the needs of each observatory. VLBI observations that include ALMA will require that the gaps between VLBI scans are of sufficient length to allow for the performance of the calibrations normally required for standard interferometric observing, including flux, gain, bandpass, and polarimetric calibrators (see below). Since all VLBI science targets for Cycle 6 are required to be sufficiently bright to permit phase self-calibration ($\geq 500$ mJy), observations of a gain calibrator will be required only every 20-30 minutes to allow calibration of the amplitudes.

It is important to recognize that most VLBI stations currently record circular polarization (either left, right, or dual), in contrast to ALMA which records dual linear polarizations. Consequently, every VLBI mode

\textsuperscript{12}See http://www3.mpifr-bonn.mpg.de/div/vlbi/globalmm/

\textsuperscript{13}See http://www.eventhorizontelescope.org/

\textsuperscript{14}http://www.aoc.nrao.edu/software/sched
**observation with ALMA must be treated as a polarization observation** in order that a transformation from linear to circular polarization can be made correctly during post-correlation processing. A special software tool, **PolConvert**, is available for this purpose. Robust execution of **PolConvert** requires a minimum session duration of 3 hours (session length, not time on science target) in order to achieve adequate parallactic angle coverage on the polarization calibrator sources for the computation of polarimetric “leakage” or D terms (see Section 8.7). Visualization and display panels enable user to pan and zoom in/out the region of interest.

Other VLBI observatories have receiver capabilities that differ from ALMA. One complication is that the sample rate used at ALMA is non-traditional for VLBI, so care must be taken in the experiment set-up to ensure that the VLBI data can be correlated with the other stations. The VOM spectral set-ups chosen for each band in Cycle 6 were selected to address these constraints and therefore are fixed. For Band 3 (at 3mm) there are four 1.875 GHz wide bands centered at 86.268, 88.268, 98.268 and 100.268 GHz, respectively. In Band 6 (at 1 mm) the bands are centered at 213.1, 215.1, 227.1 and 229.1 GHz, respectively.

A second complication is that not all millimeter VLBI sites are capable of matching ALMA’s full recording bandwidth. For the first time in Cycle 6, the EHT is planning to offer 64 Gbps recording at all sites and is therefore expected be able to record all four bands and match for the first time ALMA’s full frequency coverage and bandwidth. However, stations of the GMVA are limited to 2 Gbps recording rates. The GMVA sites will therefore be tuned to a 256 MHz band centered at 86.124 GHz, so only the first ALMA baseband is available for VLBI.

Calibration of VLBI data generally requires observations of VLBI calibration targets interleaved with the main science target. At ALMA, these VLBI calibration measurements are carried out in a manner analogous to the observations of the science targets (i.e., using the VOM and active phasing of the ALMA array). The VLBI calibrators are generally chosen for their utility across the entire global array to meet a particular calibration need (e.g., fringe-finding, bandpass calibration, polarization calibration) and need not be the same calibrators as observed for calibration of the (ALMA-only) standard interferometric data.

During VLBI observations, the observatories are usually not interconnected by a network, so recordings (onto disk modules) of the antenna baseband signals are forwarded to common correlation facilities managed by the appropriate network (MPIfR Bonn and MIT Haystack Observatory for the Band 3 and Band 6 data, respectively). These sites are responsible for the correlation of the VLBI data and delivery of the correlated VLBI data products to the observer, as discussed in Section 8.11.2. Correlation is currently performed using the DIFX correlation software.\(^{15}\)

### 8.11.2 ALMA Considerations for VLBI

As mentioned in the previous section, all VOM observations should be viewed as polarimetric observations. The reader is referred to Section 8.7 for guidance on this.

In order to participate in a VLBI observation, there are two special requirements at ALMA. The first is that the ALMA control system must tune receivers appropriately and point to the targets specified in the VEX schedule at the appropriate times. The second is that the signals from the ALMA antennas must be coherently summed up, decimated, and recorded on media suitable for delivery to the common VLBI correlator. In order to form this coherent sum, the (arbitrary) phases of the signals at the antennas need to be adjusted during the observation so that they can be added “in phase”. The necessary signal processing for this takes place within special cards and circuitry within the ALMA correlator as well as with special software within the ALMA control system which were developed by the ALMA Phasing Project (APP).\(^{16}\) Further information on the APP and the ALMA Phasing System (hereafter APS) is described in Matthews et al. 2018, PASP, in press. Details of this are also presented in the following sections, with particular emphasis on those aspects which are essential either for the proposer or data analyst. The discussion begins with a description of the VOM Scheduling Block in the next section. Section 8.11.2 discusses the scan sequences used by the VOM. That is followed by a section with some important details about how the APS works (Section 8.11.2). Finally, Section 8.11.2 makes a few comments about the analysis of VLBI data (a full discussion is outside the scope of this document).


\(^{16}\) An NSF MRI/ALMA North America Development Project.
8.11. VLBI OBSERVING MODE

VLBI Observing Mode Schedule Block and Execution

As with other modes, VOM observations are executed with a Scheduling Block which provides observation-specific details to an SSR observing script. In the case of the VOM, the script is called StandardVLBI.py, although, as noted elsewhere, VLBI is a non-standard observing mode in Cycle 6.

The Scheduling Block is created from the observing project and the VEX file using a special tool called VEX2VOM. This tool reconciles the Phase 2 Schedule Block with the schedule and targets specified in the VEX file and thus provides the detailed operating instructions for the SSR observing script. The preparation of the final schedule block with VEX2VOM is done by the Friend of VLBI or the AoD prior to execution when all the necessary details of the observing array are known.

Like the StandardInterferometry.py script, the StandardVLBI.py ensures that the necessary calibrations as specified by the OT are performed with the desired repetition rate. Unlike the StandardInterferometry.py script, it must give precedence to the scheduled VLBI scans which need to occur at the appointed times. As mentioned above in Section 8.11.1, the VEX file will have provided gaps of various durations which are created to allow the observing script to make these ALMA-specific calibration observations. Thus, the script organizes its work by first noting when the next VLBI scan will occur and then noting which calibrations might fit into the time until the VLBI scan and finally executing those that fit. The VLBI scans in the VEX file include a start time specified in so-called VEX time format, YYYYyDOYdHHhMMmSSs, so the execution of the StandardVLBI.py script by the SSR will look something similar to:

```
perform initial ALMA calibrations
VLBI Scan 2016y089d07h00m00s on Target-X for 300s
perform some ALMA calibration
VLBI Scan 2016y089d07h10m00s on Target-X for 300s
perform some ALMA calibration
VLBI Scan 2016y089d07h20m00s on Target-X for 300s
perform some ALMA calibration
VLBI Scan 2016y089d07h30m00s on Target-X for 300s
perform some ALMA longer calibrations
VLBI Scan 2016y089d07h50m00s on Target-X for 300s
perform some ALMA calibration
VLBI Scan 2016y089d08h00m00s on Target-X for 300s
... perform final ALMA calibrations
```

VLBI Scan Sequence

The VEX file specifies the observations in terms of VLBI scans which begin and end at some appointed times. The VOM translates this request into a sequence of (ALMA) correlator scans which typically start prior to the specified time in order to allow the phasing system to stabilize on a good solution. Each of these correlator scans appears in the ALMA meta-data, and is processed by the telescope calibration system (TelCal) according to the specified intents. The relevant ones for this discussion are intents for phasing and for WVR correction. The phasing intents direct TelCal to calculate the phases of the signals (i.e., the signals of each polarization in each band of each antenna) relative to those of a designated “reference” antenna. These phases are calculated from “channel-averaged” data products calculated within the correlator. Presently there are eight channel averages calculated during one of these correlator subscans; see Fig 8.7. With those phases known, commands are issued to the tunable filter banks (TFBs) within the station cards in the correlator to adjust the antenna phases by exactly the computed values so as to bring all signals into phase with those of the reference antenna. Once these adjustments have been made, a “residual” phase correction typically needs to be made on the following correlator scan, and indeed, such corrections continue to be made every scan until the end of the recording. The phasing loop is closed in the sense that small errors in the phase adjustments may be corrected on subsequent scans. This is the so-called “slow” loop which is illustrated schematically in Figure 8.7.
Figure 8.7: VLBI Scans in the Phasing System. Each VLBI scan is partitioned into “subscans” for correlation and “slow” timescale processing (seconds) in TelCal. WVR adjustments are made in the CDP on a “fast” timescale (every second). See text for full discussion and explanation.

While the correlator scans are shown as time-contiguous in Figure 8.7 (green bar), there are in fact small gaps to allow the data to be dumped out for subsequent correlator processing. Even so, the solutions from TelCal (purple) do not arrive at the TFBs (red) until after the correlator scan has started—the time of application is noted and the early part of the scan is excluded from the subsequent phase calculation. Usually one phase application is needed to get an acceptable solution, so the recording should start no sooner than the start of the third correlator subscan. These parameters are all programmable, and judicious choices for them are made when the Scheduling Block is created. For Cycle 6, the subscans are likely to be 16 s long with 4 s (channel average) integrations; there are then about 30 s of observations prior to the start of the recording. Note that the ALMA correlator does its processing by subscans—each subscan must be dumped and processed by the correlator and ultimately archived—with a gap of several seconds between subscans. On the other hand, the signals from the antennas flow continuously to the summation logic and recorders, so there are no comparable gaps in the VLBI data.

In addition to the “slow” loop, there is also a “fast” loop that may be enabled and which uses the WVR data from each antenna to correct the antenna phases in response to the wet component of the atmosphere. These corrections are performed at a frequency of \(\sim 1\) Hz (blue in Figure 8.7). This loop is open (corrections are continuously made, but feedback is only achieved through the “slow” loop). Note also that all WVR corrections are made online in the ALMA correlator, but obviously after the WVR data was collected, so the phasing corrections are at their best in stable atmospheres, and deteriorate with increasing atmospheric instability. There is no limit to the duration on the operation of these loops (other than general ALMA ones with regard to total data collected).

**VLBI Phasing System**

Moving beyond the timing of the phasing system, there are several other aspects of the system that an ALMA VLBI observer should be familiar with. All ALMA observations operate with an array of antennas (i.e., those available to a scheduled observation). Of these, a portion are controlled by the phasing system to form the coherent sum signal which is used for the VLBI recordings. The previously mentioned reference antenna is merely a designated member of this “phased array”. At least two antennas from the active array (designated “comparison antennas”) are held outside the phased array. These comparison antennas are thus available for diagnosing the performance and efficiency of the phased array system. Finally, after construction the summed signal is decimated to two bits per sample (as is the case with all of the ALMA signals in the VOM), so it...
is possible to have the ALMA correlator correlate the sum signal during online processing. This is done by co-opting some antenna that is not used by the VOM and placing the sum signals into the correlator as if it had come from this “sum antenna”. The sum signal therefore appears as antenna “APP001” in any ASDM file containing ALMA data acquired with the VOM. There are a few things to note about the data and metadata for this antenna:

- it has no useful metadata, as it is not a physical antenna
- it has no WVR data, for the same reason
- it is highly correlated with other antennas in the phased array since the non-physical receiver noise is present in both.

The selections of antennas for the phased array and comparison array are made when the Scheduling Block is created, although these may be adjusted if necessary during the observation. Normally, the antennas to be phased are all chosen to lie within a certain radius from the array center. The efficiency of the phasing system decreases on longer baselines, so in practice, maximum baseline lengths in the phased array will typically be \( \leq 1 \) km. The reference antenna is in general selected to be centrally located, and typically at least one of the comparison antennas is chosen to be within the radius of the phased array. An additional detail is that because of the decimation to two-bits, the number of phased antennas must always be odd.

A final important note concerns delays between antennas. The ALMA control system estimates the total delays between all antennas and adjusts the signals through a variety of techniques (in hardware and software) so that the visibilities found in the dataset have essentially zero delay on each baseline, i.e. there is no slope in a plot of visibility phase as a function of frequency on any baseline. Unfortunately, a (significant) component of the delay is removed with a frequency dependent phase rotation in the online correlation processing, not in the hardware, so there is some significant delay present in the signals at the TFBs where the phase corrections are made and correspondingly in the logic which creates the “sum antenna” signal. Thus in order for the phasing system to work properly, it is necessary to turn off this component of software delay correction and to instead correct it in the phasing calculations. To do this, the observing band is subdivided into a relatively large number of channel averages (at least 8) and the normal phase-solving procedure produces different, independent phases that can be applied to the TFBs corresponding to each channel average. This does remove most of the delay, but it does leave a small residual of phase-slope within each channel average. The only restriction is that the source needs to be bright enough in each of the channels so that a usable phase solution can be found. The flux limit of 500 mJy in Cycle 6 ensures that this is the case.

**VLBI Observing Mode Analysis**

As pointed out in Section 8.11, any VOM observation actually involves two concurrent interferometric observations: one on the ALMA scale that results in correlated data products between all participating ALMA antennas and one on a global scale, where ALMA serves as a station in a global VLBI array. A description of the reduction and analysis of ALMA VLBI data is beyond the scope of this document\(^{17}\). This process is, however, informed by the analysis of the ALMA-scale observation with regard to observatory performance (SEFD, system temperature, etc.), which for the most part may be obtained in the usual way. However, since the VLBI signal that is recorded is that of the “sum” antenna whose properties in turn depend on the phasing performance, a few details should be pointed out. The most significant point is that the effective area of the “ALMA dish” is roughly proportional to the square root of the number of phased antennas. (In Cycle 6, only 12 m dishes are to be used, and they have similar effective areas. If the 7-m dishes were to be included, the relation is more complicated.) In practice, a number of effects conspire to lower performance from this ideal relation\(^{18}\) to about 60\% of what would be expected in a perfect system (see Matthews et al. 2018, PASP, in press for discussion). Part of this efficiency loss stems from the fact that the sum antenna signal is decimated to two bits from the individual antenna signals, which were in turn also decimated to two bits. Thus there is an additional reduction

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\(^{17}\)A VLBI analysis handbook is in preparation.

\(^{18}\)ALMA-05.11.63.03-0001-A-REP.pdf
in sensitivity by 0.88 that must be taken into account in the VLBI analysis. In any case, the net effective area of the ALMA “sum” dish may be calculated from the ASDM on a per-scan basis.

The performance of the phasing system is captured in the metadata in ASDM CalAppPhase tables. There is one table entry per scan per baseband. In each table, there are temporal bounds on validity (i.e., the range of time within the scan of stable phase) and entries to list the disposition of the antennas amongst the different categories (phased antennas, reference antenna, and comparison antennas) and whether this represents a change from the previous scan. The table primarily reports a number of phase values, Nv, which is dependent on the “packing mode”. In Cycle 6, there is one value per channel average per polarization per antenna. Additionally, the CalAppPhase table reports lower-level details of the phasing system, including online per-antenna assessments from TelCal of the antennas drawn from the quality of the phasing solution. This table may be used in conjunction with the ALMA data for a particular scan to calculate the phasing efficiency and thus to make an appropriate correction for the absolute VLBI correlation amplitudes.

As indicated in Section 8.11.1, the correlation products are delivered to the PI for analysis. Traditionally for the GMVA this has been in the form of FITS-ID1 files. A complication for work with ALMA data is that the polarization conversion is performed with a software tool (PolConvert) that requires input from the CASA analysis of the ALMA data set. This conversion can be performed on the FITS-ID1 file or at the correlator using the raw correlator output if the the CASA analysis is available. Arrangements for doing this will be worked out between the PI of each project and the supporting observatory collaborations (GMVA and EHT) following the observation.

8.12 High Frequency Observing

High-Frequency Observations (specifically Bands 8, 9, and 10) require some special observing techniques, mostly because of the combined ill effects of lower atmospheric transmission and fainter and more widely spaced calibrators at these frequencies. These techniques, which are just addenda to standard observing modes, are being progressively implemented in the ALMA observations for Cycles 6 and beyond, after some extensive testing is carried out. It is expected that all of them will improve not only the quality of the final data products, but also the overall observing efficiency at the ALMA highest frequencies. For reference, a brief description of the nomenclature that could be used by Contact Scientists and the ALMA Observatory in general when matters regarding high-frequency projects are discussed, reported or communicated to PIs is given below:

- **High-Frequency Cone-Searches and Calibrator Surveys**: Observations in the high-frequency ALMA bands require that the availability of calibrators near the target source is assessed near the time of the planned observations. Searches for phase calibrators and check sources (see Chapter 10 for details) within 10 degrees of the science targets (i.e., the so-called “Cone-Searches”) in each SB are carried out by ALMA staff and Allegro, the EU ARC Node in the Netherlands, within 90 days of the high-frequency observations. The results of those measurements are subsequently ingested into the ALMA Calibrator Catalogue, so that they are available for queries at the time the actual observations take place.

- **Bandwidth Switching (BWSW)**: For high-frequency observations (Bands 8, 9, 10) with spectral set-ups including only narrow-band spectral windows, the PI will be asked to add an additional wide-band spectral window. This approach increases the accuracy of the gain calibration, allowing a solution across the combined spectral windows and/or transferring of the wide-window solutions to the narrow spectral windows. When adding extra spectral windows is not possible (i.e., projects that have occupied all the spectral windows with narrow-band set-ups), it is possible to use bandwidth switching, in which the phase calibrator is observed with wider spectral windows for increased signal-to-noise. The wide-to-narrow spectral window complex-gain offsets are measured from observations of the bandpass calibrator, allowing the phase cal solutions to be transferred to the target field. For more details, please read Chapter 10.

- **Sideband separation by 90-degree Walsh phase-switching (Bands 9 and 10)**: The Band 9 and 10 receivers are of DSB design, i.e. do not separate or suppress either sideband, and signals from both sidebands are superposed on a single IF output per polarization (see Chapter on Receivers). Prior to Cycle 6 the only option available for interferometric observing in these bands has been to use orthogonal LO
frequency-switching sequences to suppress the contributions of one sideband to the visibilities. Since Cycle 6 a means of sideband separation is also offered, resulting in twice the output bandwidth and number of channels. The separation uses orthogonal sequences of 90-degree phase shifts applied at the first LO in each antenna, combined with synchronised accumulation over the sequence duration of the three phase combination states of the cross-correlations of each baseline by the correlators. The result is subsequently combined to produce the two sideband spectra in the correlator data processing software. There are 128 patterns in a switching sequence, each pattern switching on 16 ms intervals, and therefore total length of 2048 ms. This mode therefore adds a constraint: the dump and integration durations of the correlator must be a multiple of 2048 ms. Due to an unrelated requirement that subscans must start and end on 48 ms Timing Event (TE) boundaries, and 2048 ms not being a multiple of 48 ms, the allowed values of subscan durations are somewhat limited, as will be seen when adjusting times in the OT. In Cycle 6, only a single pair of spectral windows (LSB+USB) in each baseband is offered, and either the TDM (2 GHz, 256/Npol channels) or widest bandwidth FDM (1.875 GHz, 7680/Npol channels) modes must be used (additionally for FDM, the spectral window should not be moved from the center of the baseband). The modes in the four basebands remain independent, e.g. one baseband can use FDM to observe a spectral line, while the other basebands can use TDM to detect continuum.
Chapter 9

ALMA Sensitivity Calculator

The main tool for calculating the sensitivity of ALMA is the ALMA Sensitivity Calculator (ASC). This is an application contained within the ALMA Observing Tool (OT) that is used to perform various conversions between sensitivity (in either flux or temperature units) and time. Its main use is to take the user-requested sensitivity and calculate the necessary observing time for each source in a Science Goal using parameters such as the representative frequency, source declination, number of antennas and the expected precipitable water vapour (PWV). Many of the inputs cannot be changed by the user and, for example, the OT will always assume a fixed number of antennas for the particular cycle and will always use the standard PWV octile that is appropriate to the frequency of observation.

In addition, a separate version of the ASC is available in the OT which allows a user (via a GUI) to experiment with various sensitivity options such as the number of antennas and PWV octile. An almost identical version is also available as a web application that can be found in the ALMA Science Portal\(^1\). This differs from the GUI version incorporated in the OT in that it is possible to enter some parameters that are outside the ranges that are applicable to the current observing cycle.

9.1 Calculating the System Temperature

When determining the time required to achieve a particular sensitivity, the system temperature, $T_{\text{sys}}$, is a fundamental parameter as it takes into account various sources of noise that make it difficult to detect the very weak astronomical signals that ALMA is trying to detect. The most prominent sources of noise are from the receivers and from the atmosphere. The latter is highly variable, both in time and frequency, and thus dynamic scheduling and careful placement of spectral windows are crucial.

9.1.1 Sky Temperature

The OT’s estimate of both the atmospheric zenith opacity, $\tau_0$, and the sky temperature, $T_{\text{sky}}$, are calculated using the Atmospheric Transmission at Microwaves (ATM) code\(^2\). This provides values of the opacity and the atmospheric “output radiance”, in steps of 100 MHz, for the seven different octiles of PWV. The sky temperature is converted from the radiance using the Planck function and includes the contribution due to the CMB.

The ATM code only provides measurements of the sky temperature at the zenith, $T_{\text{sky}}(z = 0)$, and therefore the OT must account for the greater atmospheric emission at lower elevations. It does this by assuming that

\(^1\)http://almascience.org/

\(^2\)See Pardo, J. R., Cernicharo, J., Serabyn, E., 2001, ITAP, 49, 1683. This calculates the sky temperature by integrating the atmospheric temperature profile, this having been formed from the average of 28 radiosonde measurements taken at the ALMA site during November 1999.
Table 9.1: Octiles of PWV measured at the ALMA site from years of monitoring data and used in the ASC. The first octile corresponds to the best weather conditions and shows that 12.5% of the time, PWV values at least as good as 0.472 mm can be expected. Subsequent octiles give the corresponding value for 25%, 37.5%, etc.

<table>
<thead>
<tr>
<th>Octile</th>
<th>PWV (mm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>0.472</td>
</tr>
<tr>
<td>2</td>
<td>0.658</td>
</tr>
<tr>
<td>3</td>
<td>0.913</td>
</tr>
<tr>
<td>4</td>
<td>1.262</td>
</tr>
<tr>
<td>5</td>
<td>1.796</td>
</tr>
<tr>
<td>6</td>
<td>2.748</td>
</tr>
<tr>
<td>7</td>
<td>5.186</td>
</tr>
</tbody>
</table>

The octiles characterize the amount of PWV that can be expected at the ALMA site, i.e. a value of PWV at least as good as the first octile value can be expected 12.5 per cent of the time, a value at least as good as the second octile 25 per cent of the time, and so on. The octiles corresponding to the ALMA site (determined from many years of monitoring) are shown in Table 9.1.

When estimating the time for a project, the OT will always select a PWV octile that is appropriate to the frequency being observed. It does this by calculating the time required for each octile and then choosing (and reporting) the highest (worst) octile for which the increase in time relative to the first is less than 100 per cent. A consequence of this definition is that the octile also depends on source declination, i.e. sources at low elevations will require better weather conditions. The resulting curve of octile versus frequency is shown in Figure 9.1, for a source declination of zero degrees. A user can override this choice in the GUI or web application versions, but submitted projects will always use an automatic choice.

### 9.1.2 CMB Temperature

The temperature of the Cosmic Microwave Background is included in $T_{\text{sky}}$ and thus does not feature further in this document.
9.1. CALCULATING THE SYSTEM TEMPERATURE

Figure 9.1: Plot of PWV octile assumed by the ASC as a function of frequency, for a source declination of zero degrees. The vertical lines separate different bands, the numbers of which are shown at the top of the plot. The water line at 183 GHz (Band 5) is particularly prominent. In general, higher frequencies require drier observing conditions.

9.1.3 Receiver Temperatures

For ALMA Bands 4 and 10, the calculator currently only uses the specifications for the receiver temperatures (required over 80% of the receiver bandwidth) and not the actual measured values. For ALMA Bands 3, 5, 6, 7, 8 and 9, however, typical values measured in the laboratory are used as these are usually significantly better than the specifications. The measured values are somewhat conservative and so are in between what might be expected at the middle and edges of the bands. The values used in the ASC are given in Table 9.1.3.

Note that single sideband noise temperatures are reported for Bands 1-8 and double sideband temperatures for Bands 9 and 10.

At the moment, no attempt is made to incorporate the frequency dependence of the receiver temperature, i.e. only a single value is used per band. Ultimately, it is the intention to use the actual measured values for all receivers and to incorporate the frequency response across the band.

Note that the calculator doesn’t concern itself with the so-called “zero-point fluctuations” as the requisite half photon of noise ($h\nu/2k$) has already been included in the noise measurements provided by the various receiver groups (A. Kerr, private communication).

The receiver temperatures are already expressed in terms of the Planck expression and thus do not require the correction given in Equation 9.4.
Table 9.2: Receiver temperatures assumed in the ASC as a function of ALMA band. In general, these adopted values are conservative numbers towards the upper 10% of the actual measured $T_{rx}$ averaged over the whole band and over all antennas. For many bands, these are considerably better than the original specifications in Table 4.1.

### 9.1.4 Ambient Temperature

Ambient temperature is essentially spillover from the sidelobes of the antenna beam corresponding to emission from the ground and the telescope itself. This is held constant at 270 K (median value as measured from many years of monitoring data at the ALMA site). However, the value used by the ASC is converted to a noise temperature according to Equation 9.4 and thus its total contribution is frequency dependent and can vary between the different sidebands.

### 9.1.5 DSB Receivers

Due to the way that the astronomical signals are down-converted to an intermediate frequency, every heterodyne radio receiver simultaneously detects radiation from two sidebands. This means that, if nothing were done to prevent it, a spectral window processed by the correlator would contain two sets of astronomical signals mixed in with one another, one from the “signal” sideband and an undesirable one from the “image” sideband. In the case of SSB receivers, the contribution from the image sideband (emission from the source and noise) is suppressed to a very high degree. However, for DSB receivers (Bands 9 and 10) it is only possible to remove the source contribution (either by LO offsetting or 90-degree Walsh switching) and so the noise cannot be neglected.

One important consequence of this is that, if a spectral window has its image counterpart in an area of very poor atmospheric transmission, it can greatly increase the system temperature and lead to very long on-source times. Therefore, it is important to avoid areas of bad atmospheric opacity in the image spectral windows and the OT therefore shows the location of these in the spectral visual editor (Figure 9.2).

One subtlety is that the tuning software will always place spectral windows in the upper sideband if possible. Therefore, to take a simple case, a single spectral window centred at 637 GHz will find its image equivalent in the middle of the zero transmission feature at ~621 GHz. Where possible, this situation can be avoided by defining dummy spectral windows such that the line of interest is forced into the other sideband (Figure 9.2).

### 9.1.6 System Temperature

The OT version uses two distinct formulas depending on whether a double sideband receiver is being used, or not. The DSB equation is the following:

$$T_{sys,dsb} = \frac{1}{\eta_{eff} c^{-\gamma_0 sec^2}} \left( 2 \times T_{rx} + \eta_{eff} \left( T_{sky,s} + T_{sky,i} \right) + \left( 1 - \eta_{eff} \right) \times \left( T_{amb,s} + T_{amb,i} \right) \right)$$

(9.5)
9.1. Calculating the System Temperature

Figure 9.2: A Band 9 spectral setup as displayed in the spectral visual editor in the OT. The top figure shows an example of how the image equivalent of a single spectral window can fall into an area of poor atmospheric transmission, leading to much higher $T_{\text{sys}}$ than necessary. Placing a dummy spectral window at a higher frequency can remedy the situation (bottom).

where

- $T_{\text{rx}}$ – receiver temperature
- $T_{\text{sky},s}$ – sky temperature at the requested frequency in the signal sideband
- $T_{\text{sky},i}$ – sky temperature in the image sideband
- $T_{\text{amb},s}$ – ambient temperature in the signal sideband
- $T_{\text{amb},i}$ – ambient temperature in the image sideband
- $\eta_{\text{eff}}$ – the coupling factor, or forward efficiency. This is equal to the fraction of the antenna power pattern that is contained within the main beam and is currently fixed at 0.95
- $e^{-\tau_0 \sec z}$ – the fractional transmission of the atmosphere, where $\tau_0$ is equal to the zenith atmospheric opacity and $\sec z$ is the airmass at transit (the ASC always assumes that the source is being observed at transit).

The terms $\eta_{\text{eff}}$ and $e^{-\tau_0 \sec z}$ both attenuate the source signal and thus we must divide through by them in order to obtain a measure of the system noise that is relative to the unattenuated source. Note that this is always done at the signal frequency.

For SSB and 2SB receivers, the equation is the following:

$$T_{\text{sys,ndsb}} = \frac{1}{\eta_{\text{eff}} e^{-\tau_0 \sec z}} \left( T_{\text{rx}} + \eta_{\text{eff}} T_{\text{sky},s} + (1 - \eta_{\text{eff}}) \times T_{\text{amb},s} \right)$$

(9.6)

where the terms are all the same as in Equation 9.5.
9.1.7 System Temperature (web application and OT GUI)

The equations to calculate the system temperature used by the web application and the OT’s GUI are similar to the above, however, both are unaware of the details of the tuning setup and so do not know the location of the image sideband and cannot perform a rigorous calculation of its contribution to the system temperature. In this case, the same equation is used for all receiver types and the DSB noise contribution is simply double the single sideband case. This is controlled via the sideband gain ratio, \( g \):

\[
T_{\text{sys}} = \frac{1 + g}{\eta_{\text{eff}}} e^{-\tau_{\text{sec}} z} \left( T_{\text{rx}} + \eta_{\text{eff}} T_{\text{sky,s}} + (1 - \eta_{\text{eff}}) \times T_{\text{amb,s}} \right).
\]

(9.7)

For Bands 1 and 2 (Single Sideband; SSB) and 3-8 (Sideband Separating; 2SB), \( g = 0 \). For DSB receivers, \( g = 1 \).

9.2 The Sensitivity Calculation

Once \( T_{\text{sys}} \) has been determined it is possible to calculate the point-source sensitivity given a requested amount of on-source observing time or vice versa. At no point is any account made for the expected level of loss in sensitivity due to problems with the telescopes such as residual pointing and focus error.

Under certain circumstances it will not be possible to achieve the theoretical sensitivity, particularly if the source being observed is very bright and/or the \((u,v)\) coverage relatively sparse. In these situations the image can be dynamic-range limited due to residual source flux scattered across the map as a result of calibration errors and imperfect source deconvolution.

9.2.1 12-m and 7-m Arrays

When dealing with the 12-m and 7-m Arrays, the point-source sensitivity, \( \sigma_S \), is given by the standard equation:

\[
\sigma_S = \frac{w_r 2 k T_{\text{sys}}}{\eta_q \eta_c A_{\text{eff}} (1 - f_s) \sqrt{N(N - 1)} n_p \Delta \nu t_{\text{int}}}.
\]

(9.8)

The various parameters are:

- \( w_r \) – robust weighting factor. Pipeline imaging and subsequent QA2 assessment is performed assuming that the visibilities are weighted using robust weighting, specifically a Briggs robustness factor of 0.5. Simulations have shown that this factor is equal to 1.1.
- \( A_{\text{eff}} \) – effective area. This is equal to the geometrical area of the antenna multiplied by the aperture efficiency \( \eta_{\text{ap}} = R_0 \exp (-16 \pi^2 \sigma^2 / \lambda^2) \) where \( \sigma \) is the rms surface accuracy of the antenna. The latter is set to the design goal of 25 \( \mu \)m and 20 \( \mu \)m for the 12 and 7 m antennas respectively\(^3\). \( R_0 \) is the product of a number of different efficiencies and is equal to 0.72. See Table 9.3 for values of antenna efficiencies in various ALMA bands.
- \( f_s \) – shadowing fraction. For the more compact 12 m configurations and the ACA 7-m Array, antennas can block the field-of-view of other antennas in the array and thus reduce the total collecting area. The shadowing fraction is a function of source declination as shown in Fig. 7.5.
- \( \eta_q \) – quantization efficiency. A fundamental limit on the achievable sensitivity is set by the initial 3-bit digitization of the baseband signals. This is equal to 0.96.
- \( \eta_c \) – correlator efficiency. This depends on the correlator (64-input or ACA) and correlator mode, although the efficiency of all 64-input Correlator modes is equal to 0.88. The ACA efficiencies do depend on the mode, but this is only taken into account in the OT, not by the web application or OT GUI, which therefore assume a value of 0.88.

\(^3\)Note that not all antennas might achieve this specification. The performance of a given antenna will also vary with the thermal conditions and the length of time between surface realignments.
### Table 9.3: Aperture efficiencies at typical continuum frequencies for both the 12- and 7-m antennas. The effective area, $A_{\text{eff}}$, is equal to $\eta_{\text{ap}}$ multiplied by the physical area of the dish i.e. 113.1 m$^2$ and 38.5 m$^2$ for the 12 and 7 m antennas respectively.

<table>
<thead>
<tr>
<th>Band</th>
<th>Frequency (GHz)</th>
<th>$\eta_{\text{ap}, 12 \text{ m}}$ (%)</th>
<th>$\eta_{\text{ap}, 7 \text{ m}}$ (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>100</td>
<td>71</td>
<td>71</td>
</tr>
<tr>
<td>4</td>
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<td>5</td>
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<td>230</td>
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<td>7</td>
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<tr>
<td>8</td>
<td>405</td>
<td>60</td>
<td>64</td>
</tr>
<tr>
<td>9</td>
<td>690</td>
<td>43</td>
<td>52</td>
</tr>
<tr>
<td>10</td>
<td>870</td>
<td>31</td>
<td>42</td>
</tr>
</tbody>
</table>

- $N$ - number of antennas. This defaults to 43 for the 12-m Array and 10 for the 7-m Array.
- $n_p$ - number of polarizations. $n_p = 1$ for single polarization and $n_p = 2$ for dual and full polarization observations.
- $\Delta \nu$ - resolution element width. As already mentioned, this should be equal to 7.5 GHz for continuum observations. This is due to the maximum usable bandwidth of a spectral window being limited to 1.875 GHz by the anti-aliasing filter through which the baseband signal passes. $n_p \Delta \nu$ is often referred to as the effective bandwidth.
- $t_{\text{int}}$ - total on-source integration time.

The associated surface brightness sensitivity (K) is related to the point-source sensitivity (Jy) by

$$\sigma_T = \frac{\sigma_S \lambda^2}{2k \Omega}$$

(9.9)

where $\Omega$ is the beam solid angle. This is related to the user-entered spatial resolution, $\theta$, by

$$\Omega = \frac{\pi \theta^2}{4 \ln 2}.$$  

(9.10)

This assumes that the telescope beam is a circular Gaussian with a half power beamwidth of $\theta$.

### 9.2.2 Total Power Array

In the case of the TP Array, a different equation is used

$$\sigma_{TP} = \frac{2k T_{\text{sys}}}{\eta_0 \eta_k A_{\text{eff}} \sqrt{N} n_p \Delta \nu t_{\text{int}}}$$

(9.11)

This is similar to Equation 9.8, but there is only a factor of $\sqrt{N}$ in the denominator and there is no need to take shadowing into account.

Particularly for continuum observations, the above equation is likely to be too optimistic due to rapid fluctuations of the receiver gain and atmospheric opacity. These require extremely demanding calibration strategies and, as these have not yet been commissioned, only spectral line total power projects are currently possible.
Figure 9.3: Screenshot of the GUI version of the ALMA Sensitivity Calculator as implemented in the ALMA Observing Tool. The white area at the bottom is for displaying error messages i.e. parameters out of bounds. The example here shows the achievable sensitivity for all three arrays for an on-source time of 10 minutes.

9.3 User Interface

The main way that a user interacts with the Calculator is through a GUI in the OT or via the web application in the ALMA Science Portal – both are essentially identical. By entering various parameters, the time required to achieve a particular sensitivity (in either Jy or K) can be calculated, or vice versa. The inputs that affect the sensitivity or time are given below; a screenshot of the OT’s GUI version is shown in Figure 9.3.

- Source declination – this is used to calculate the maximum elevation of the observation and thus the minimum airmass i.e. the ASC assumes that the source is transiting. It is also used to calculate the amount of shadowing that is likely to affect the ACA 7 m and the smaller 12 m configurations.
- Observing frequency – this sets the receiver temperature, antenna efficiency and the PWV octile.
- Bandwidth per polarization – this otherwise straightforward parameter should be set to 7.5 GHz for continuum observations (see Section 9.2.1). For spectral line observations, it is usually set to the frequency/velocity resolution that one requires in one’s spectrum.
- Water column density (PWV). The user is able to enter one of the seven octile values, or the calculator will set this automatically depending on the frequency entered.
- Number of antennas – the ASC currently assumes 43 from the 12-m Array, 10 from the 7-m Array and 3 from the TP Array.
Angular resolution – this affects the time estimates when sensitivities are specified in temperature units. The calculator will not perform any calculations when Kelvins have been specified, unless a non-zero value for angular resolution has been entered. The calculator will also issue a warning if the angular resolution falls outside of the range corresponding to 125 m and 1 km baselines.

The calculator reports the values of $T_{\text{rx}}$, $\tau_0$, $T_{\text{sky}}$ (including the correction for the source elevation) and $T_{\text{sys}}$ that correspond to the entered frequency and PWV.

9.4 Total Time Estimates

Note that the time calculated by the ALMA Sensitivity Calculator does not account for telescope overheads (calibration, software and hardware latencies, etc.) and is only used to calculate the expected on-source time. However, the time estimates calculated by the OT do include these various sources of overhead and Chapter 5 of the OT User Manual should be consulted for details on how the total time required to observe and calibrate an ALMA project is calculated. The web application and OT GUI will only provide on-source times.
Chapter 10

Calibration and Calibration Strategies

10.1 Introduction

This chapter describes the methods used by ALMA in order to calibrate and edit data for imaging and for scientific analyses. First, a review of aperture synthesis, visibility functions and calibration assumptions is given. Then, the organization and descriptions of all ALMA calibrations are separated into four parts: Section 10.2 for initial calibrations and improvement of antenna properties; Section 10.3 for antenna integration/calibrations after an antenna move or major configuration changes; Section 10.4 for Execution Block (EB) Calibrations; Section 10.5 for additional calibrations that the PI could consider, including self-calibration. Short sub-sections 10.5.8 and 10.5.9 contain information on where information about ALMA total power and solar calibrations and imaging can be obtained.

Since about 2014 an increasing number of EB executions have been calibrated using the ALMA pipeline and its processing steps¹ are referenced in this chapter. For non-standard PI observations and many test observations, a python script is run manually². The pipeline reduction methods evolved from the manual script so they are similar.

Section 10.4 and Section 10.5 may be of most interest to PI’s to check their ALMA calibrated data and to improve the calibrations. But, first we review the steps and assumptions as applied to ALMA data.

10.1.1 Complex Visibility Function

An ALMA observation is made in one of (currently) eight frequency bands, at specified frequency resolutions, with an array in a given configuration capable of producing images at different resolutions, with up to four polarization states. ALMA contains 54 antennas of 12 m diameter and 12 antennas of 7 m diameter and their general properties and uses are described in Chapter 2.

The fundamental theory of aperture synthesis is given in Chapter 3 and the correlation of the signals between any antenna-pair is described in Chapter 4. These correlation quantities are called the complex visibility function, \( V(u, v) \), where \((u, v)\) is the projected east-west and north-south separation of the two antennas as viewed from the source direction. Its name is often shortened to visibility.

It is assumed that the emission from celestial objects is random and incoherent; that is, there is no correlation between the emission between two different times or between two different points in the sky. This lack of coherence permits the simple integration of the intensities (sky and visibility) with no interference terms that would be associated with coherent emission processes.

²https://github.com/keflavich/SgrB2_ALMA_3mm_Mosaic/blob/master/script_7m\uid___A002_X85b7b_Xb3.ms.scriptForCalibration.py
The angular distribution of emission from the sky can be described in a coordinate system of directional cosines $I(l, m)$. The van Cittert-Zernike theorem (Section 3.4) shows that the emission distribution is related to the complex visibility functions by a two-dimensional Fourier transform. If the visibility functions are sampled in the $(u, v)$ plane sufficiently well, an estimate of the sky distribution can be reconstructed with a resolution that depends on the maximum length of $(u, v)$. In practice, the visibility functions are sampled and averaged over a short period of time (a few seconds), with each antenna pair sample corresponding to a $(u, v)$ separation. In this case, the approximate emission distribution is then given by the sum of these visibility functions over the $k$ measurements as,

$$A(l, m)I(l, m) = \sum_k V(u_k, v_k)w_k e^{-2J\pi\nu(u_k l + v_k m)} \quad (10.1)$$

where $J = \sqrt{-1}$. The term $w_k$ is a visibility function weight that is related to the density of the sampling in the $(u, v)$ plane. See Chapter 7 for interpretation and use of the visibility weighting.

### 10.1.2 ALMA Visibility Data Flow

All PI observations are made from a scheduling block (SB) that contains the observation schedule needed to obtain an image of the target (see Section 8.1). All relevant calibrations are included in the SB so the data can be completely calibrated. An execution block (EB) is the use of an SB for one observation that rarely exceeds two hours. During the EB execution an assessment of the data quality using electronic monitoring results and some initial calibrations is made. This assessment is called QA0 and is described in Chapter 11.2. If the data quality passes or provisionally passes, the data are stored in the ALMA archive.

Many PI projects need multiple EB executions of an SB at one configuration in order to meet the rms or uv-coverage specifications of the experiment. These executions need not be observed consecutively. However, polarization experiments require an SB to be run two or more time in succession, and these linked EB’s are called a session.

The combination of data from several configurations including the 7-m Array and the TP array may be needed to meet the science goals of a proposal, and the collection of these data may cover more than one year. This chapter, however, concentrates on the calibration for a single EB, but the last section will discuss methods and suggestions for combining many EB’s to make an image.

### 10.1.3 True and Measured Visibility Data

Before the measured set of visibility functions can be used to determine the intensity distribution in the sky, the data set must be calibrated. This is a catch-all term that is used to correct for the changes of the radio signals on their path from the radio source through the atmosphere to each antenna; the changes as they pass through the complicated electronics and long signal paths of ALMA; then multiplied pair-wise in the correlator where the visibility function from each antenna pair are formed and then stored in the archive.

The measured complex visibility for any single frequency stream at any time $t$ and frequency $\nu$, associated with antennas $i$ and $j$ can be split into a visibility amplitude $a_{ij}^m(t, \nu)$ and visibility phase, $\phi_{ij}^m(t, \nu)$,

$$V_{ij}^m(t, \nu) = R(V_{ij}^m(t, \nu)) + JI(V_{ij}^m(t, \nu)) = a_{ij}^m(t, \nu) + J\phi_{ij}^m(t, \nu) \quad (10.2)$$

$R$ and $I$ signify the real and imaginary parts with $J = \sqrt{-1}$. The archive also contains information about each antenna and other array parameters needed to interpret the set of visibility data. The true visibility function in the same notation is

$$V_{ij}^o(t, \nu) = a_{ij}^o(t, \nu) + J\phi_{ij}^o(t, \nu) \quad (10.3)$$

Notice that the dependence of the visibility on the $(u, v)$ coordinate has been replaced by the antenna pair $(i, j)$ associated with the visibility data. The $(u, v)$ coordinates can be determined from the time, frequency, and location of the antennas.
10.1.4 Calibration Simplification

The simplest conversion of the observed visibility function to the true visibility function is made with a complex baseline calibration gain $G_{i,j}(t, \nu)$. It is a function of antenna pair $(i, j)$, time $(t)$, and frequency $(\nu)$,

$$V_{i,j}^o(t, \nu) = V_{i,j}^m(t, \nu) \ast G_{i,j}(t, \nu). \quad (10.4)$$

The time sampling depends on the averaging period of the visibility data and the time-scale for significant changes in the calibration parameters.

Although this calibration approach is valid, the number of complex gain calibration terms is proportional to number of the baselines which is about 1000 for a typical ALMA observation. However, there are three factorizations of the complex baseline calibration gain that are valid and generally utilized:

**Antenna-based**: Nearly all of the changes that corrupt the visibility function (atmosphere, system noise, amplitude changes, etc) can be decomposed into the two complex antenna-based gain factors associated with a baseline. This reduces the number of gain correction terms for an $N$-element array from $N(N-1)/2$ to $N$ with a considerable increase in signal to noise per antenna-based gain compared with the baseline-based gain. The antenna gain separation property is called closure$^3$.

**Amplitude and Phase-based**: The calibration of the antenna signals are nearly always associated with changes that can be separated into an amplitude and a phase correction, rather than independent real and imaginary corrections.

**Dispersive and non-Dispersive**: Many antenna-based phase changes are associated with effective path-length variations along the antenna signal routes (including the troposphere) which produce antenna-based phase changes that vary linearly with frequency. These changes are called non-dispersive. Their linear phase-frequency relationship means that these corrections need to be measured at only one frequency.

Using the above simplifications concerning the calibration properties of the ALMA (and most interferometric) system, the conversion of the measured visibility function into the true visibility function becomes:

$$V_{i,j}^o(t, \nu) = V_{i,j}^m(t, \nu) * A_i(t) * A_j(t) * B_i(\nu) * B_j(\nu) \exp(2i\pi[\alpha_i(t) - \alpha_j(t) + (\beta_i(\nu) - \beta_j(\nu))] + \alpha_i(t) \exp[2i\pi \delta_{i,j}(t, \nu)]$$

where $A_i(t)$ is the temporal amplitude variations of antenna $i$; $B_i(\nu)$ is the frequency amplitude variations of antenna $i$; $\alpha_i(t)$ is the temporal phase variations of antenna $i$; $\beta_i(\nu)$ is the frequency phase variations of antenna $i$. The non-closing part of the amplitude and phase is given in the last line and is negligible for most astronomical applications.

It is assumed that the sum of all antenna signal changes produced by many different effects are linear in their amplitude product or in their phase sum. For example, the antenna signal from an amplitude change of an amplifier and that from an atmosphere absorption change will combine appropriately in their summed correction factor.

10.1.5 Reference Antenna

The observed visibility phase data for any baseline, frequency and time is corrected by applying the difference in the antenna-based phase calibration of the two relevant antennas at the relevant frequency and time. It is clear that adding an arbitrary temporal phase to all antenna-based phases will not change the corrected baseline visibility phase. Hence, to bring specificity to the antenna-based phase corrections it is most convenient to

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$^3$see e.g. Taylor et al. 1999 Synthesis Imaging in Radio Astronomy, ASP Conference Series, Vol. 180, 1999
choose one antenna as the reference antenna which is defined to have zero phase at all times for all relevant frequency, time and polarization. This then determines unambiguous antenna-based calibration phases over the EB.

Since the antenna phases are used for many types of reduction and analysis algorithms, choosing a reasonably phase-stable reference antenna will make such analysis easier and more robust since the non-reference antenna phases will be more continuous. The pipeline task `hif_refant` picks a set of possible reference antennas by using several criteria, such as the amount of data flagged, and distance from the array center. It is strongly recommended that the same reference antenna is used for all data streams and calibration steps associated with an EB.

10.1.6 Calibrator Sources

Some antenna calibrations can be accurately predicted from algorithms associated with the environmental or electronic metrics associated with ALMA. After these corrections are made, additional calibrations are necessary, and these can be estimated from short observations using test sources of emission with known and measured radiative properties. Nearly all additional calibrations are made with interferometric observations, even if parameters are needed only for one antenna, because the correlated signal between two antennas removes much of the single antenna-based noise that interfere with many single antenna calibrations.

The absolute flux density scale for many experiments is derived from an observation of selected solar system objects that have accurate emission models and are not too large in angular size, Section 10.4.8. Every approximately 15 days, there are short ALMA observations of a solar system object with 40 brightest compact sources, called `grid` sources to determine their variable flux densities at 100, 250 and 350 GHz, and the current list is shown Table 10.1. Because these sources have a straight power-law spectra, the extrapolation of the measured 100 and 350 GHz flux densities to higher frequencies can be estimated to about 10% accuracy. They are used for determining the bandpass amplitude and phase shape for all observations.

<table>
<thead>
<tr>
<th>J0006−0623</th>
<th>J0237+2848</th>
<th>J0238+1636</th>
<th>J0319+4130</th>
<th>J0334−4008</th>
</tr>
</thead>
<tbody>
<tr>
<td>J0423−0120</td>
<td>J0510+1800</td>
<td>J0519−4546L</td>
<td>J0522−3627L</td>
<td>J0538−4405</td>
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<td>J1159−2914</td>
</tr>
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<td>J2253+1608</td>
<td>J2258−2758</td>
<td>J2357−5311</td>
</tr>
</tbody>
</table>

L = Large scale emission; R = Resolved source

Table 10.1: Bright Quasars Used For Bandpass and Flux Density Calibrators

For phase referencing (see Section 10.4.9) calibrators, the ALMA catalog\(^4\) contains a list of about 5000 small-diameter sources (mostly quasars) with accurate positions obtained from VLBI catalogs\(^5\). This large list is needed to find a relatively close calibrator and check source to any ALMA target.

10.2 Initial Calibrations

These calibrations are made as soon as possible after incorporation of an antenna into the array or after major overhauls. They are executed by the ALMA staff.

\(^4\)https://almascience.eso.org/sc/
\(^5\)Petrov VLBI catalog: astrogeo.org/vlb/solutions/rfc_2017c
10.2. INITIAL CALIBRATIONS

10.2.1 Antenna Surface Setting and Efficiencies

The surface accuracy of an antenna is obtained using interferometric holographic techniques by measuring the relative amplitude and phase of a bright quasar as the antenna pointings are offset from the quasar position\(^6\). At higher ALMA frequencies, experiments are made at several elevations. After correcting the surface panel deviations from the desired antenna parabolic shape, the nominal antenna efficiencies at frequencies less than 300 GHz are about 70%, and decrease at 690 GHz to about 43%, 52%, respectively for 12 m and 7 m antennas (see Table 9.3). The minimum loss of efficiency is 30%, and is caused by the aperture blocking of the feed and subreflector and by small residual panel offsets. At the higher frequencies, the rms surface accuracy of 20 \(\mu\)m for the 12 m antennas and 15 \(\mu\)m for the 7 m antennas cause a further decrease in efficiency. These efficiencies are only achieved during the night when thermal antenna deformations are a minimum.

10.2.2 Antenna Focus

The emission reflected from the main antenna surface is directed toward the feed for each band by using a subreflector that can be tilted and moved. The appropriate focus curves (subreflector offsets and tilt) are generated and stored for each antenna, and remain fixed for weeks. Additional information is given in Appendix A and elsewhere\(^7\). However, daytime observations may contain a short focus-check observation of a bright source at the beginning of the EB. The FocusModel.xml and Focus.xml files attached to the ASDM data set contain the parameters used for each scan file.

10.2.3 Antenna Primary Beam Pattern

To first order, the relative response of the antenna sensitivity as a function of angular offset from the pointing center varies with the antenna diameter and frequency. Assuming that the antenna surface and focus are accurately determined, the field of view, \(\theta\) (defined as the angular width between half-power sensitivity) in radians is inversely proportional to the number of wavelengths (\(\lambda\) in the antenna diameter \(D\)).

\[
\theta = 1.13\lambda/D
\]  

(10.6)

More details on the shape and characteristics of the primary beam are given in Baars(2003)\(^8\), and in Chapter 3. The beam pattern shape is important for constructing accurate mosaic images.

10.2.4 Antenna Pointing

The antenna pointing correction is the difference between the true sky location of the peak sensitivity of the antenna primary beam and the location shown on the azimuth and elevation encoders. After an antenna has been assembled or had mechanical work performed, interferometric observations of about 40 quasars around the sky are made with this antenna, plus several others. The measured all-sky pointing offsets are then fit to 19 parameters that describe the expected azimuth and elevation dependencies. These interferometric observations are called all-sky pointing, and are usually made at nighttime. Each quasar scan of about 25 sec consists of five sub-scans; at the nominal source location; and at the four nominal half power points. The TPOINT analysis program\(^9\) is used, with details of the ALMA implementation are given by Mangum & Lucas(2007)\(^10\). The pointing fits have an accuracy of about 1.5" over the sky, and are generally made at 230 GHz.

The pointing difference between the ALMA bands are azimuth offset and an elevation offset above the general fit, and these offsets are accurate to about 1.5". They are determined from a short pointing observation of a strong quasar at the relevant band. The pointing parameters are nearly constant over days and weeks as long

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\(^6\) Mangum et al, PASP.118,1257,2006
\(^7\) https://arxiv.org/pdf/1210.1899.pdf
\(^8\) http://legacy.nrao.edu/alma/memos/html-memos/alma456/memo456.pdf, (ALMA Memo 456)
\(^9\) http://www.tpointsw.uk/index.htm
\(^10\) https://safe.nrao.edu/wiki/pub/ALMA/CalExamples/PointingCalStepByStep.pdf
as the antenna remains at the same pad. However, some small pointing updates may be needed for daytime observations and these are discussed below.

10.3 Antenna Integration Calibrations

There are presently 10 ALMA configurations of the 12 m antennas with maximum baselines from 160 m to 16.2 km. At least 43 antennas are needed for PI observation in any configuration. There is a configuration change every 6 weeks for which about 15 antennas are moved over a period of about 10 days. Sometimes the hybrid arrays can be used for PI science projects. After an antenna is moved, calibrations are needed to be integrated into the array, and these major calibrations are listed below.

10.3.1 Improved Antenna Pointing

Additional pointing observations are needed after an antenna has been moved to a new pad. Five of the 19 pointing parameters are expected to change: the azimuth and elevation offset, two antenna tilt terms and one overall rotation term. This relevant observation is part of the antenna integration process and uses a short all-sky-pointing observation, described in Section 10.2.4. This short observation will provide a nominal pointing accuracy over the sky of about 2''.

10.3.2 Antenna Delay

The signal paths from an antenna to the correlator will differ considerably among antennas. If these path length differences are not removed, then the correlated visibility between the relevant baselines will show large phase slopes. For example, at an observing frequency of 100 GHz, if the residual baseline path length offset is 10 m, then the resultant phase slope will be (10 m)/(3 mm) = 300 rad. Thus, in a 2 GHz spectral window, there will be a phase slope of (2 GHz)/(100 GHz) * 300 rad = 6 rad or about one revolution. To remove these large and smaller residual delay offsets, an observation of a strong source can measure a phase slope over 2 GHz to a 40'' accuracy, corresponding to a residual delay offset 1.5 cm will lead to a phase slope of 40'' over a 2-GHz bandwidth. An even more accurate phase slope is determined when the bandpass shape observation associated with each EB (See Section 10.4.7).

10.3.3 Antenna Positions

In order to make high quality images, the relative antenna locations should be known to 0.2 mm accuracy in order to limit systematic phase errors between the calibrator and the target that will produce distortions in the target image. However, the estimated antenna position after placed on a new pad may be in error by up to 5 mm, so its position must be more accurately determined. Unlike antenna pointing and focus offsets that not only decrease the effective antenna sensitivity but distort the derived image, an antenna position offset affects only the visibility phase which can be corrected by the offline calibration with no deleterious effect on the image quality. Hence, the accurate determination of antenna positions for a configuration need not be determined before the initial PI observation are made for a configuration, as long as the improved antenna positions are measured before the offline reductions are made.

The accurate antenna positions are determined from a sequence of many short observations of bright quasars that are observed over the sky. Using a script called all-sky-delay, each quasar is observed for 20 sec and about 40 sec are needed to move to the next quasar that is usually be more than 40'' away in order to sample the entire visible sky in about 20 min. The observations are made at Band 3 where the quasars are strongest, in TDM mode with a total bandwidth of 8 GHz in the four spectral windows. Often, the entire ALMA array is used to improve the SNR of the solutions, even though only 10 to 20 antenna positions need be improved. For a typical observing time of about 45 min, the accurate antenna-based delay for about 50 scans are obtained for each antenna and scan at the < 0.001 psec accuracy. If the assumed antenna position is accurate, then the
residual delay for all scans will be constant, independent of quasar elevation and azimuth. An analysis script in
the reduction tool TelCal, called tc_antpos is used to fit the set of the residual delays to obtain the estimated
position offset and error estimate for each antenna. An example of the fit for one antenna of the scan delays
before and after the correction of the fitted residual antenna offsets are shown in Figure 10.1.

Figure 10.1: Antenna Position Determination: The antenna position fit for DA46, at 3.5 km from the
reference antenna, from a 50-min all-sky-delay observation on Dec 25, 2017. Each of the 45 points show the
delay solution from one scan, separated by about 1 min. The observations were made at Band 3 where a delay
change of 0.001 ns is equal to a phase change of 40°. The red points show the original delay measurements
and the scatter is caused by antenna position offsets as a function of quasar elevation and azimuth. The black
points show the delay after correction of the DA46 antenna position error by about 4 mm. After correction, the
spread of phases between sources over the sky is about only 20°, so the expected phase difference from a residual
antenna position offset between a target and calibrator that are just a few degrees apart will be considerably
less.

As part of the antenna integration procedure, a shorter all-sky-delay run of a small number of antennas is
now made in order to obtain the antenna position to an accuracy of 1 to 2 mm. This does not replace the full
all-sky-delay run, but decreases the position errors so that test data reductions and PI reductions will not be
significantly impaired by errors larger than 2 mm.

After the all-sky delay run, the updated antenna positions be placed in the on-line data base (TCMBD)
that is used by the delay server. The updated antenna positions are also currently placed in a file which can be
accessed by the pipeline and other offline reduction systems from the Analysis Utilities depository of useful
ALMA offline scripts. The antenna position uncertainties depend on the distance from the array center (the
reference antenna is usually chosen near the array center), because the atmospheric phase variations, caused by
large-scale turbulence, generally increase with baseline length. The typical antenna position errors are about
0.05 mm per km of baseline length. Hence, at the longest baselines, a typical antenna position could be as large
as 1 mm.

10.4 Execution Block Calibrations

This section describes how each ALMA PI execution block are calibrated. A pipeline calibration system is used
for standard ALMA executions11. For non-standard observations (currently solar, polarization, astrometry,

band-to-band and bandwidth switching) a similar but flexible reduction methods, called scriptForClibration.py, is used manually. This was the early ALMA reduction script that evolved into the ALMA Pipeline.

10.4.1 Making the Measurement Set

The archived data from the array are stored only at the ALMA site and at JAO, Joint ALMA Operation Center in Santiago, Chile. These data are transformed into an ASDM (ALMA Science Data Model) data structure\textsuperscript{12} using asdmExport. This format is then converted into a measurement set (ms) using the CASA task importasdm. These steps will not be needed by most PI's, but could be useful for special tests and calibrations before official reductions are made. The pipeline task hifa_importdata makes this ASDM to ms conversion.

10.4.2 Flagging Data Based on Calibrations

A significant part of the flagging of data is determined from derived calibration values that are faulty in some way. This discrimination is made relatively simple for ALMA data since most calibration values should be similar among more than forty antennas that are used for most PI observations. There are several methods used in the pipeline to determine reasonable flags, such as hifa_flagdata (which includes antenna shadow flagging), hifa_fluxcalflag and hif_rawflagchans. See the pipeline manual for implementation details.

10.4.3 Water Vapor Measurements

Variations from water vapor (PWV) in the antenna line-of-sight cause significant delay variations, up to 0.5 mm/sec, corresponding to 50\textdegree\ of phase/sec at 90 GHz. The changes are driven by wind turbulence and cloud structure, and cover a large range of spatial scales from 5 m to many km. All 12 m ALMA antennas are equipped with a Water Vapor Radiometer (Dicke-type) that measures the emission from atmospheric water vapor every 1.1 sec at four frequencies near 183 GHz, in nearly the same direction as the ALMA pointing. The radiometers were developed by B. Nikolic\textsuperscript{13} and information about their design and use is given in Section A.6.

These WVR corrections are made by the online calibration software system in TelCal\textsuperscript{14}. An offline script using the CASA task wvrgcal also converts the water vapor emission to a predicted antenna delay, and the algorithms of the two corrections are similar. In the near future only the uncorrected visibility data will be stored in the archive and the offline WVR correction will be applied as part of the calibration process. This enables small edits or scaling of the WVR emission to delay to be made if needed\textsuperscript{15}. The on-line corrected data, however, will be available for use of the QA0 (Chapter 11) analysis of each experiment, but not stored in the archive.

The use of the WVR correction reduces by more than 50\% the water vapor-induced phase fluctuations on time scales of 2 sec to many minutes, enabling more experiments to meet the minimum success level, and to produce significantly higher quality images. An example of the antenna-based phases at Band 9 for an antenna before and after WVR correction is shown in Figure 10.2. The PWV was about 0.45mm for this example, and it is the typical improvement for relatively low WVR content. The generation of the WVR delay correction is given in the reduction script mystep = 4. The pipeline tasks hifa_wvrgcal and hifa_wvrgcalflag determine the WVR correction and associated flags from WVRgcal value that are out of range or noisy.

10.4.4 System Temperature Measurements

Another important supporting calibration of ALMA determines the antenna receiver sensitivity and the radiative sky temperature. These are obtained from a 15-sec scan where a hot load and ambient load are placed in front of the antenna. The load scans are made in a 15-sec scan where a hot load and ambient load are placed in front of the antenna.

\textsuperscript{12}http://adsbit.harvard.edu/full/2006ASPC..351..627V
\textsuperscript{13}2013, A&A, 552, 104
\textsuperscript{14}http://adsabs.harvard.edu/abs/2011ASPC..442..277B
\textsuperscript{15}L. Maud, https://arxiv.org/abs/1707.03506
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Figure 10.2: **WVR correction**: The improved temporal phase by using the WVR correction is shown in the four panels for four different antennas at Band 9. The red lines show the temporal phase before correction; the green lines after correction. The rms phase fluctuations before and after correction are (in rms phase in degrees): DV16-40 to 26 deg; DV60-51 to 38 deg; DA56-36 to 26 deg; DV12-31 to 23 deg. The offsets from zero phase are arbitrary.

of the feed. Both 12 m and 7 m antennas are equipped with these Amplitude Calibration Devices (ACD), and their description and use are given in Section A.5.1.

At frequencies below about 400 GHz, where the system temperatures are relatively constant over an hour of time or 10$^\circ$ in the sky, a ADC scan is made every ten to twenty minutes. However, at higher frequencies all scans should have an associated ADC measurement. The online measurements of $T_{\text{sys}}$ and $T_{\text{rx}}$ are stored in the CalAtmosphere table in the ASDM datasets, and are displayed in real-time at the telescope to show possible measurement problems immediately for QA0 assessment. An example of the $T_{\text{sys}}$ measurements from a Band 9 observation is shown in Fig. 10.3. The two antennas show the typical range of the $T_{\text{sys}}$ measurements at this high frequency. Notice that the $T_{\text{sys}}$ in each spectral window contains narrow lines that are associated with tropospheric emission. For example, the additional line emission of 20% in the first spw near channel 60 is the emission from the Hydrogen recombination 673.91 GHz line.

The $T_{\text{sys}}$ measurements are applied to the correlated data in the reduction script `mystep=4`. The pipeline task `hifa_tsyscal` determines the $T_{\text{sys}}$ correction and `hifa_tsysflag` produces flagged antenna data when $T_{\text{sys}}$ values are out of range. Most of the atmospheric line emission will be removed in the $T_{\text{sys}}$ corrected spectrum, but there will be increased noise level at these line locations.

10.4.5 Antenna Pointing Tweaking

Additional pointing observations are sometimes included in an EB (see Section 8.2). For observations > 350 GHz, the pointing between sources more than about 20$^\circ$ apart may differ by more than 2" from the nominal pointing
Figure 10.3: System Temperature at Band 9: The measured system temperatures at Band 9 for a target source for two antennas with different properties. Top row is for DV05 for four spws, bottom row is for DV09 for four spws. The ordinate is the system temperature, the axis is the channel number in each spw. The average frequency for each spw is 673.9, 675.9, 677.9, 679.9 GHz, and the channel separation is 1.56 MHz. The two polarization T_sys differs significantly for DV05. The finite width of each T_sys curve is caused by the small time variation over the 50-min EB.

solution. Hence, a pointing observation with the bandpass/flux density calibrator scan (see Section 10.4.7) and with the phase calibrator may be included. Also, for observations that demand the highest pointing accuracy, say < 0.5", a pointing every 10 min on a source near the phase calibration or target can be done (pointing referencing).

10.4.6 Antenna Position Tweaking

The determination of the most accurate antenna positions after a configuration change is described in Section 10.2.4. Because a PI observation may be executed before the most accurate antenna position are measured, the visibility phases should be corrected for the most recently measured antenna position offset. The corrections are determined from the comparison of the on-line antenna positions in the data base to the updated position in an offline file\textsuperscript{16}. The reduction script mystep = 6 reads the new antenna positions, compares them with that used during the observations and determines the phase offset to be applied. The pipeline task hifa_antpos calculates the corrections.

10.4.7 Bandpass/Frequency Calibration

The goal of the bandpass calibration is the determination of the relative amplitude and phase of the frequency channels in each data stream. As discussed in Section 10.1.4, because of the excellent stability of the ALMA bandpass (Kameno 2015)\textsuperscript{17}, the relative frequency dependence is constant to within a few percent and one degree over an hour-long EB; hence, a bandpass calibrator is observed once at the beginning of an EB, with a

\textsuperscript{16} AIV/science/PadData/almaAntPos.txt
\textsuperscript{17} https://almascience.nrao.edu/documents-and-tools/documents-and-tools/alma-technical-notes/AlmaTechnicalNotes15_FINAL.pdf
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Figure 10.4: BP Phaseup: The phase variation for a Band 9 bandpass calibrator. (Top) Xpol and Ypol phase and their difference for 675 GHz spw, (Middle) Xpol and Ypol phase and their difference for 679 GHz spw, (Bottom) 679 GHz - 675 GHz spw for Xpol only and their difference. The WVR correlation has been applied to all data. The phase variations over time are extremely well correlated between polarization and spws, and these phase offsets are virtually constant over an EB 2-min to 30-min scan, depending on the requested bandpass accuracy. In some cases, the phase calibrator (see below) that is observed many times over the entire EB can be used, if necessary, as the bandpass calibrator.

A bright quasar is generally chosen from the list of grid sources (see Table 10.1) that is closest to the target, and it is called the bandpass calibrator. The choice of quasar will also depend on the observing band, channel frequency width and the bandpass accuracy that is needed for the science goal. The nominal accuracy goal is 2% per relevant antenna channel-width.

There are two steps in obtaining the bandpass calibration. First, the temporal phase variations over the bandpass scan are determined. An example of the phase changes over the bandpass scan is shown in Fig. 10.4. These time-dependent phase fluctuations are contaminated by tropospheric delay changes that are polarization independent and scale with the small frequency difference between spectral windows. This similarity can be used to combined frequency channels for subsequent temporal phase calibrations, described in Section 10.4.9.

After removal of the bandpass scan phase changes, the second step determines the bandpass amplitude and phase dependence for all channels in each spw/polarization. The bandpass solutions are usually normalized to average amplitude 1.0 and phase 0.0 for each spw to produce a relative bandpass across each spw. The pipeline task hita_bandpass and manual script determine the bandpass. The bandpass solution for four antennas at Band 9 for four spw is shown in Fig. 10.5 and Fig. 10.6. The accuracy of each channel solutions is < 1% and < 1°. The bandpass response is flat over a 2-MHz width and up to this frequency width can be averaged to improve the S/N of the bandpass. The 200 MHz ripples in the figure are real and may have several causes related to the antenna structure reflections and very small reflections in some of the antenna paths in the system. Finally, near one of both edges of many spectral windows, there may be significant deviations from the nominal bandpass...
behavior. This part of the spectrum tend to have a much larger noise component and be variable over time, so they are generally flagged. The pipeline task for bandpass determination and flagging is `hifa_bandpassflag`. In this case channels 0-9 and 123-127 should be removed for all antennas/spw.

![Figure 10.5: Bandpass Amplitude](image)

The bandpass amplitude at Band 9 for four antennas, Xpol are shown as blue dots and Ypol as green dots. Xpol channels below 10 and Ypol channels above 120 should be be flagged. Notice that the line emission near channel 60, seen in the $T_{\text{sys}}$ solution (Fig. 10.3), has been significantly reduced after multiplication of the raw visibility amplitude by the $T_{\text{sys}}$ factor.

### 10.4.8 Absolute Flux Density Scale

The application of the system temperature correction, Section 10.4.4, converts the visibility amplitudes from correlator units to Kelvin (K). The conversion factor from K to Jy (SEFD: Source Equivalent Flux Density) can then be found by using an astronomical object with a known flux density model. For ALMA, the SEFD is stable with time and pointing direction, so the value found from the flux calibrator source is valid to a few percent (see last part of this section) and the flux density of other calibrators can be derived.

After the bandpass calibration is applied and the phase variations are removed for all calibrator sources in the EB, the following calibration steps are made using the pipeline to obtain the flux density of all calibrator sources.

**Solar System Object:** If the primary flux scale object is one of the six solar system objects that are used normally (Mars, Uranus, Neptune, Ceres, Pallas, Juno), the source model is obtained from CASA task `setjy` and the task `gaincal` will determine the SEFD ratio for each antenna and spectral window. For larger ALMA configurations and high frequencies, only the antennas near the array center, where the source model does not drop to less than 50% of maximum, should be used.

**Transfer SEFD:** The next step made in the task `fluxscale` that transfers the above measured SEFD to
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Figure 10.6: Bandpass Phases: The bandpass phase at Band 9 for four antennas, Xpol blue dots and Ypol green dots.

other calibrators, most important is the (grid) bandpass calibrator. Sometimes the transfer to the phase calibrator and check source to obtain their flux density has a noise bias value if they are weak.

Grid Source: If an appropriate solar system object is not available, then one of 40 (variable!) grid bandpass sources can be substituted as the amplitude calibrator. Their flux density and spectral index between 100 and 350 GHz are measured every 15 days (Section 10.1.6) so that their flux density can be estimated at any observing frequency between Bands 3 to Band 10. This SEFD is then transferred to phase calibrator, check source and target.

Narrow Bandwidth with Low Signal-to-Noise: If some of the spectral windows are too narrow to derive an accurate SEFD, then the SEFD derived from one wide band spectral window can be transferred to the frequencies of the narrow bands using the known spectral index of the solar system object or grid calibrator.

Flux Density Scale Accuracy: There are many variables in the determination of the flux density scale for any observation, with most effects becoming more severe with higher frequencies. The approximate absolute flux scale limits are 5% at Band 3, rising to 10% at Band 7, and then to 20% at Band 9 and 10.

The SEFD distributions derived from many PI, grid and test observations, now being compile, confirm that the expected SEFD stability for nearly all antennas and frequencies is stable to a few percent. The spread of SEFD for more than half of the antennas have an rms distribution (1-sigma) about 2% at Band 3 to 5% at Band 9 around an average value that is stable. Those antennas with SEFD out of the range generally have already known problems. These results suggest that the accurate flux density scale of an ALMA observations may be determined apriori from the SEFD value clustering for half of the array, especially below 400 GHz. At higher frequencies, the elevation dependence of the antenna gains and other effects must be incorporated, regardless of the method for determining the flux density scale of an observation.
10.4.9 Phase Referencing Calibration

A critical calibration for obtaining good quality images from an EB is called the phase referencing calibration. The calibration goal is the removal of temporal changes of the amplitude and phase in the target visibility measurements over the EB. The phase referencing technique alternates quickly between the target and the phase calibrator that are separated by less than a few degrees in the sky. The CASA task gaincal is then used to determine the amplitude factor needed to transfer the measured (or corrected so far) visibility amplitude to the known calibrator flux density, and the visibility phase offset needed to transfer the measured (or corrected so are) visibility phase to zero — for each antenna and spw and polarization. This antenna amplitude and phase is then applied using a simple interpolation between phase calibration scans, to the target.

The properties of a good quality phase calibrator are:

**Calibrator-target Separation**: The closer the calibrator is to the target, the more likely that the calibrator measurements will apply accurately to the target. This includes short-term phase variations of time scale of 10 sec and systematic phase differences of many minutes that are associated with large-scale atmospheric structure and antenna position offsets.

**Calibrator Detection**: The calibrator antenna-based solution from each scan and antenna and spw must be determined with sufficient sensitivity. For example, an $S/N = 10$ for an antenna solutions produces a 10% amplitude error and $5^\circ$ phase error. The SNR can be increased by a factor of 2.8 by summing the four spws and two polarizations (after obtaining their phase offsets derived from the bandpass observation (Section 10.4).

**Calibration Cycle Time**: The cycle time is defined as the time between subsequent calibrator scans. For successful phase referencing, the calibrator phase difference between adjacent cycles should be less than about 60°. The cycle times as a function of frequency and configuration are currently fixed to nominal values, so cannot be adjusted to the phase conditions at the time of observations.

For configurations longer than 5 km, the current cycle time is 80 sec, with the phase calibrator scan length of 18 sec. For smaller configurations, the cycle times increase to 240 sec with the phase calibrator scan length about 30 sec. These cycle times may be updated in the future. At frequencies above 400 GHz, the cycle times can be made longer in order to have sufficient SNR on the phase calibrator. For ALMA configurations less than 2km and frequencies below 400 GHz, the best calibrator (closest to the target and above a detection limit of 15 per antenna solution) is chosen automatically from the ALMA catalog. For longer baselines and higher frequency observations, a candidate list of calibrators is observed in short observation sessions called a cone-search that (1) determines if the calibrator flux density at the high frequency is above the detection limit and (2) determines if the calibrator not extended. Cone-searches are also used for selection of the check source.

For any EB with maximum spacing greater than 5 km, and/or at frequencies greater than 400 GHz, check source observations are included in an EB. The criteria for choosing a check source from the ALMA catalog are: (1) about the same distance from the phase calibration as the phase calibrator is from the target and (2) strong enough to produce an image $SNR > 15$. This flux limit is about a factor of three less than that for a phase calibrator. Most of the catalog sources are unresolved even at the longest ALMA baselines, but exceptions do occur. Although there are no firm rules yet, a check source is observed every fourth cycle with a length the same as the phase calibrator. At least two or possibly three check source scans should be included in each EB unless it is shorter than about 30 min.

The CASA task gaincal determines the antenna-based amplitude and phases from the phase calibrator for an EB. The pipeline task hifa_timegaincal is used. An example for two antennas from a Band 9 observations for a 1-km configuration is shown Fig. 10.7. The X- and Y-polarized phases follow each other since the variations are mostly caused by tropospheric delay changes over each antenna. A good check on the signal-to-noise of the calibration scan and other problems is given by the X-Y phase difference. It should be relatively flat with variations that are consistent with the expected noise. From the amplitude and phase solutions, periods of unstable gain and phase conditions can be easily recognized with the appropriate data flagging made.

Data should be flagged based on the temporal calibration anomalies. The cause of the bad data and its duration should be investigated but often requires a long investigation that is not possible. A conservative
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Figure 10.7: Calibrator Phases: The calibrator temporal phase at Band 9 for two antennas p(top) 1.0 km from the reference antenna and (bottom) 0.3 km from the reference antenna. Over the 4000-sec EB, 10 phase calibrator scans were made, each of length 60-sec with three 20-sec points connected to show the internal scan phase variation. The X and Y polarizations show virtually the same temporal changes, as expected. Within each scan the phase change is about $20^\circ$, between scans separated by 7 min the difference is about $50^\circ$ and is a little larger for the longer baseline. The X-Y phase difference (green (points/line) shows the approximate SNR of each solution, and is much less than the atmospheric variations. The offset is arbitrary for clarity.

scheme is to flag all data for the offending antenna between the time of the preceding good calibrator value to the time of the next good calibrator value. Please report to the ARC any serious flagging conditions that were missed by the ALMA reductions.

The antenna-based calibration amplitude (based on the estimate of the phase calibrator flux density discussed in Section 10.4.8 and phases) are then interpolated to all of the other scans using the task applycal. Since the amplitude variations are generally much smaller and smoother in time than that of the phase, the interpolations are often made independently for amplitude and phase. For the phase a simple two point (from calibrator scan to calibrator scan) linear interpolation is generally used, although more complicated interpolation and smoothing times are available. For the amplitude correction, a solution over many calibrator scans are often made (after applying the phase corrections), because the amplitude stability of ALMA over tens of minutes is often a few percent or less.

Finally, an image of the phase calibrator is provided as part of the calibration package. See hif_makeimlist and hifa_makeimages. It should be a good quality point source with a peak intensity and integrated flux density that agree to $2\%$ of the flux density assumed in the calibrations. The residual sidelobes or rms noise should be $<1\%$ of its peak intensity. If the phase calibrator image does not meet these expectations, then the calibration and editing process is suspect.

10.4.10 Check Source

The analysis of the check source image defects and position offsets are a good indication of what to expected for the target image, even though the integration time on the check source is considerably less. The scaling of the check source problems will depend mainly on the separation ratio of the check source to phase calibrator compared with that of the target to the phase calibrator. For example, if the target is half the distance from the calibrator as that of the check source, then to first order the decoherence of the target and its position offset should be about half found from the check source.

Check source images are not routinely included in the pipeline or script PI packages. Thus, it is recommended to make a check source image with the same parameters as that made for the phase calibrator and target in order to derive information about the EB quality from the check source.
Since the check sources are point sources of known position (about 5% of the them may have some extended structure), their cleaned image is a good indication of the quality expected for target image quality. If the PI data package does not already contain a check source image, it should be made using the imaging parameters as that for the phase calibrator and target. The characteristics of the check source image to be noted are:

What is the ratio of the peak flux density to the integrated flux density, an approximate measure of the EB de-coherence? A rough guideline is: above 0.9 is good, below 0.5 is not good.

What is the position offset of the check source from its assumed position? This is an indication of the positional uncertainty of the phase referencing. An offset up to 0.2 times the resolution is normal, but an offset larger than > 0.5 resolution is suspect. Such an offset could be caused by an inaccurately known check source position or structure, or caused by a low elevation (< 45°) observation.

Is the Check Source Image Distorted? If the check source image is more than symmetrically broadened, with clear substructure, then the phase referencing quality of the experiment is suspect.

Can the Check source be used as a secondary Calibrator? Sometimes the check source is nearly as strong as the phase calibrator and can be used as a secondary phase calibrator. The target image made from only the check source as the calibrator will provide information on the quality of the phase referencing. For astrometric observations, the check source should be observed as often as the phase calibrator. See Section 10.5.2

10.4.11 Polarization Calibration

The feed at the focus of each ALMA band receives two orthogonal linear polarization streams, called X (n/s direction) and Y (e/w direction) polarizations. The four resultant correlation pairs are designated as XX, XY, YX, YY. The initial calibration procedures are the same as that for the normal calibration of the XX,YY polarizations described above. However, the additional calibrations using all four of the polarization streams to obtain the total, linear and circular polarization emission of the target require more reduction steps that are discussed in Section 8.7. These data must be reduced with the manual reduction script.

10.5 Additional Calibrations

This section contains suggestions for additional analyses and calibrations beyond that normally done with standard ALMA reductions. Depending on the science goals, the observing conditions, and the complexity of the target image, further processing by the PI may improve the image quality. Several common techniques are described, but additional information and support can be obtained from your local ARC.

10.5.1 Self-Calibration

For each EB, the pipeline or manual script produce an image of the target. For about 25% of the targets, especially at configurations larger than about 2 km, the self-calibration algorithm may improve the quality of the image, sometimes by a factor of five or more as measured by the peak flux density to the rms noise in the image, and also increase the reliability of fainter image features. Hence, this is the first extra-calibration step that is recommended and described. This method is more robust than with other high frequency arrays because of the more than 40 antennas used for ALMA observations.

The self-calibration algorithm uses the CASA task gaincal much like that for the phase referencing of the target, Section 10.4.9. In this case, the phase calibrator model is known: usually a point source of measured flux density at a known position. Using gaincal, the antenna-based amplitude and phase corrections that are needed to make the phase calibrator visibility data more consistent with the model are determined. When these calibration corrections are applied, the resultant calibrator image is much high quality, and additional calibrator structure down to the few percent level that was not included in the original calibrator...
model, will appear in the new image. This demonstrates that the antenna-based calibration procedure to improve the image is not very sensitive to the input source model. This is because any additional calibrator structure, which changes the closure properties of the visibility data (see Section 10.1.4) will largely average out when using 43 antennas, and thus have little effect on the antenna-based corrections that produce an improved image. This structure versus antenna calibration isolation depends on the ratio of the number of baselines to the number of antennas, and is about 20 for an ALMA observation of 43 antennas. If the source structure is very complicated and extended, then this isolation is less strong even if the target is relatively strong.

The starting input image for the target self-calibration is usually the target image determined from the normal calibration and editing. Whether to use the continuum image or a spectral line image is related to which image has a larger peak/rms ratio. The success of the target self-calibration will also depend on the quality of the initial phase-referenced image, especially if the target is extended. Finally, the quality of the check source image, if available, is also a guide to whether self-calibration of the target will be useful. For example, if the check source image is almost point like with little loss of coherence, then the target image quality should be even better and self-calibration may not improve the image.

A recommendation is to try self-calibration on a target when the normally calibrated image shows emission with a peak of at least 20 times the rms noise. Use the CASA task gaincal with the target.model as input, and for this initial attempt use solint = 'inf' to obtain one antenna solution for each target scan. This time-scale of several minutes will not remove the image degradation from faster changing phase variations from the troposphere, but will remove longer term systematic phase errors that also decrease image fidelity. Use the same reference antenna as that used for the normal calibrations. For continuum self-calibration, use calmode='p', gaintype='G'. Combining spectral windows for this first self-calibration try is recommended combine='spw' all spws have virtually the same phase deviations.

From this attempt, the likelihood that self-calibration can improve the image quality can be assessed. The antenna-based phase solutions (using task plotcal) should show continuity between scans. Also, a measure of the each solution noise error is obtained with the plot of the (X-Y solution) poln = '/'. Since all true atmospheric phase variations will be the same at either polarization. For sources with significant extended emission, the self-calibration noise phase scatter will increase with the longer spacings because of the decrease of the target correlated flux density. But, even with an rms scatter of the X-Y phases of 15° per scan/antenna solution, the atmospheric phase variations are often larger, and this additional random phase component will not significantly reduce the self-calibration image quality.

If the results from this initial attempt suggests that self-calibration may be successful, then the continued reduction path will vary depending on many properties of the data and source. A few suggestions for the case of a moderately strong and not too extended target follow:

**Initial Phase-only self-calibration:** It will be clear from the above first execution if the solution interval can be shortened to (considerably) less than one target scan. If the scan solution interval produces noisy solutions, then self-calibration may not be useful.

**Next phase-only self-calibration:** Assuming good results from the above step, put in the improved target model, and phase self-calibration again, with a shorter solution interval if possible. After this phase self-calibration, deep cleaning, especially of an extended source, is recommended. Even cleaning a large region down to only 1.5-sigma will still leave a significant amount of residual flux density in the image. This will leave symmetric image feature in the residual image that are not real. If in doubt, keep cleaning the known emission regions.

**Later amp and phase self-calibration:** If you reach a 100:1 peak to residual image quality, another self-calibration iteration that includes amplitude can be useful. Put in the model image from the previous phase self-calibration, and then use a relatively long solint (about 5 to 10 min) since ALMA generally has excellent short-term amplitude stability. Use the option calmode='ap' rather than calmode='a' because this option gives better stability amplitude solutions even though the phase corrections are minimal.

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18 from a small number of channels, suitably averaged over velocity using the CASA task mstransform
19 For the initial image model to use, the natural weighting option will give better image SNR and less resolution which may produce a more accurate self-calibration first iteration.
In summary, the self-calibration of moderately strong sources that are not too extended is robust with ALMA data and should be attempted by the PI. The above suggestions are meant as a guideline to get started, but may require advice from your local ARC or from the JAO, especially if image peak/rms levels over 1000 are possible to achieve. In a few years, the ALMA reduction systems may include self-calibration techniques to improve image quality.

10.5.2 Astrometric Observations/Calibrations

The main goals of astrometric observations are to determine the accurate position, proper motion, and/or parallax of relatively compact objects. The \((u, v)\) coverage is not critical, but the most stable phase conditions and long baselines (as long as the object is not too resolved) are needed. The two major astrometric limits are 1) the SNR of the peak of the relatively compact object; 2) the systematic position error associated with normal phase referencing results. Obtaining higher precision involves using more than one calibrator. A discussion of ALMA astrometric results can be found in\(^{20}\) and are summarized below.

The astrometric limit precision of an object on an image of often limit by the SNR of the object. This limits, \(\Delta \rho\), in mas is approximately,

\[
\Delta \rho = 60 \text{mas} \times \left(\frac{100 \text{GHz}}{FREQ}\right) \times \left(\frac{10 \text{km}}{BslMax}\right) / \text{SNR}
\]  

(10.7)

where \(\text{SNR}\) is the intensity ratio of the peak to rms of the image, \(FREQ\) is the observing frequency in GHz and \(BslMax\) is the maximum baseline in km of the configuration. If two objects are in the same image, then the accuracy of their relative separation is given by the above thermal noise limitation if they are both within the \(\sim 70\%\) primary beam power sensitivity. If they are further apart, then their separation accuracy will decrease, and this will depend on the phase calibration stability. In order to eliminate image gridding errors, the image pixel size should be no larger than about 10% of the image resolution. If the field field contains one or a few compact objects that dominate the total emission in the field, then \((u, v)\) model fitting (CASA task \texttt{uvmodelfit}) may provide an accurate determination of the astrometric parameters.

The absolute ALMA astrometric limit is the precision of an object in an ALMA image with respect to the ICRF fundamental reference frame that is used for ALMA observations and all calibrators. The phase stability during the EB and the proximity of the phase calibrator to the target field are the major influences of the absolute image astrometric accuracy. Of course, the SNR limit of astrometric precision is also relevant. With a stable atmosphere and a calibrator-target separation less than a few degrees, an absolute astrometric limit of about 2% of the synthesized beam is possible, but this also requires an SNR=50 for the object. To reach this astrometric accuracy under less ideal conditions, cycling among several phase calibrator surrounding the target will improve the accuracy. These techniques is described in the above reference.

The year-to-year ALMA configuration cycling is not conducive for parallax determinations that require periodic observations over the year period. Proper motion accuracy are not critically sensitivity to specific configurations, so can be scheduled with some flexibility to obtain the relevant time coverage needed.

10.5.3 Calibrator Structure

Some calibrators (and check sources) are not perfect point sources and there use may effect the image quality. However, for those calibrators where > 80% of the emission in contained in an unresolved component, the effects on the antenna-based phase solutions are minimal, generally less than a few degrees. If these calibrators are used to improve amplitude calibration of the target, errors up to 5% and be made, and it is recommended to rely on the SEFD values obtained from the bandpass calibrator. An analysis of the imaging errors obtained from non-point sources will be available\(^{21}\). In the cases when a significantly resolved phase calibrator has been used (egs. a solar system object), the visibility data should be recalibrated using the a priori model or that obtained from the phase calibrator data in the experiment.

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\(^{21}\)Image Errors from Calibrators with Structure, H. Messias, in preparation
10.5.4 Band-to-Band Phase Calibration

For phase referencing at frequencies above 400 GHz, finding a suitably strong calibrator closer than 10° to a target is not common. As noted in Section 10.4.9, the closer a calibrator is to the target, the more likely that 1) the short-term (< 1 min) phase fluctuations will be similar between the phase calibrator and target, and 2) the longer term (2 to 60 min) phase differences between them will be decreased.

The goal of band-to-band phase transfer is to calibrate a target observation at high frequency by using phase referencing with a phase calibrator at a much lower frequency where there is a greater chance of finding a detectable calibrator nearer to the target. The success of this calibration requires that; 1) all phase changes among all of the spectral windows (high frequency and low frequency) are non-dispersive (delay-like) and scale with frequency, regardless of their time-scale and origin; 2) there is a constant instrumental phase offset for all of the spws that varies less than a few degrees over the EB period. These constraints are illustrated in Eqn.(10.8)

\[
\phi_h(t) = \phi_l(t) \ast \left( \nu_h / \nu_l \right) + \psi
\]

The instrumental phase difference \( \psi \) is determined by alternating observations of a bright source at the low and high frequencies, and is called the DifGainCal observation (DGC). Tests are in progress to determine the optimal frequency switching time in order to remove the short-term atmospheric phase changes between the two frequencies, to determine guidelines for the length of DGC observation to measure the \( \psi \) for each spw to an accuracy of a few degrees, to confirm that this spw-phase offset is constant with time up to two hours, and to confirm that the instrumental phase is independent of sky position. There are also tests to determine at what distance ratio between the nominal high frequency calibrator from the target compared with a closer low frequency calibrator to the target will produce significantly higher quality B2B target images.

10.5.5 Band-Width Switching

When observing a target within narrow spw bandwidths (typically < 250 MHz), a suitable phase calibrator, especially at the highest frequencies, may not be detectable within a one to two minute calibrator scan. Two methods of calibrating these narrow bandwidths are available:

**One wide spw:** If the aggregate bandwidth in the combined spws in the experiment has a bandwidth > 1 GHz (usually because one of the spws is wide band), then there will be sufficient sensitivity to determine the antenna-based gains.

**Band-Width-Switching:** If all of the spws are narrow bandwidth, it is likely that a reasonably nearby phase calibrator cannot be detected within a scan. In this case, the strategy for phase calibration is to observe the phase calibrator with wide-band spectral windows and the target in the desired narrow bandwidths by switching the spectral setup between the relevant scans. In a method similar to band-to-band calibration, a strong DGC calibrator is initially observed at both the wide and the narrow bandwidths in order to determine their phase difference among the spws that should be constant with time. There will be a slight delay change among the different spws, but they will be much closer in frequency than that for band-to-band observations.

Tests are now in progress to determine the stability of the wide-to-narrow bandwidth phase difference.

10.5.6 Multiple EB Combination

Many PI projects require several EB’s from one configuration in order to obtain sufficient SNR or \((u, v)\) coverage to meet the scientific goals. The guidelines are evolving, but usually the pipeline processing does not begin until all of the EB executions from one configuration have been completed. All EB’s are then individually calibrated, then concatenated into one measurement set, imaged, and then sent to the PI, and the calibrated data sent to
the archive. For weak targets (where the peak/rms intensity per EB is less than 10), the concatenation of the EB's are needed to obtain an image with reasonable SNR.

For those targets in which the SNR is greater than about 20 from each EB from an SB, it is useful to determine the effective quality associated of each EB before combining them to produce the target image. The easiest check of the EB to EB consistency and quality is to make an image of the target separately for each EB. Even with slight differences in their resolution and integration time, the images should be similar. The differences that might arise among the images are:

**Registration among the Images:** This is determined by comparing the position of the peak emission on each of the images, as described in Section 10.4.10. Offsets at the level of 0.2 times the resolution are typical because of difference of the phase conditions among the EB’s. Significantly larger offsets are usually associated with relatively poor phase stability or low elevation associated with the relevant EB.

**Quality among the Images:** Unless the images are dynamic range limited (usually when peak/rms intensity > 100), the rms noise in each EB should be similar, whereas the peak intensity can vary considerably depending on the average atmospheric phase variations over the EB. As a rough guide, for EB’s in which the peak flux density is less than 70% of the nominal peak value from the other EB’s, it is likely that the phase stability during the EB was poor and this EB should be down-weighted.

The default image presently made from a concatenated data set containing several EB’s uses the relative weight in each EB based on the average $T_{sys}$ and on-target integration time. This weighting does not include the phase stability that affects the image quality, but probably should. Guidelines for this additional weighting among EB’s are being discussed, and the PI should consider down-weighting some EB’s (or even ignoring them) that are of poor quality.

### 10.5.7 Combining Configurations

For targets that contain a wide range of angular structure, two different 12 m configurations and the possible addition of the 7 m configuration will be observed. A description of how the data are combined and weighted is based on the rms noise of the data points. The concatenation of all of the interferometric data, up to two 12 m configurations and the 7 m configuration into one measurement set, can be made using the CASA task *concat* and is the responsibility of the PI for subsequent imaging. The entire interferometric data set is then cleaned and examples are given described in Chapter (7.9). With data that cover a large range of angular size, it is likely that images at a variety of resolutions will be useful to display the relevant information for a complex source.

### 10.5.8 Total Power Observations

Total power (TP) observations are made using up to four of the ALMA 12 m $PM$ antennas for determining the structure of sources that are larger in extent than the 7 m antenna primary beam size. These data are usually associated with mosaic interferometric images. The reference for the calibration of these data can be found in, and the observing modes are discussed in Chapter 8 and the imaging in Chapter 7.

Good quality TP images require two important calibrations: First, an off position is needed to define the zero sky temperature in the vicinity of the target. Second, using a fast TP raster to map the target area, the system temperature difference between all the target raster points and the off-position is gives the target temperature distribution. Third, the conversion from Kelvin to Jansky can be made from the relatively accurately known SEFD of each of the PM antennas used for the TP, or by observing a solar system object with a known thermal temperature.

The TP image is combined with the cleaned image made from all of the interferometric data using a feathering method, CASA task *feather*. In most cases, these images are mosaics that have been obtained

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22https://casaguides.nrao.edu/index.php/DataWeightsAndCombination

23https://casaguides.nrao.edu/index.php/M100_Band3_SingleDish_4.5
from many interferometric pointings and a corresponding TP image area. The feather method uses a Fourier technique in which each image distribution is converted to the \((u, v)\) domain, and then combined. The two controls are a scaling factor between the TP and interferometric (FFT) image, and a low pass filter as needed (See Chapter 7). This total Fourier data is then transformed back into the image plane to produce the \textit{combined} image.

A second combination method is based on a software program called \textit{tp2vis}, and is being developed for future ALMA use\textsuperscript{24}. This method converts a total power image into appropriate \((u, v)\) data which is then concatenated to the much higher resolution interferometric data. This entire data set can then be cleaned to obtain an image that contains all angular scales. This combination approach gives more flexibility with the total power data and the interferometric data combination and weighting. Another advantage is that the interferometric data are cleaned after the TP image has been included. This avoids some of the large-scale sidelobes that are often associated with interferometric data that do not include the large-scale structure.

\subsection*{10.5.9 Solar Observations}

ALMA provides a unique facility for research of the solar chromosphere and the structure and dynamics of the many small-scale solar phenomena. The document\textsuperscript{25} provides the most up-to-date information on the present aspects of ALMA solar observations. These experiments are significantly different from other ALMA experiments, so contact your local solar-expert at the appropriate ALMA ARC for information about proposal submission, observations and support for calibration and imaging of the data. At the present time only Bands 3 and 6 are supported, and both the 12-m and 7-m Arrays are observed simultaneously, usually with total power using the PM antennas. Configuration 3 (maximum baseline of 0.5 km) is the largest configuration allowed at present. It is also the preferred configuration since the 7-m Array is well-displaced from the center of Configurations 1 and 2 that give non-circular \((u, v)\) coverage.

\textsuperscript{24}https://github.com/tp2vis/distribute
\textsuperscript{25}Shimono et al. 2017, Solar Physics, 292, 87S
Chapter 11

Quality Assurance

The goal of ALMA Quality Assurance (QA) is to ensure that a reliable final data product is delivered to the PI, that is, the product has reached the desired control parameters outlined in the science goals (or is as close to them as possible), it is calibrated to the desired accuracy, and calibration and imaging artifacts are mitigated as much as possible. The QA analysis will be based on a calibration plan that specifies which observations must be acquired and at which intervals to monitor system performance and environmental factors as they evolve with time. Furthermore, it will also be used in assessing the merging of data within each science goal taken with different configurations, the inclusion of 7-m Array and TP Array data, and the final image quality. Errors introduced by user supplied parameters, such as incorrect source coordinates, inadequate frequency setting (e.g., an incorrect redshift) or inadequate sensitivity limits (leading to an inadequate integration time or inadequate uv plane coverage) are outside the scope of the ALMA QA, unless the error occurred due to faulty information or tools provided by the Observatory.

To be more efficient in detecting problems, ALMA QA has been divided into several stages that mimic the main steps of the data flow. The broad classification of this multi-layered QA approach is:

**QA0:** Monitoring of calibrations and overall performance during and just after observations

**QA0+** A simplified fast pipeline imaging QA. This typically is completed 30-90 minutes after the observations are taken, and is mainly used when the results from QA0 are unclear or marginal

**QA2:** Full Calibration and generation of science products

**QA3:** Issues found with the data by the PI or the ALMA contact scientist after data delivery

The QA0, QA0+ and QA2 stages will be handled by the Program Management Group (PMG) and the Data Management Group (DMG) (with contributions from ARC personnel) using the ALMA Quality Assurance (AQUA) Tool (see Section 11.7). Responsibility for data quality assurance rests with the Data Manager and his Deputy, within the Department of Science Operations, drawing upon the resources of the Program Management Group and the Data Management Group. The final output of the ALMA QA0-QA2 process is a “QA Report” per ObsUnitSet (Member or Group\(^1\)) that summarizes all the relevant QA information for each of the different QA stages up to, and including, the final imaging. This report is included in the data package delivered to the PI. The QA3 stage will be handled separately, by the ARCs, via internal tickets created by the ARC personnel (see below). A more detailed description of the different stages of QA is given below.

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\(^1\)Group ObsUnitSet processing, that is, the combination of datasets taken in different configurations/arrays, will only be done sporadically during Cycle 6. It is expected that group processing will be implemented as a regular data reduction strategy in later ALMA cycles.
11.1 Cycle 6 Quality Assurance Goals

The ultimate goal for ALMA QA is that the delivered products are considered "Science Ready", suitable for publishing with little need for the user to reprocess. However, ALMA cannot guarantee that all the scheduled Cycle 6 projects are completed (especially those with grades B and C), or that specific criteria like sensitivity or angular resolution are precisely met, given the restrictions of the array configurations, system performance at the time of the observations, etc. The observatory will therefore attempt to meet the sensitivity and resolution stated by PIs in their Science Goals, within certain tolerances as described later in this chapter. In principle, each project component ("Scheduling Block" or SB) is scheduled for the number of executions that are expected to reach these goals (based on nominal Array performance values). Additional executions may be needed during the observing cycle if specific executions do not pass QA0 or QA2 (see below).

Observers may need to invest their own time and expertise to ensure that the data products are of the appropriate quality and to re-reduce the raw data if the quality is not satisfactory. This may include the need to visit the relevant ARC or ARC node to get help and to assist with quality assurance and potential data reduction.

11.2 QA0

The first stage of quality assurance, known as QA0, is a near-real-time verification of data quality for each SB execution, or execution block (EB). It deals with performance parameters on timescales of an SB execution length or shorter, and thus is performed at the time or immediately after each execution. Initial QA0 assessment is performed by AoDs (Astronomers on Duty) at the OSF, using the software tools AQUA, AosCheck, and QuickLook, based on semi-real time output from the calibration scans obtained by the real-time TelCal ALMA software. This information is complemented with reports derived using Monitor and Control display tools to monitor specific parameters not directly tracked by the calibrations (e.g., total power level variations, weather parameters, etc).

QA0 metrics/parameters have been selected to check the health of the whole signal path from the atmosphere down to the correlators, as well as issues connected with the observation itself. AosCheck goes through the calibration data, looking for outliers or out-of-spec results, and other problems. Together with additional checks in AQUA itself, these are combined to provide a QA0 assessment for each EB. These parameters can be grouped into the following general categories:

- **Atmospheric Effects**: Weather Parameters, Sky Opacity, System Temperature, Phase Fluctuations, WVR Outputs.
- **Antenna Issues**: Antenna Gains, Antenna/pad delays, Relative/Offset Pointing, Geometric Shadowing, Antenna/Pad positions.
- **Front-End Issues**: Bandpass, Sideband Ratios, Receiver Temperatures, Phase variations, variations in $T_{sys}$ spectral shape.
- **Connectivity Issues**: IF Delay Measurements, System Temperatures, unusual relative phase or amplitude variations between spws.
- **Correlator Issues**: Bandpass, Delay Measurements.
- **Observation Issues**: Calibrator fluxes, Incomplete datasets or incomplete mapping, mismatch between desired and actual synthesised beam shape.

Tolerances for individual parameters in the above list (for example the maximum range of $T_{rx}$) are defined in aoscheck. These are combined to assess the QA of individual sub-units of the EB, such as antennas, spws,

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Some additional hardware problems can occur which are automatically flagged out by the system and are not seen in the data or at QA0 unless this hardware flagging has failed. Examples of these are LO lock problems, antenna not on-source, or the calibration device not in position.
scans, periods of time, or targets. An overall assessment of QA0 of the EB is then made based on combining these results. The overall tolerances for this QA0 assessment that have been adopted by ALMA for this Cycle are listed in Section 11.2.1. Apart from these, checks are made in AQUA on other observatory calibration data, including a recent measurement of the flux of the amplitude calibrator, and a relevant measurement of the antenna/pad position.

Based on these results, each SB execution is classified into three main categories:

- **QA0_PASS** are datasets that comply with all the QA0 criteria and that will be used for the final imaging.
- **QA0_SEMIPASS** are datasets that do not fulfill all the QA0 criteria, but contain data that is deemed of some scientific value. For instance, the presence of usable bandpass/amplitude calibrator data. QA0_SEMIPASS data are not included in the final data products, but PIs can access those data from the ALMA Science Archive and reduce them if they wish.
- **QA0_FAIL** datasets are those that are not included in the other two categories. Since they represent unusable/uncalibratable data, they are not visible to the PIs in the Project Tracker or SNOOPI.

In addition, other temporary states can be set at this time, which require further intervention before they can be re-set to one of the above main states:

- **QA0_PENDING_OTHER** is used if the state based on the initial assessment is unclear; it requires additional QA, using QA0+ as well as possibly further checks (see Section 11.3).
- **QA0_PENDING_FLUXCAL** is used if the EB is a Pass but the amplitude calibrator (normally also the bandpass calibrator) has no recent flux measurement (within the last 7 days). Once a flux measurement is available in the archive, this will automatically be set to a Pass.
- **QA0_PENDING_CORRECT_ANTPOS** is used if an antenna was used in the array which had not been fully integrated - usually because the antenna positions in the observatory database (Telescope Monitor and Control Database, or TMCDB) have not yet been updated after an antenna move. Once the positions are measured and updated, this will be set to a Pass.

### 11.2.1 QA0 criteria

QA0 pass/semipass/fail criteria that have been adopted by ALMA during Cycle 6 are the following:

- **Antennas:** Situations that will be considered to render an antenna as not available include issues with the antenna itself, such as $T_{\text{sys}}$ values $> a$ factor of 2 higher than the others in all spws, or $T_{\text{sys}} > 2000$ K in Band 3-6, delays $> 0.5$ ns, relative amplitude on the calibrators $< 10\%$ of the median of other antennas, or if the gain calibrator amplitude or phase fluctuations are significantly higher (factor of $> 3$) than other antennas at the same distance from the array centre. This detects outliers because of antenna problems, which are different from high phase noise from poor weather conditions, noted below.

- **Weather:** If the atmospheric fluctuations are such that the antenna-based phase noise\(^3\) is larger than 1 radian, then this is regarded as too high for this antenna to be safely usable. The number of high phase noise antennas in the dataset is determined. For the EB to be a PASS, the minimum number of usable antennas (ie with low phase noise) is then the same as above; so if fewer than 20 antennas in the 12-m array have phase noise $< 1$ radian, the EB will normally be a Semipass. In practice, such datasets will initially be set by the AoD to Pending Other, meaning that the data quality is marginal. Further checks on the image quality are then carried out using QA0+ or offline before deciding whether the EB should be a Pass or Semipass (see Section 11.3).

\(^{3}\)The antenna-based phase noise is determined in the following way: data from spws and the two polarisations are phased up and combined to get a single phase solution per antenna per phase cal scan. The phase noise for that antenna is then the median value of the phase difference between adjacent scans.
- **Bandpass**: QA0\_SEMIPASS if the bandpass calibrator signal is too weak. The limit on signal/noise ratio is that a sufficiently high SNR (>30) is available per channel only if the data is spectrally averaged down to ≤2 channels.

- **Gain/phase**: QA0\_SEMIPASS if the gain/phase calibrator signal is too weak, with an SNR <5 after averaging all spws or a detected flux < 10% of the predicted value.

- **Execution**: QA0\_SEMIPASS All datasets whose execution failed at some point and contain less than 20% of the expected time on the science target for a given SB execution, or where the Execution Fraction (see 11.2.2) is less than 0.2.

- **Calibrations**: QA0\_FAIL Datasets missing critical calibrations that render them uncalibratable. Execution Blocks only containing usable calibrator data shall be labelled QA0\_SEMIPASS.

- **Storage**: QA0\_FAIL Data that could not be read from the Archive

- **Time-Critical Observations**: QA0\_SEMIPASS if observations have not been carried out at the requested dates and times.

### 11.2.2 Execution Fraction

For operational and scheduling purposes, individual SB executions during Cycle 6 will effectively be weighted by measuring their data quality. This weighting is quantified by an Execution Fraction (EF), where 1.0 is the predicted value, 0.0 means that the data is not usable, and > 1.0 means that the EB is better than predicted. The EF is a single number representing the normalized fraction of the theoretical sensitivity (based on the radiometer equation) that a given execution should have reached, combined with a measure of the expected image quality (see below). It is calculated using the number of unflagged antennas compared with the number that should be in an array in Cycle 6 (i.e., forty-three 12 m antennas on the main array, ten 7 m antennas on the Morita Array, and three TP 12 m antennas), the actual time on-source compared with expected, the mean $T_{sys}$ at the representative frequency compared with that assumed by the OT, and the fraction of antennas with phase noise below a set limit (of 1 radian)\(^3\). Executions with lower data quality than the reference will have fractional executions between 0 and 1.0, while those with better quality (such as observations carried out in unusually good weather with lower $T_{sys}$, or with more than 43 antennas), will have a fractional execution value above 1.0. For Cycle 6, the EF of an EB will be capped at 2. The scheduler then sums the EFs of executions to determine how many more are required.

### 11.3 QA0+ and imaging quality

QA0+ is a second stage of QA0, which runs a mini-pipeline reduction of the channel-averaged data, producing continuum-only images of the science target, phase calibrator, and check source (when available). The idea behind QA0+ is that if an execution is left in a QA0 'Pending Other' state (perhaps because the phase noise is above the QA0 limit, or is unclear for any reason), the QA0+ images can be used as a second level of data quality assessment. The EB can then potentially be set to 'Pass' if the images meet the QA specifications. Typically the QA0+ pipeline takes 20-90 minutes to complete. To speed up reduction, it does not run bandpass or Tsys calibration, and uses the online-corrected WVR datastream. It uses flags from the QA0 process and combines both polarisations and all spws with bandwidth >400 MHz. For imaging it will run a single iteration of CASA clean with a centred simple clean box on the first science target and, for targets with SNR >20, it will try self-calibration on a scan-based time interval and will produce a self-calibrated image.

For EBs requiring QA0+, checks of image quality are performed where possible: the metric used is the peak to integrated flux ratio for a point source (either a check source or the science target if this is compact). Also the peak to residual ratio after a single CLEAN iteration can be used: both indicate how much flux from a point source is distributed over the map due to phase decoherence. Whereas completely random phase fluctuations simply reduce the target flux and scatter power over the whole map, semi-systematic or large-scale atmospheric
phase changes over the array will distort the image in a semi-coherent way. These can be more egregious during long-baseline or high-frequency observing, and their effect can be monitored by observing a check-source as part of the observations\(^4\); QA0+ is then used to assess the image quality. A perfect image has a peak/integrated value of 1.0, ‘good’ image quality is peak/integrated > 0.8, and datasets with values less than 0.5 are regarded as marginal for a QA0 pass. The QA0+ criterion is that datasets with a peak/integrated flux ratio on the check source of >0.5 are set to Pass. For those without a check source, in some cases the science target is bright enough that the images can be compared with previous executions of the same EB; those judged of similar quality are also set to Pass. For executions where the check source peak/integrated ratio is <0.5, or there is no check source or clear science target, the QA0 status is maintained at (or set to) Pending Other. These are unclear cases, and require a further stage of checking, usually offline by the full CASA pipeline, before setting to either Pass or Semipass.

### 11.3.1 Relationship between image quality and phase noise criteria

The quality of images is strongly affected by the phase stability of the atmosphere\(^5\), and so measuring the phase noise does provide a good proxy for the final image quality. However, there is a ‘grey area’, where imaging is required. ALMA uses nearby phase calibrators to measure and correct for residual sky phase fluctuations (phase referencing). But this correction is not perfect. The image quality of these phase-referenced observations is affected by spatial differences in the line-of-sight delays to the target and phase calibrator, as well as baseline errors thought to be due to large-scale, slowly-varying atmospheric structure\(^6\). The spatially-dependent differences between target and phase calibrator are linked to the time-dependent phase fluctuations through the turbulent structure of the atmosphere, so monitoring just the time-dependent fluctuations should still give a reasonably good measure of the overall data quality, even without checking the images. Figure 11.1 confirms that image quality (measured by the ratio of peak-to-integrated flux on a point source, taken from QA0+ results on the check source during the 2017 long-baseline campaign) is correlated with the phase fluctuations (measured by the median phase difference between antennas\(^5\)). The phase is measured in degrees, therefore just depends on the on the changing pathlength compared with the wavelength, independent of the actual observing band. A perfect image has a peak/integrated value of 1.0, ‘good’ image quality is peak/integrated> 0.8, and datasets with values less than 0.5 are regarded as marginal. It shows that datasets with phase fluctuations less than 1 rad almost always give good image quality and so can be set to QA0 Pass immediately. The QA0+ imaging stage is for the marginal cases; some are then set to Pass as the images are deemed acceptable, but a significant number give very poor imaging.

Time-dependent phase noise is not the only parameter that affects image quality, however. Figure 11.1 indicates that observations where the target-calibrator separations is more than 3 deg on the sky tend to have worse image quality - even if the time dependent phase fluctuations are low (<1 radian). This is illustrated further in Figure 11.2, where the image quality is plotted against the phase calibrator separation. For reasonable conditions (phase rms <1 radian), generally there is a slow dependence on separation; ‘good’ image quality is generally achieved when the separation is small (< 4 deg) and phase fluctuations are low (< 1 radian).

### 11.4 Observatory calibration and QA

As well as data taken during the SB execution, the observatory tracks array and antenna performance parameters which vary slowly (typically on timescales longer than a week) and which can affect data quality. They are measured by AoDs and System Astronomers at predefined periods as “Observatory Tasks”, and after major interventions such as antenna moves, Front-End swaps, subreflector fixes, etc, or if significant deterioration of performance is detected during operations. Typically after such interventions, antennas are set to non-integrated state and not used for regular science observing. Antenna integration into the array then requires measurement

\(^4\)Check sources are QSOs within a few degrees of the phase calibrator, and are normally observed for a short time during long baseline or high frequency executions.

\(^5\)Although many of these effects can be removed using self-calibration, when the target is bright enough to allow this - see Chapter 10.5.1

\(^6\)The baseline errors only become significant on the longest baseline observations.
CHAPTER 11. QUALITY ASSURANCE

Figure 11.1: Check source image quality (measured by peak to integrated ratio) shown as a function of phase fluctuations (the mean antenna-based scan difference on longer baselines, in degrees). Data are taken from 2017 long-baseline observations in Bands 3-7 (no dependence on the band was found). Blue points are QA0 pass, red crosses are QA0 semipass. Encircled symbols indicate that the phase calibrator is more than 3 deg separation from the check source. The highest image quality (peak/integrated >0.8) generally occurs for phase fluctuations < 40 deg and for the closest calibrators. The cutoff for an immediate QA0 Pass from AosCheck is < 57 deg (1 rad); higher phase noise datasets are passed on to QA0+ for assessment of image quality.

and ingestion of these parameters. Observations and reduction of such data is done jointly by the AoDs, the DSOANT group, and System Astronomers within the DMG. The product is a set of parameters that are ingested into observatory databases, in particular the TMCDB (Telescope and Monitor Control DataBase) which records the hardware status and history. The TMCDB gives the up-to-date view of hardware parameters, so that they can be used by the system during observations. Problems that would significantly downgrade the quality of the data are solved by dropping the offending antenna from the active array, fixing, re-measuring and updating the problematic parameters. Particular parameters which form these observatory calibrations are:

**Array Calibrations:** Baseline and antenna position measurements, pad Delays

**Antenna Calibrations:** Pointing models, Focus model curves, Surface measurements, Beam patterns (including polarization observations), beam squint (offset between polarizations), IF delays (relative delays between basebands and polarizations), Front-End delays (relative delays between frontends)

**Source Calibrations:** Monitoring of solar-system flux standards, and secondary quasar flux standards
11.5 QA2

QA2 deals with QA at the level of data reduction and imaging using the Science Pipeline or performed manually with the Script Generator by the ALMA Data Reducers Team. It is only at the stage of data reduction that the science goals set by the PI can be compared with the actual value in the data products (i.e., RMS, angular resolution, SNR, dynamic range, etc). The official software for ALMA data reduction is CASA.

During Cycle 6, it is expected that, for the ALMA standard observing modes, the automated Pipeline (see the ALMA Science Pipeline Reference Manual on the Science Portal\(^7\)) will be used for the calibration of the data and imaging. The Pipeline provides a Weblog page with detailed information on the quality of the calibration, which will provide the basis for Pipeline QA2. This information will also be made available to SnooPI during Cycle 6, and the weblogs will also be included in the data packages sent to the PIs. For manually-reduced datasets, the QA2 will be carried out with special purpose scripts whose outputs will be included in the data packages delivered to PIs. (Note: Some datasets can be calibrated by the pipeline, but need to be imaged manually, or with manual intervention in the pipeline. In these cases, a weblog will be provided for the calibration portion, and an imaging script, or imaging weblog provided for the imaging.) The QA2 metrics which determine the success of an observation are given in Section 11.8.

A summary list of QA2 parameters/issues checked during data reduction (calibration of individual Execution Blocks and joint imaging) are:

**Calibration Issues:**

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• Bandpass quality: is the spectral profile flat, well behaved and devoid of spectral lines?
• Flux scale calibration: is the absolute flux scale accurate enough?
• Phase transfer and astrometry: Cycle time and sky separation between phase reference and target; typical and extreme (unflagged) phase differences between phase reference scans.

Final Data Characterization:

• Longest baseline, visibility coverage and time on target (after flagging).
• Synthesised beam (spatial resolution) for specified weighting scheme; specific imaging requirements (if stated by the PI). See Section 11.8 for tolerances.
• Spectral resolution and channel ranges used to make sample images.
• RMS noise in target images: Values are compared with those predicted from data after flagging and with those requested. See Section 11.8 for tolerances. (Note: For Solar Observations, this criteria does not apply as there is no blank sky in most of the images.)
• Residual imaging artefacts: Sidelobe Levels, effects of “missing spacings” and possible dynamic range limitations.
• Combinations of array configurations and/or total power, e.g. amplitude scale consistency.
• Mosaicing and/or contamination by bright sources outside FOV or aliasing in clean.
• Polarization purity: minimum fractional polarization detectable, and polarization angle accuracy (if relevant).

The data reduction scripts sent to the PI will include any modifications to the standard scripts which were found to be essential during QA2.

There are three possible QA2 states a reduced dataset can be placed into. QA2_PASS implies that the scientific goals, as defined by measurable parameters such as noise RMS, LAS and angular resolution, have been achieved at the representative spw within the specifications listed in Section 11.8. QA2_SEMIPASS refers to those datasets that fall short of meeting the PI requested science goals, but are otherwise of good quality (or as good as possible with the observations that could be achieved). QA2_FAIL is a temporary state used during an observing cycle when an observation fails to meet the PIs goals by a significant margin (see Section 11.8 of this Handbook) and needs to be scheduled for additional observations. Although the Observatory will try to take a similar quality of data in all spws, those measurable parameters in non representative spws will not be used for the judgement of QA2_PASS or QA2_FAIL.

If data fail to obtain enough additional observations to pass QA2 by the end of an observing season and they are not grade A proposals, their QA2 state is changed to QA2_SEMIPASS and they are delivered to the PI.

11.6 QA3

QA3 is post-reduction evaluation of the data products delivered to the PIs. It is advisable that PIs check the data products themselves and report any problems that they find to their Contact Scientist via the ALMA Helpdesk (https://help.almascience.org/). The QA3 process will be triggered by the PI, the contact scientist or other ARC personnel. They will open an ALMA Helpdesk ticket reporting a problem with the data products which may reflect an underlying problem with the data, observing procedure or calibration. The ARC receiving the Helpdesk ticket will retrieve the data from the archive and evaluate the nature of the problem. The evaluation by the ARC should include an assessment on whether the problem is present only in a particular dataset or whether others taken under similar set-ups and conditions also show it. If the problem is deemed to reflect a problem with the performance of the array, the calibration or data reduction processes, or the QA process, the ARC will
communicate their findings to the observatory, which will work on solving the problem in collaboration with
the ARCs. The result will be communicated back to the reporting investigator. An extension of the proprietary
period of delivered datasets with QA3 issues will be granted based on the policies in the Cycle 6 User Policies
document.

11.7 The Quality Assurance Report

During Cycle 6 the Quality Assurance Reports of QA0 and QA2 will be accessible to the PIs using SnooPI.
These reports will be generated using the AQUA tool. For Pipeline calibrated datasets, additional reports
(Weblogs) will be sent to the PI as part of the data delivery. For non-standard observing modes that require
manual data reduction, script-generated QA2 reports will be delivered instead.

The basic unit of a Report is the ObsUnitSet, which represents part of the scientific goal stated by the PI
during Phase 1 (project creation and review). An ObsUnitSet will typically contain several executions of an SB.
For each execution, a QA0 report is generated by the AoDs using the information available at the time of the
observations, which includes TelCal outputs and other monitoring data (weather, total power levels, Corr GUI
outputs, etc). A given execution is only cleared for reduction if it has passed QA0. There will be only one QA2
report for the whole ObsUnitSet generated by the JAO/ARCs at the end of the data reduction process; this
also has to be approved before the data products are delivered to the PI. It is expected that sometime during
Cycle 6, a complete interface between the Pipeline and AQUA will be made available. From that point onwards,
complete QA Reports will be generated using the AQUA software and delivered as part of the data products
package (including the Weblog information). The final report per ObsUnitSet delivered to the PI in that
case be a concatenation of all the relevant QA0 reports per execution with the QA2 report. Comments on each
stage of the QA process (with supporting images, if required) would be added to the Report.

The standard policies for QA0 failures are that the observations of those Execution Blocks that failed have
to be repeated. Failures to pass QA2 may trigger additional observations if the achieved RMS or imaging
parameters do not fulfill the QA2 pass criteria. Additional observations may not be possible in circumstances,
such as projects with very tight weather, very specific configurations, or time constraints. If the available data
are insufficient to reach the required sensitivity, but are otherwise of good quality, they will be released to the
PI at the end of the cycle.

11.8 Overall QA Criteria

For an execution to be considered QA0_PASS it must have a phase RMS (post-calibration) of \( \leq 1 \) rad in
most antennas, enough calibrator data to be able to calculate necessary calibration terms (bandpass/amplitude,
complex gain, and, if relevant, polarization), and enough science data to be useful (see above). Furthermore, the
execution has to be "significant" in terms of the fractional execution (i.e., more than 0.2). If it does not meet
these criteria, but has some useful data (calibrator or source), it is classified as QA0 SemiPass. Such data are
available through the ALMA archive, but they will not be used in the generation of the imaging data products.
If neither of these hold, the execution is deemed to have no useful data and declared QA0_FAIL.

For any other situation, the data will be accepted, although it may require some additional flagging for
mis-behaving antennas, baselines, etc.

For QA2 at the Member OUS, MOUS, level, the main criteria are:

- The achievement of the requested noise RMS in the images (it must be within 10\%, 15\% and 20\% of the
goal for Bands 3/4/5/6, 7/8, and 9/10, in flux units per beam, respectively)
- Achievement of the synthesized beam shape (In order to satisfy the QA2 criteria, the angular lengths of
the major and minor axes of the synthesised beam must be within the angular resolution range requested

\[8\]As stated above, processing of group ObsUnitSets will not be regularly done during Cycle 6. PIs should assume that all
statements here apply to member ObsUnitSets.
by the PI, taking into account the modifications applied by OT to ensure schedulability of the SB (See Proposer’s guide). This angular resolution range can be viewed at SB level in the OT in Phase 2 (Min. and Max. Angular Resolutions).)

- Reaching the required calibration quality (phase RMS of 5 degrees after calibration, and absolute flux scales within the ALMA specifications\(^9\)).

Observations that do not satisfy the QA2 criteria will be released as QA2 semipass if there is no realistic chance of further observations improving the data quality to the level of a QA2 pass, or if further observations are timed out. For the individual MOUS that must be combined at the GOUS level, the RMS values quoted above should be scaled by the relative sensitivities of the configurations, using the most compact 12-m Array configuration as reference.

\(^9\)Absolute Calibration accuracy is regularly checked as an observatory task.
Chapter 12
Data Flow and Logical Data Structure

This chapter describes the data flow process from the observations until raw data is ingested into the ALMA Archive. It includes a brief description of all the main software subsystems involved in the data acquisition and archiving, as well as a summary of the data structure adopted by ALMA.

12.1 Data and Control Flow

This section describes the overall control of the ALMA system and the flow of data during observations. A summary of the main actors and operations involved in the observations is shown in Figure 12.1. Each of the light blue boxes in Figure 12.1 represent an ALMA subsystem involved in the observations. The rest of the boxes, colored according to the actor involved, include labels for the actions performed either by external agents (actors) or by those subsystems.

The software subsystems run within a cluster environment known as a Standard Test Environment (STE), with the STE used for the real array operation referred to specifically as an ALMA Production Environment (APE). The STE incorporates central server nodes, embedded computers for hardware control, dedicated correlator data processing nodes, and console nodes for user interaction.

A typical observing session would be started by the Telescope Operator interacting with the Executive subsystem. The Executive subsystem is in charge of starting up the ALMA Common Software (ACS) and its Common Object Request Broker Architecture (CORBA)-based services and then initializing all of the various software subsystems involved in the observing and data storage process. Once all the components are ready, the Executive also handles asynchronous events from several of the subsystems and responds to them accordingly. Among the events, the Executive also publishes a list of error conditions to the attention of the operator and the requests for status of the Control, Telescope Calibration, Correlator and Scheduling subsystems.

The actual observations start by the operator creating an array, which means selecting the antennas, photonic reference and correlator (if needed by the observations\(^1\)) to be used. The antenna selection is largely automated by filtering based on antenna types, connection to either correlator, and status of the antennas in the ALMA Dashboard\(^2\). Up to six independent arrays can exist and operate at one time. Once an array is successfully created, Scheduling Blocks (SBs) can be queued and run on it. SBs can either be manually chosen, e.g. for array calibration or test purposes, or selected from a ranked list produced by the dynamic scheduler for regular science operations.

An SB is the smallest, calibratable element of a project. SBs are XML documents which are part of the ALMA Project Data Model (APDM), and are kept in the archive along with the proposal and other ancillary information. SBs have an associated state which among other things controls whether they can be executed.

\(^1\)Most ALMA observations use a correlator, either for interferometry or single dish spectroscopy, but single dish continuum observations can be made using only power detectors in the antennas.

\(^2\)http://dashboard.alma.cl/, Pietriga et al. 2014, SPIE, 9152, 1B.
For science SBs to be executable the Phase-2 process\(^3\) should have been completed for them, and they should not have been fully observed or cancelled/timed-out at the end of an observing cycle.

The dynamic scheduling software considerations include the target elevations as a function of time, observing conditions (PWV and estimated path length variability), actual array configuration, project completion status, proposal grade and rank, among other things. The dynamic scheduler will also suggest array calibration SBs at appropriate times. It is intended that by Full Operations the algorithm will be optimized to the point that the dynamic scheduler can fully automatically queue SBs during normal operation periods. At present the Astronomer on Duty will select SBs in the dynamic scheduler GUI to be added to the array’s scheduling queue for execution.

When an SB reaches the head of the array’s scheduling queue it starts an execution. The SB contains the name of a Python observing script, which is run by the Control subsystem and passed a Python representation of the SB content. The observing scripts, which reside in the Science Software Requirements (SSR) module of the software, use the SB content to determine the necessary execution sequence. The observing script commands the Control subsystem with sequences of scans and subscans. The scan and subscan specifications specify for each subscan the tuning, phase center and antenna pointing, calibration device position, intent metadata (to convey the purpose of the scan/subscan), and other parameters that may need controlling from the observing script level. Control executes the scan/subscan sequences by commanding all relevant hardware, the relevant correlator subsystem (Baseline (BL) or Atacama Compact Array (ACA)) and the Total Power Processor (TPP). Subscans, and most actions performed by the Control and correlator subsystems, start and end on 48 ms (Timing Event)
boundaries which are accurately synchronised among the whole system. The subscans result in raw data files and metadata rows being sent to the archive subsystem and made available to the online Telescope Calibration (TelCal) subsystem. TelCal publishes results from calibration scans it reduces, which are sent to the archive as calibration tables, displayed by GUIs to allow Operator and AoD evaluation of the observation progress, and received by the observing script to apply online corrections e.g. for antenna pointing and focus. For details

Each run of an SB produces an Execution Block (EB) that is stored into the archive through two parallel paths: metadata tables compiled by the "Data Capture" component of the Control subsystem, and binary data files which are streamed directly from the data producers via the "Bulk Data" transmission system. The metadata contains the relevant references to the binary data files to allow them to be exported from the archive along with the metadata. The metadata tables are defined by the ALMA Science Data Model (ASDM, see Section 12.2). The tables are primarily stored as XML documents, but larger tables are written in a binary format and stored in the same way as the raw data binary data files. Both the metadata and binary data are received by an online archive component running on the STE, which provides buffering and handles the communication between the STE and the archive for the data storage. For more information on the archiving, please read Chapter 13 of this Technical Handbook.

The Data Capture software is as an interface between the real-time domain of the data taking and the storage side (see e.g., 2006 ADASS contribution by Hafok, Caillat and McMullin). It receives many inputs from the Control, Correlator and TelCal subsystems for each subscan and scan, which it compiles together into the ASDM tables. The Control software itself is also compiling values from many hardware devices into the information it provides to Data Capture. At the end of each scan the current table contents pertaining to the scan (an ASDM "slice") are passed to TelCal to allow it to produce any necessary calibration results (combined with binary data received directly from the Bulk Data streams), and when TelCal finishes it passes back results for Data Capture to append to calibration tables. When the execution finishes, Data Capture sends the finished ASDM tables to the online archive component for storage (in future the storage may be made more incremental to improve scalability), waiting for the last scan results from TelCal if needed. Data Capture is also responsible for monitoring the storage of the binary data to the archive at the end of the execution and reporting when the data is fully archived or if there was an error.

The archiving at the end of an execution can be carried-out in parallel to further executions, so the next SB from the queue can start shortly after the last scan of the previous execution. This is achieved by having an independent Data Capture component run for each execution, which remains alive until all the data from the execution is archived. At present the array is not allowed to be destroyed until all Data Capture components for executions in the array have completed archiving.

A summary plot of the main elements involved in data flow is shown in Figure 12.2.

12.2 The ALMA Science Data Model

The ALMA Science Data Model (ASDM; Viallefond, F. 2006, Astronomical Data Analysis Software and Systems XV, 351, 627) defines the collection of information recorded during an observation that is needed for scientific analysis. The ASDM is a common standard shared by ALMA and the EVLA. As described above it contains both binary data files (BDFs) and metadata tables. The metadata tables contain links to other tables and references to the binary files in the Archive.

The ASDM contains 16 core tables that are common to all observing modes, and up to 23 additional tables that are only created for specific observations. On top of these, TelCal also creates associated tables whenever it processes any calibrations. All tables are organized with a similar structure, with the columns listing the contents and the rows including the actual values. The core tables have been defined to outline some of the following: hardware characteristics, array configuration, antenna tracking, targets, tuning setups and resulting spectral windows, scan and subscan timing and intents, and flagging information.

The tables produced by TelCal have names prefixed by "Cal" (there is also the CalDevice table which is

\[^4\text{http://adsabs.harvard.edu/abs/2011ASPC..442..277B}
\[^5\text{see e.g. https://casa.nrao.edu/casadocs/latest/reference-material/the-science-data-model} \]
unrelated to TelCal). The list of these calibration tables expands as new observing modes and calibrations become available (see Chapter 10).

The ASDM metadata implementation in the ALMA software uses automatic code generation to provide common libraries for Python, Java and C++ to read and write ASDMs. These are used in the various parts of the software, both online and offline, which need to interact with ASDMs.

The BDF format is a mix of XML headers and actual binary data, encapsulated in multipart MIME. Data recorded in each file can include cross-correlations, auto-correlations, and flags (“BDF flags”). For a typical subscan with a correlator, four BDFs are produced for different types of data: spectral correlator integrations, channel-average correlator integrations, Water Vapour Radiometer (WVR) integrations, and baseband Square Law Detector (SQLD) integrations. The BDF structure is designed to support incremental writing. Currently all but the WVR BDFs are written incrementally. Typically one incremental write includes data from one integration or all integrations within a second when the integration time is shorter than a second. BDF file sizes can be up to ~100GB at present.

Most users interact with the ASDMs via the Common Astronomy Software Applications (CASA) data reduction package. CASA uses a different internal storage format called the Measurement Set (MS)\(^6\). A piece of software distributed with CASA known as the “Filler” is responsible for producing an MS from an ASDM. It is invoked either from the command line as `asdm2MS` or by the CASA task `importasdm`. The filler uses the same ASDM libraries as the ALMA software for parsing.

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\(^6\)see e.g. https://casa.nrao.edu/casados/latest/reference-material/measurement-set
Chapter 13

Data Archiving

13.1 Introduction

The ALMA Archive is at the center of the ALMA data flow (Fig. 13.1). It is a combined database and binary data storage system that is accessed by the different software subsystems through the same software layer. The ALMA Archive stores all metadata and data of ALMA including the user accounts, proposals, antenna configurations, monitoring, raw, and reduced science data. The science data and metadata are made available to PIs and archival researchers for query and download following ALMA's data access policy (http://almascience.org/documents-and-tools/latest/alma-user-policies).

The ALMA Archive is divided into two parts. The ALMA Frontend Archive (AFA) and the ALMA Science Archive (ASA). The AFA provides the core persistence functionality. The ASA holds a small subset of metadata of the AFA in a relational database and provides access to external interfaces like the Archive Query interface and in the future Virtual Observatory (VO) tools. The storage architecture is based on the Next Generation Archive System (NGAS) with Oracle technology for replicating the metadata.

Each of the three ALMA Regional Centers (ARCs) in North America, Europe and East Asia holds a copy of the entire ALMA Archive for backup, user support and data distribution to the ALMA PIs and archival researchers. The ARCs provide a completely identical user experience to their communities and a user can download data from any ARC.

An ALMA Science Archive Manual is available from the ALMA Science Portal.\(^1\)

13.2 Data Flow and Archive

Data from the correlator, together with monitoring and weather data, are sent via dedicated optical fiber links to the OSF, where they are archived. A peak data rate of 66.6 MB/s can be sustained for short periods of time, i.e. days. This peak rate is a technical limitation of the data capture and data flow systems and will be imposed by the Observing Tool at the proposal validation stage.

The Pipeline processing system and Archive storage system have been designed to cope with an average data rate of 10% of the peak rate, i.e. 6.6 MB/s leading to a yearly amount of 200 TB of data. The volume of data products from the pipeline are expected to approximately equal the raw data volume, resulting in about 400 TB/yr being ingested into the ALMA archive.

As soon as data are taken and have passed (or semipassed) the QA0 quality control step (see Chapter 11.2), their metadata are harvested from the AFA into the ASA and made available for search. This allows archival researchers to see which data will become public in the future. Metadata harvested include all the information

\(^1\)http://almascience.org/documents-and-tools/latest/science-archive-manual
needed to describe the observations, including date and time, source coordinates, frequency settings for each spectral window, and spectral and spatial resolution.

The ALMA Archive at the OSF is designed to provide up to a year of temporary storage for the instrumental data (in the form of files in the “ALMA Science Data Model”, or ASDM, format) and the monitoring data. The instrumental data are then transferred to the main archive at the SCO, where the pipeline is run and from where the data and pipeline products are distributed to the three ARCs. At this stage, a sizeable fraction of the data are processed at the ARCs and then mirrored back to SCO. The process of copying the data to any of the archives involves a replication of the metadata (support data) and of the bulk data (ASDM and FITS files). All data transfer is done over the network. Metadata replication from SCO to the ARCs happens within a few seconds, transfer of bulk data can take longer (up to several hours), depending on the amount of data to transfer to the individual ARC.

Once the pipeline has processed an ObsUnitSet (hereafter OUS, see Chapter 8 for details on the OUS structure) and the science products have passed QA2 they are ingested into the ASA.

PIs of large programs are expected to return processed data products to the project as described in the Proposer’s Guide. These data will be transferred to JAO from where they will be ingested into the ASA.

13.3 PI Data and Data Delegation

The unit of data delivery to the PI is the OUS. As soon as an OUS has passed QA2 and the data products have been replicated to the user’s home ARC, PIs will receive an email notification containing a link to their data and supporting information. This sending of the notification to the PI triggers the start of the proprietary period of 12 months for standard proposals and 6 months for Director’s Discretionary Time (DDT) proposals. Extensions of the proprietary period can be granted under certain circumstances (see Section 8.4.4 of the ALMA User’s Policy). When a Group OUS consisting of several Member OUSs is processed (for example, a combined TP, 7-m and 12-m Array observation), the Group OUS products are released separately, with their own 12-month proprietary period.
13.4 Archive Query

The data deliveries consist of one or more bundles of raw (ASDM-format) data and one "products" bundle tar file which includes the FITS files, logs, scripts, QA information and calibration tables.

Often PIs want to make their proprietary data available to collaborators, e.g. CoIs. To this end, a data delegation service is available so that PIs do not have to give away their Science Portal password to anyone. Instead, PIs can give access rights to the data of a project to any registered ALMA user. To do so, PIs need to log into the Science Portal, go to their user profile page in the top right corner of the Science Portal page and then add delegates in the "Project delegation" tab.

13.4 Archive Query

At http://almascience.org/alma-data/archive users can query the holdings of the ASA (Figure 13.2). They then can download the data corresponding to their queries, if those data are public or if the users are authenticated and have the proper access rights to those data.

![ALMA Science Archive Query](http://almascience.org/alma-data/archive)

Figure 13.2: The ALMA Science Archive Query Form

Queries can be made by physical quantities along the Position, Energy, Time and Polarisation axes. Help on how to query is provided in the tooltips of the query fields as well as through the "Query Help" link. Query constraints using standard operators for strings (*, ?) and numbers (>, <, .. ) can be placed into the form fields. There is also a name resolver available for non-solar system objects (Sesame), which queries the Simbad, NED and Vizier databases. No operators can be used in the name resolver field.

By default, the Archive Query Interface (Fig. 13.2) will present users with the metadata of the observations, although they can choose to also see the metadata per project or per publication. On the results page (Fig. 13.3), users can sort and subfilter the results and add or remove columns from the result table to narrow their search. The sky coverage of the observations is visible in a graphical window.

The results page also offers to download the results in VOTable, TSV (Tab-separated values) or CSV (Comma-separated values) format for further processing in tools like topcat². The URL used for the exporting

²http://www.star.bris.ac.uk/~mbt/topcat/
of the query results can also be modified to access the ALMA Archive queries programmatically. Astroquery\textsuperscript{3} can be used to encapsulate and simplify programmatic access to the ASA. Neither Topcat nor Astroquery are part of the official ALMA software, and support of users of these tools will be limited to best efforts.

### 13.5 Request Handler

Once data of interest are defined, they can be selected via checkboxes on the results page and submitted to the ALMA Request Handler for download. The data are displayed in the full OUS hierarchy. This allows users to check immediately if there are data from other Member OUS or higher-level Group OUS products available. By default, only data products (i.e. images and cubes as well as ancillary processing documentation) will be selected for download. Users who want to download the raw data as well are asked to select “include raw” before hitting the Download button or to select the desired data manually. The sources row can be expanded to show all the source names that are part of the package. Furthermore, a readme file is provided summarising the structure of the data and providing the full links to the OUS.

Four possibilities exist for the download itself:

- The first option is to use the download script. This script runs under Linux and MacOS and downloads files in parallel streams and is also adapted for downloads to a processing environment where no web browser is available.

- The second option is to use the download manager applet with the FireFox browser. This method is very convenient and allows for parallel downloads, too. It requires a Java Browser Plugin.

- It is also possible to use the download manager through Java Webstart technology.

\textsuperscript{3}https://astroquery.readthedocs.org/en/latest/alma/alma.html
Finally, a page with the links to all selected files can be displayed which then can be conveniently downloaded, e.g. using a browser plugin like "DownThemAll".

The Request Handler also allows selecting individual files for download. Only the data deliveries that users have permission to download can be selected. If they are not authorised to access any data delivery in the request, no "Download selected" button appears.

If the user is authenticated before requesting data for download, the request will be stored. This allows users to go back to previous requests. Note that these requests are stored only at the ARC the user is currently accessing. If data should be downloaded from a different ARC, then a new request has to be issued. For very large data requests PIs or archival researchers have the possibility to ask via the ALMA Helpdesk for data delivery on hard media i.e. USB hard-disks. The ARCs may have different policies regarding the details of the shipping of the hard-disks.
Appendix A

Antennas

A.1 Design and Properties

ALMA is composed of a total of 66 antennas, 54 with a diameter of 12 m and 12 with a diameter of 7 m. The four 12 m antennas used for total power observations and the twelve 7 m antennas together form the Atacama Compact Array (ACA). The ALMA antennas are manufactured by three different contractors. These are VertexRSI (North America) which provided 25 12 m antennas, Alcatel Alenia Space European Industrial Engineering MT Aerospace (AEM, Europe), which also provided 25 12 m antennas and Mitsubishi Electric Corporation (MELCO; East Asia), which provided the four 12 m total power antennas and the twelve 7 m antennas (Figure A.1).

All antennas have been designed to meet very stringent ALMA performance criteria, and to successfully operate under the extreme environmental conditions at the Array Operation Site (AOS), i.e. strong winds, large temperature ranges and gradients, solar irradiation and snow. The primary operating conditions are the following:

- Range of Ambient Temperatures: \(-20 ^\circ C \leq T_{amb} \leq +20 ^\circ C\)
- Gradient of temperature: \(\Delta(T_{amb}) \leq 0.6/1.8 ^\circ C\) in 10/30 minutes
- Wind Velocities \(\leq 6/9 m/s\) (day/night)
- Full solar loading

The antennas have the following specifications within the Primary Operating Conditions:

**Antenna Surface:** RMS deviation of 25 (20) microns or less for 12 m antennas (7 m antennas) relative to an ideal parabola.

**Pointing:** Absolute pointing \(\leq 2.0\) arcsec all-sky. Offset pointing \(\leq 0.6\) arcsec within a 2 degree radius on the sky.

**Primary Beam:** The total power pattern response of each ALMA antenna shall be determined to a measurable and repeatable precision better than 1% at frequencies <400 GHz and 2% at frequencies >400 GHz.

**Subreflector:** 6 degrees of freedom to allow for alignment with the corresponding receiver beam.

**Subreflector Motion:** Maximum horizontal (X) and vertical (Y) displacements of \(\pm 5\) mm. Maximum focal displacement (Z) of \(\pm 10\) mm. The maximum rotation around the axes is 1.2 degrees. Positioning must be accurate to 5 microns.
Antenna Location: The phase center position of the ALMA antenna shall be determined to a radial precision of 65 microns (including the antenna structure and pad), stable over two weeks.

Configuration: The ALMA antennas shall be relocatable.

Lifetime: a minimum of 30 years.

Antennas used during ALMA Cycle 5 have both 12 meter and 7 meter diameters, with the receivers mounted at the secondary (Cassegrain) focus. The 12 m dishes have a focal length of 4.8 meters, but the distance from the secondary focus to the plane of the subreflector of the 12 m antennas is 6000 mm, giving an effective focal ratio f/8, with an effective secondary focal length of 96 m and a plate scale of 2.15 arcsec per mm. The subreflector has a diameter of 750 mm. The 7 m dishes have a focal length, to the primary focus, of 2.572 meters. Given an effective focal ratio f/8, an effective secondary focal length is 56 m. The subreflector has a diameter of 457 mm.

The main reflectors of the ALMA 12 m and 7 m antennas are composed of individual panels. The size and number of panels varies between the different types of antennas:

VertexRSI: 264 panels spanning 8 rings with 12 (rings 1 and 2), 24 (rings 3 and 4), and 48 (rings 5 through 8) individual panels which are roughly a half-meter-square in area.

AEM: 120 panels spanning 5 rings with 8 (ring 1), 16 (ring 2), and 32 (rings 3 through 5) individual panels which are roughly one-meter-square in area.

Melco 12 m: 205 panels spanning 7 rings with 5 (ring 1), 20 (rings 2 and 3), and 40 (rings 4 through 7) individual panels which are roughly one-meter-square in area.

Melco 7 m: 88 panels spanning 5 rings with 4 (ring 1), 12 (ring 2), and 24 (rings 3 through 5) panels which are each roughly one-meter-square in area.

Figure A.1: The four different ALMA Antenna designs: Vertex 12 m, MELCO 12 m, AEM 12 m, and MELCO 7 m (from left to right).

Each panel has up to 5 adjustment screws, which can be used to optimize the surface accuracy of the individual antennas (based on holographic measurements). The surface of the panels are etched to scatter optical and near infrared solar radiation.

The antennas are equipped with a movable aluminium subreflector. Subreflector adjustment is used to maximize the transfer of power into the receivers by compensating for changes in the focus position due to gravitational- and temperature-induced deformations. The backplane of the subreflector is attached to a hexapod that controls its position and orientation. The hexapod has six degrees of freedom, displacement and tilt around the three axes, horizontal (X), vertical (Y) and along the optical axis (Z).
### A.2 Antenna Foundations

The antennas are placed on specially-designed concrete pads to guarantee stable orientation and location (Figure A.2). All antennas are attached to the pads at three points at the vertices of a triangle. The three points (inserts) are located on a circle centered at the antenna pad with a spacing of 120 degrees.

This interface guarantees a position repeatability error of the antenna, considered as a rigid body, not exceeding the values below:

- $X/Y$ plane < 2 mm (peak to peak)
- Rotation around $Z$ < 30 arcsec (peak to peak)
- Parallelism with respect to $Z$ +/- 10 arcsec with respect to Zenith

#### Table A.1: Design Properties of the Different ALMA Antennas.

<table>
<thead>
<tr>
<th>BUS</th>
<th>Number Rings/ Panels</th>
<th>Panel Material</th>
<th>Quad type</th>
<th>Cabin</th>
<th>Drive System</th>
<th>Metrology System</th>
</tr>
</thead>
<tbody>
<tr>
<td>Vertex</td>
<td>CFRP 8/264</td>
<td>Al +</td>
<td>Steel</td>
<td>Gear</td>
<td>4 linear displacement sensors + 1 two-axis tiltmeter (above the azimuth bearing)</td>
<td></td>
</tr>
<tr>
<td>Invar</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Melco 12 m</td>
<td>CFRP 7/205</td>
<td>Al +</td>
<td>Steel</td>
<td>Direct</td>
<td>Reference Frame metrology</td>
<td></td>
</tr>
<tr>
<td>Steel 5/88</td>
<td>Al +</td>
<td>Steel</td>
<td>Direct</td>
<td>Thermal (main dish), Reference Frame metrology</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Melco 7 m</td>
<td>CFRP 5/120</td>
<td>Nickel x</td>
<td>CFRP</td>
<td>Direct</td>
<td>86 thermal sensors + 2 tiltmeters in yoke arms</td>
<td></td>
</tr>
<tr>
<td>Invar</td>
<td>Rhodium</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Notes: 1 Shape of the quadrupod supporting the subreflector as seen looking along the optical axis of the antennas when they are pointed to the viewer. 2 A gear drive consists of a main motor driving a series of connected reduction gears (i.e., gearbox) that do the actual precision work. A direct drive system does not require such gears and takes the power directly. The direct drives used in ALMA antennas are magnetically supported. 3 Jointly used to correct in semi-real time the pointing of the antennas, under a wide range of environmental conditions, to meet the ALMA specifications.

All antennas have a Cassegrain cabin that is kept at a constant temperature of 20 degrees Centigrade and contains the receivers, the amplitude calibration device and associated electronics.

A shutter protects the inside of the Cassegrain cabin when the antenna is not operating. A membrane transparent to the frequencies that can be observed with ALMA is located below the shutter to prevent airflow from the cabin to the outside when the shutter is open. The current design uses a 0.5 mm thick Goretex membrane.

The different antennas use a combination of steel, aluminium, Carbon Fiber Reinforced Polymer (CFRP) and Invar to achieve the best compromise between stiffness, robustness, smoothness, and low thermal expansion (see Table A.1 for a summary of properties). Common to all antennas is that they have a steel pedestal.

All antennas have built in metrology systems which allow thermal and wind deformations to be computed and corrected. For these purposes, the antennas are fitted with thermal sensors, linear sensors and inclinometers (tiltmeters).

The Vertex antennas have a drive system that is gear-driven whereas the AEM and MELCO antennas have magnetically supported direct drives.

The antennas are controlled using the ALMA Control Software (ACS). ACS sends instructions to the Antenna Bus Master (ABM) computer, which are then sent to the Antenna Control Unit (ACU) through a CAN bus.
Figure A.2: Structure of an antenna pad (actual pad at the OSF) (left) and detail of antenna anchored to a pad (right).

The minimum stiffness which the foundation must exhibit at each insert is:

- Vertical stiffness (Z) > 13 x 10^9 N/m

- In X/Y plane > 9 x 10^9 N/m

Stiffness includes the inserts, the concrete pad and the soil. This does not include the kinematic mount lower part nor the foot of the antenna. The position of the pads are measured to a precision of 65 microns, and then monitored for stability for over two weeks. The pads are equipped with two vaults that contain the power, communication, Local Oscillator (LO) and data transmission cables that are connected once the antenna is placed on the pad.

Figure A.3: The ALMA array with eight 12 m antennas (left), and an antenna being transported to the AOS (right).
A.3 Antenna Transportation

Antennas are moved from one pad to another using a specially-designed transporter (Figure A.3, righthand panel). ALMA has two of these vehicles. They are 20 meters long, 10 meters wide and 6 meters high, and each has 28 tires. The transporter positioning system performs a fine positioning of the antenna before setting it down on the foundation in the 3 in-plane degrees of freedom (x, y, rot-z) and 2 in-tilt (rot-x, rot-y). Adjustment in each of the 5 adjustment axes can be done independently. The adjustment range of the antenna positioning system compensates for the inaccuracy of the vehicle position with respect to the antenna foundation (which must be smaller than 10 cm) to achieve the required antenna positioning accuracy. The antennas can be positioned to within a few millimeters, ensuring accurate placement on the antenna foundation pads. More information on the transporters can be found on the ALMA EPO pages\(^1\).

![Figure A.4: Side view of ALMA front end showing cryostat assembly, with room temperature unit below.](image)

A.4 Cryostat

The ALMA front end consists of a large closed-cycle 4 K cryostat containing individual cold cartridge assemblies (CCA) with mixers and LO injection for each band, along with room temperature electronics for the IF and LO for each band (the warm cartridge assembly, WCA) and fore-optics and entrance windows for each band.

\(^1\)http://www.almaobservatory.org/en/technology/transporters
The water vapor radiometer (WVR) is mounted to one side of the cryostat using a pickoff mirror to direct the antenna beam into the WVR. The Amplitude Calibration Device (ACD) is mounted above the front end, and is described in Section A.5. Figures A.4 and A.5 show overviews of the front-end unit, with the cylindrical cryostat on top and the room temperature electronics beneath.

All of the receiver cartridges are in the same cryostat, with the mixers thermally-coupled to the same 3-stage Sumitomo cryocooler (Figure A.6). The three stages have nominal temperatures of 4 K, 15 K and 110 K. To avoid overloading the cooler, only three bands can be switched on at a time. It takes about 1 minute to switch between any of the bands that are switched on at a given time. For bands that are off, the time to fully thermally-stabilize them from an off state is 15 minutes – this is mainly to ensure a flat bandpass shape. All of the receivers are mounted off-axis to avoid extra rotating band-selection mirrors, which necessitates a pointing offset of the antenna to change band. The band pointing offsets are known and well-measured; the reference band for pointing is Band 6, and all offsets are with respect to this band. The four higher-frequency bands (Bands 7-10) are mounted close to the central boresight to minimize aberrations.

Figure A.5: Bottom view of ALMA front end, showing WCAs.
Figure A.6: Views of cryostat assembly, showing different windows (top) and the portholes for the WCAs for each band (lower view).
A.5 Amplitude Calibration Device

The ALMA specification for relative amplitude calibration repeatability\(^2\) has been set to be better than 1% for frequencies below 300 GHz and better than 3% for all other frequencies covered by the ALMA Front End. To achieve this goal, ALMA has adopted a two-load amplitude calibration approach.

The Amplitude Calibration Device (ACD) is located above the cryostat. It consists of a robotic arm attached to the top plate of the front end (Figure A.7). The arm holds two calibration loads, one at ambient (i.e., receiver cabin) temperature and the other one maintained at 80°C (353 K). In addition, this arm also holds a solar filter to attenuate solar radiation during observations of the Sun. The arm is designed to allow the two loads to be placed in the path of any of the receiver beams (Figure A.8). Typically it takes 2 seconds to move the arm from the park position to the position where one of the loads is in the beam, and also 2 seconds to change between loads.

![Figure A.7: Lateral view of the ACD on top of the ALMA front ends.](image)

To accurately calibrate radio astronomical data to a temperature scale, the actual brightness of the two loads has to be precisely known. Critical to this calibration precision is the coupling of the load to the beam of a given band. This coupling must be very good at any telescope elevation and free of reflections of the load emission. This is because any reflection from the loads back into the cryostat would be terminated at a different temperature and would cause standing waves. Both loads have thus been designed so that the actual effective brightness temperature and that computed from the measured physical temperature (with sensors embedded in the loads) using known emissivities differ by, at most, ±0.3 K and ±1.0 K for the “ambient” and “hot” loads, respectively. This requirement also sets a limit to the fluctuations and departure from the set temperature that are allowed for the “hot” load. Furthermore, the return loss specifications for these loads are -60 dB and -56 dB, respectively.

\(^2\)“Calibration Repeatability” means being able to make repeated measurements of the same flux densities (or brightness temperatures) for the same source under different conditions (weather, telescope elevations, front-end status, etc.).
Figure A.8: Top view of an ALMA front end showing the robotic arm of the ACD retracted during normal observations or on top of one of the front-end inserts for calibration. The current design has been improved by placing all the loads in a wheel.

A.5.1 Atmospheric Calibration Procedure

The ACD is used to measure the receiver temperature and the sky emission by comparing the signals on the sky, ambient and hot loads. This is known as atmospheric calibration (ATM calibration), and is required to correct for differences in the atmospheric transmission between the science and the celestial amplitude calibrators. Normally ATM calibration is done during observations, both near the science target, as well as near the amplitude calibrator.

Traditionally, most mm and submm observatories have used the single-load calibration method, but several simulations have shown single-load calibration is not capable of reaching the relative amplitude calibration accuracies required by ALMA at all of its observing frequencies. However, that method has the very desirable feature that it is only weakly dependent on the opacity of the sky at the time of the observations. A method, using the two calibration loads within the ACD, has been devised in the past to try to achieve the same weak dependence on the opacities at the time of the observation. This method (“the α method”) uses the voltage outputs from the observations of both loads to simulate a single load with a brightness temperature close to that of the atmosphere at the observing frequency. This fictitious single load is defined as a weighted sum of the voltages of the “hot” and “ambient” loads so that the temperature calibration factors are almost independent of the optical depth. The fictitious load voltage output, \( V_L \), is defined as:

\[
V_L = \alpha V_{L_1} + (1 - \alpha) V_{L_2} \tag{A.1}
\]

where \( \alpha \) is the weighting factor, and \( V_{L_1}, V_{L_2} \) the output voltages when the two loads are measured. From this definition and some algebra, one can find the optimum weighting factor needed to minimize opacity dependence, and the corresponding resulting calibration factors are:

\[
\alpha = \frac{\eta J_M + (1 - \eta) J_{SP} - J_{L_2}}{J_{L_1} - J_{L_2}} \tag{A.2}
\]

\[
T_{Cal} = (J_{M_s} - J_{BG_s}) + g \eta e^\tau_s (J_{M_i} - J_{BG_i}) \tag{A.3}
\]

where \( \eta \) is the forward efficiency of the antenna, \( g \) the sideband ratio, \( \tau \) the opacity, and \( J_M, J_{SP}, J_{L_1}, J_{L_2} \) and \( J_{BG} \) are the emissivity temperatures of the average sky, the spill-over, the two loads and the background radiation, respectively. The subscripts \( s \) and \( i \) represent the signal and image bands, respectively. The system temperature is then derived using the formula:

\[
T_{Sys} = T_{Cal} \frac{V_{Sky}}{V_L - V_{Sky}} \tag{A.4}
\]
For ALMA it has been found that with the current system, the non-linearities are the dominant source of error for this calibration. The system electronics and SIS mixers are not fully linear and dominate the relative amplitude calibration accuracy that can be achieved for Cycle 5.

### A.6 Water Vapor Radiometers

In the mm and submm regions, variations in the water vapor distribution in the troposphere that move across an interferometer cause phase fluctuations that degrade the measurements. ALMA uses the so-called "Water Vapor Radiometry" technique to correct for these phase fluctuations. Water Vapor Radiometry involves estimating the excess propagation path amount due to water vapor along a given line-of-sight by measuring the brightness temperature of the sky at frequencies near the atmospheric water vapor resonances. These temperatures can then be transformed into a path length and the difference between any pair of antennas in the array gives the final phase fluctuations to be corrected for a given baseline. ALMA has implemented this technique by placing a Water Vapor Radiometer (WVR) on each 12 m antenna (the 7 m antennas do not have WVRs). For the WVRs to be effective, the measurements have to be taken with a cadence that is fast enough to map the actual variations in the atmosphere. The relevant shortest timescale is the antenna diameter divided by the wind speed as the path delay is averaged over the whole antenna beam and cannot therefore be corrected at any finer time resolution than that. The effective diameter is about 10 m for the ALMA antennas and the relevant windspeed is usually 10 m/s or a bit less so the fastest necessary sampling speed is 1 Hz. On timescales shorter than this 1 Hz timescale, the water vapor path fluctuations are expected to lead to small apparent pointing fluctuations which are analogous to the seeing effects in single-aperture optical telescopes. ALMA selected the 183 GHz line because it is quite bright and allows a more compact design than would the 22 GHz water line. It was decided to measure the temperature of the 183 GHz line in four regions offset from the center using filters of different bandwidths. The positions of the filters are indicated as blue boxes superimposed on the profile of the water vapor line in Figure A.9. The sensitivity specification for the WVRs is 0.08–0.1 K per channel RMS.

![Figure A.9: WVR filters superimposed onto the 183 GHz water vapor emission line.](image)

It is very important that the WVR illuminates the same area of the sky as the ALMA band receivers in the near-field region. This is because the origin of the water vapor fluctuations is usually located in the lower troposphere (i.e., near the observatory), with one to several layers of water vapor clumps encompassing a wide range of sizes. Since the ALMA backends are located at the Cassegrain focus, an offsetting optical system (see Figure A.10) had to be designed to allow the WVR to measure along the optical axis of the antennas.

The WVRs are only able to detect the variations in atmospheric brightness temperatures due to the “wet”
A.6. WATER VAPOR RADIOMETERS

Figure A.10: Offset optics used to collect the sky emission along the optical axis of the antenna into the WVR.

atmosphere (i.e., the precipitable water vapor). There are also variations due to the changes in bulk ambient temperature at different heights above the observatory. It is expected that these could become significant during day time and some techniques are being currently studied to try to measure them (including thermal sounders of the atmosphere that use the profiles of the emission of the oxygen molecules). The brightness temperature variations of the sky that the WVRs have to detect are sometimes quite small, so the quality of the receiving system becomes very important. In fact, the current specification for the ALMA WVRs is that they need to allow corrections of the path fluctuations (in $\mu$m):

$$L_{\text{corr}} \leq \left(1 + \frac{w}{1 \text{mm}}\right)10 \mu m + 0.02\delta L_{\text{raw}}. \quad (A.5)$$

where $w$ is the precipitable water vapor (PWV) content in mm along the line of sight, and $L_{\text{raw}}$ the total fluctuations observed at any given time. Therefore, this formula includes the expected error of about 2% in measuring the total fluctuations, and states the total resulting path errors after correction ($L_{\text{corr}}$). For a 1 mm PWV, the residual term in the formula would be 20 $\mu$m. The stability specification for the WVRs is very stringent (0.1 K peak-to-peak over 10 minutes and 10 degree tilts). To achieve this, a Dicke-switching-radiometer approach was adopted. The input into the mixer is switched periodically (5.35 Hz) between two calibrated loads (the “cold” and “hot” loads at 293 K and 351 K, respectively), and the sky using a rotating vane embedded in the light path as shown in Figure A.11.

Calibration of the measurements is done following the usual method for a 2-load system. The ratios of the output powers when observing the “hot” and “cold” loads can be used to determine the receiver temperatures. Furthermore, these output powers from the loads are also used to extrapolate to a virtual load that has a brightness temperature similar to that of the atmosphere. The specification for the absolute accuracy of the calibration is 2 K (maximum error). The mixer system is an un-cooled DSB Schottky diode pumped by an
LO at 15 GHz that undergoes 2 stages of multiplication. The receiver noise temperature is about 1000 K. After amplification, the IF signal is split into four complete chains (one per filter) and a bandpass filter is applied to select the four desired sampling regions in the profile of the water vapor emission line. In each IF chain, the signal is detected with diodes and after a Voltage-to-Frequency conversion, sent to the Control section for accumulation and control. There is a possibility of LO leakage out of the WVRs that could affect the ALMA receivers in the same antenna and others nearby. To avoid coherence, all the WVRs are tuned to a frequency slightly different (offsets by consecutive integer multiples of 10 kHz up to the total number of WVRs available). The final products sent to the ALMA Control system are time-stamped, calibrated measurements of the brightness temperatures in the 4 filter regions. The path length error due to the PWV can be calculated from these brightness temperature measurements and used to correct the data. Corrections at the scales of the sampling rates of the WVRs are possible at the correlator and refinements for longer timescales can also be done offline in CASA using the wvrgcal tool. Currently both streams of data are being recorded (WVR phase corrected and WVR non-corrected) in order to rigorously assess effectiveness.
Appendix B

The LO and IF System

In this Appendix, the signal path, LO chain, and the relation between these systems and the observer’s spectral setups are described. To the system, a spectral setup effectively consists of the settings of the local oscillators and correlator in the system such that each spectral window (spw) covers the desired lines and/or continuum frequencies. To the end-user, the spectral setup is normally defined in the Observing Tool (OT) just in terms of the observing frequencies and spectral resolutions, and there is no need to worry about the details of each LO setting. For full details of the OT and how to use it, see the OT User and Reference Manuals, available from the ALMA website1 (and also in the OT itself).

The following sections show how the LO system works. For those only interested in the spectral setups and not the details of the components in ALMA, please see Chapter 6.

B.1 Functions of the LO and IF System

In the signal path from the front end to correlator, ALMA uses three frequency conversions. The associated LO and IF systems perform multiple functions:

1. Down-conversion of the sky frequencies to basebands in the range 2–4 GHz, which then alias down to 0–2 GHz for digitization.

2. Amplification and adjustment of the correct power levels into the digitizers.

3. Adjustment of the spw center frequencies (in the Correlator FDM modes) within the basebands. This is actually done in the correlator using the TFB LO, but can effectively be treated as a 4th stage of the LO system.

4. Application of frequency corrections for fringe rotation, and compensation for the slight differences in the Doppler shifts at each antenna due to the differential line-of-sight velocities with respect to the target.


6. Separation of the signal and image sideband. In the case of DSB receivers, selection of the wanted sideband(s). This is done through frequency offsets and phase modulation at each antenna using Walsh patterns.

7. Suppression of spurious signals and reduction of the effects of DC drifts in the samplers. This is done using phase modulation of the LOs using Walsh patterns.

1http://almascience.org/documents-and-tools/
APPENDIX B. THE LO AND IF SYSTEM

Frequency down-conversion therefore effectively occurs in four stages: two hardware Local Oscillators (LO1 and LO2), a 4 GHz sampler/LO and a digital LO synthesised in the tunable filterbanks (TFBs) in the Correlator\(^2\). Section 6 shows how to setup the system to observe spectral lines (particularly multiple spectral lines) and continuum. Some other aspects of frequency setups are then discussed, including the usable bandwidth, spurious signals, and rules and limitations pertaining to this observing Cycle. An overview of the LO and IF operation in ALMA is given in Sections B.2 and in Section B.3, the hardware and how the LO frequencies are synthesised and distributed around ALMA are described.

B.2 Summary of Operation

Figure B.1: Overview of ALMA frequency downconversion, LO mixing and delay corrections. This takes place in the front end, back end, and correlator. Example frequencies are given for an observation at a sky frequency of 100 GHz seen in the USB. Some LOs (e.g. LO1) are continuously tunable; others have quantized tuning steps, such as LO2 (which can be changed in a multiple (“N”) of 125 MHz plus a finely-adjustable offset of “fts”), the TFB LO (which uses a multiple “L” of 30.5 kHz) and the Bulk Delay Correction (which has steps of 250 ps, with a factor of “P”). See text for descriptions of each stage.

Figure B.1 shows a simplified block diagram of the ALMA LO/IF system, showing example setups for an observing frequency centered on 100 GHz. Referring to this diagram, the system operates in the following way:

1. The front-end mixer uses LO1 to downconvert the observing frequency into an IF range covering up to 4-12 GHz. This wide range is needed to cover the IFs of all the ALMA bands, since the mixers for Bands 3, 4, 7 and 8 have an output IF of 4-8 GHz, Band 6 a range of 5-10 GHz and Bands 9 and 10 a range of 4-12 GHz. Over most of the front-end tuning range, LO1 and the front-end mixer can be used in upper or

\(^2\)Note that the ACA correlator is designed to appear like the 64-input Correlator to the end user, although it does not use TFBs in the same way as the BLC.
lower sideband; although at the edges of the tuning band, only one sideband is possible. LO1 consists of a common reference component for all antennas, plus a smaller offset component generated in the FLOOG (First LO Offset Generator) which is different for each antenna (see LO1 Section B.3.4). The FLOOG is used to perform coarse fringe tracking (i.e. rough correction for the small offsets in the observing frequency at each antenna), to offset the LO1 frequencies slightly to suppress internally-generated interference, and for sideband separation or selecting the sideband. It is also used to offset the LO1 phase in conjunction with a Walsh switching pattern on the antennas. A 180 degree phase offset is used to remove spurious signals and DC offset while a 90 degree phase offset is used for sideband suppression and, in the future, for sideband separation on the DSB receivers (Bands 9 and 10).

2. In the back end (BE), the IF processor (IFP) splits the IF into basebands, each with a frequency range of 2-4 GHz, via a set of filters and tunable second LOs (LO2) (see IFS/IFP in Section B.3.5). LO2 is used to offset the individual baseband frequencies within the IF range. The LO2 and second mixer only operates in LSB, with a possible LO2 tuning range of 8-14 GHz. The LO2 signal itself is generated by a course synthesiser which can be set only in steps of 125.0 MHz, plus a second fine-tuned synthesiser (fts) which provides an offset in the range 20.0-42.5 MHz (marked as “fts” in Figure B.1). The limited fts range and the 125 MHz quantization means that LO2 setting is not fully contiguous; consequently there can be up to \( \sim 30 \) MHz difference between the desired and the set value. With a single baseband, this can be compensated by a suitable offset of LO1, but with multiple basebands this is not always possible. So without additional correction, setups with multiple basebands could have the requested lines offset from the spw center by up to 30 MHz. However, the remaining differences in the different spws are compensated by applying an opposite offset to the TFBLOs (LO4 - see below.)\(^3\). An algorithm used by the OT and the realtime system generates the best LO tuning “solution” for LO1, LO2 and LO4 which minimizes the offset of the requested observing frequencies from the centers of the spws. Other uses of LO2 are that the finely-tunable fts is used for fine fringe tracking and LO2 can also be used to offset the frequencies in conjunction with LO1 to suppress interference and select the sideband.

3. The 2-4 GHz analog IF signal from the second mixer in the IFPs is digitised (or sampled) with a 4.0 GHz clock (DGCK). A fine delay (or time) offset is applied to this clock in units of 1/16 of the clock period (250 ps) (the “fine delay correction”).

4. In the FDM correlator mode, up to 32 digital filters (known as TFBs, or "Tunable Filterbanks") are applied to each digitised baseband signal, each of which can be individually adjusted across the baseband frequency (the TFB offsetting). This is effectively applying a digital LO (the TFBLO, or LO4), which is adjustable in steps of 30.517578125 kHz\(^4\) and allows the spectral windows to be moved around within the basebands. At Phase 2, the TFB is centered on the baseband if the TFB “offset” is set to the default of 3000.0 MHz; it can be moved up to +/-900 MHz from that frequency, the range depending on the spw bandwidth. The TFB outputs are resampled and sent to the correlator. The TFBLO can also be used to offset the frequencies in conjunction with LO1 to suppress interference and select the sideband. Finally the correlator software is used to perform the finest level of residual delay correction.

B.3 Frequency Generation and Distribution in ALMA

Figure B.2 shows a summary of the main units involved in the LO generation and distribution. The LOs are generated by the Central LO (CLO) (Section B.3.1) in the AOS Technical Building (lower half of diagram). A fiber-optic system is used to distribute these signals out to the antennas (Section B.3.1) incorporating a realtime path length correction system (Section B.3.3). In the antennas, the important outputs are LO1 (FE 1st LO) (Section B.3.4), LO2 in the IF Processor (Section B.3.5) and the digitizer clock (DGCK). All of these are required in each antenna, shown in the upper part of Figure B.2. In the following subsections, some of these components are described.

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\(^3\)This is done automatically when the OT generates a spectral setup in an SB from a proposal. However, it is repeated at runtime. See [https://safe.nrao.edu/wiki/pub/ALMA/AlmaLamaMemos/lamaMemo808.pdf](https://safe.nrao.edu/wiki/pub/ALMA/AlmaLamaMemos/lamaMemo808.pdf) for more details on the tuning algorithm

\(^4\)The Phase 2 OT has an "adjust" button which quantizes the value entered by the user by this unit
Figure B.2: Summary block diagram of the LO distribution system. The components in the lower section of the diagram (below the dashed-dotted line) form the Central LO (CLO), located in the AOS Technical Building. The components above are located in each Antenna (only one antenna is depicted in this diagram). These are linked by the exterior buried fibers linking the Technical Building with the antenna pads (shown middle-left of the diagram). Thin lines with arrows represent cable distribution, thick lines represent fiber-optic distribution.
B.3. REFERENCE AND LO SIGNAL GENERATION

B.3.1 Reference and LO Signal Generation

The Central Local Oscillator (CLO) generates and distributes the reference, timing and LO signals to all ALMA components to ensure that antenna movement, electronics, and data acquisition are synchronized. These signals are distributed to the antennas through optical fiber using the light of three infrared lasers. The ALMA frequency and phase standard is a Hydrogen maser (installed in 2014), known as the Master Frequency Standard (MFS, lower right), which produces a signal at a frequency of 5 MHz. This is fed into the Central Reference Generator (CRG) module, which produces several signals at multiples of the 5 MHz signal. The 125 MHz signal becomes one of the standards used by many components in the ALMA system. At the AOS Technical building it is used by the Slave Lasers in the Laser Synthesizer modules (Section B.3.2). At the antennas, it is fed into the FLOOG, the Digital Clock (DGCK, see Section B.3.6), the IFPs, and the LO2 Synthesizers. The 2 GHz signal goes into the DTS cards at the antennas. All the reference signals are modulated onto a 1532 nm IR laser in the Central Reference Distributor (CRD) module. The CRD has an internal 48 ms (known as a TE, or Timing Event) clock that is also modulated into the same signal, and is used to synchronize many of the hardware events in the observatory. The modulated 1532 nm signal is sent to an optical distributor (with 80 outputs), the Low Frequency Reference Distributor (LFRD), that feeds it into the Sub Array Switch (SAS) modules, where it is merged with the signals from the Master and Slave lasers (see Section B.3.2). The reference and LO signals are fed to the antennas through the fiber-optic distribution system (see Section B.3.3).

B.3.2 LO Signal Generation and Distribution

LO1 is distributed through fiber-optics and regenerated photonically in each antenna front end by mixing the two infrared laser carriers (known as the Master (ML) and Slave Lasers (SL)) to produce a fixed frequency for all the antennas. Figure B.3 shows a block diagram of this part of the LO system. There are five usable (production grade) Slave Lasers (Laser Synthesizers, LSs), each containing four tunable units, that can produce six different LO1 frequencies which allow simultaneous observations at different subsets of the array (multiple arrays or sub-arrays), and may be used for observing modes that will benefit from rapid switching between frequencies (e.g., band-to-band transfer or spectral scans).

Figure B.3: LO block diagram, showing the Central LO (CLO) and the LO section in the WCA in each front end. Yellow lines represent common synchronization and reference signals fed to all antennas, blue lines are LO signals which are individual to each antenna. For a description of acronyms, see text.

The laser frequencies are generated in the CLO in the following way:

- The Master Laser (ML) generates a 1556 nm fixed optical reference signal, which feeds the Master Laser Distributor (MLD) – essentially a 6-way splitter.
- The Central Reference Generator (CRG) produces reference signals that are fed into the six Laser Synthesizers (LS) fed into the six Central Variable Reference modules, which the LSs use to control the frequency of the Slave Lasers producing a frequency offset of the SL signal of 27-121.7 GHz with respect to the
ML signal. The SL signals are added to the ML signal. The offsets between the ML and SL signals provides the beat note which is used to generate the LO1 frequency in the photomixers in the Warm Cartridge Assemblies (WCAs) in the front end (B.3.4). It is used by the software to set up the front-end observing frequency. With five LSs it is possible to generate five separate LO1 frequencies.

- The Photonic Reference Distribution (PRD) feeds the optical signals to the Sub Array Switch (SAS) module (one per antenna), which chooses the sub array to connect to.

Figure B.4 shows the three laser signals after combination in the Sub Array Switches (SAS). The Master and Slave laser signals have wavelengths of about 1556/1557 nm and the laser carrier signal for the reference signals from the CRD has a wavelength of 1532 nm. The signals are distributed via a single-mode fiber optic line to each of the antennas. The fibers are distributed in buried trenches, and fed into the Cassegrain cabin on each antenna through Az and El fiber wraps. All are fed through Line Length Correctors (LLCs), which are used to correct for changes in the optical fibers. The LLCs are described in B.3.3 below.

Figure B.4: The ALMA fiber signals. The 1532nm carrier contains the frequency reference signals, and the 1556/1555 nm carriers are used to remotely generate LO1. Within each antenna, the optical signals are split and fed to both the LO Reference Receiver (LORR) for the demodulation of the reference/timing signals, and the LO Photonic Receiver for the LO Reference signals.

B.3.3 Line Length Corrections

The LO Reference signals are generated at the AOS Technical Building and need to be distributed via optical fibers to all the antennas. In order to guarantee that the phase of the LO signals is stable during the observations for fibers of up to 15 km in length, compensation for changes has to be done in real time. The method adopted by ALMA is based on a round-trip optical interferometer. Phase fluctuations for an optical fiber transmission system are mainly caused by thermal expansion of the fiber and mechanical stresses, which produce birefringent effects and changes in the absolute polarization of the signals. These changes, in turn, cause differential group propagation delays (DGD) that show up as LO phase jitter. The method implemented by ALMA to correct for this is known as the Line Length Corrector (LLC). Part of the LLC can be seen in Figure B.3, and a more detailed block diagram of the system is shown in Figure B.5.

The two-wavelength laser synthesizer signal (master and slave lasers) is adjusted in polarization, combined at the Laser Synthesizer (LS) and then passed through a 3-port polarizing beam splitter assembly (PBS). The polarization is aligned so that all the light passes through the beamsplitter. It then passes through a piezo-driven fiber stretcher assembly and the fiber to the antenna. At the antenna end there is a 3-dB coupler, so that half of the light goes to the turnaround assembly and half to the photomixer in each WCA. The turnaround assembly consists of a fiber frequency shifter (located at the LO Photonic Receiver module) and a Faraday Rotator mirror located within the WCA of specifically the Band 9 cartridge in each front end. The frequency of the signal traveling back to the AOS technical building receives thus twice a frequency shift of 25 MHz, thus it comes back offset by 50 MHz from the original. The Faraday rotator reflects the signal but turns its polarization.

\[^5\]This figure needs an update. The Slave Laser (variable) wavelength is \(\sim 1557\) nm.
B.3. FREQUENCY GENERATION AND DISTRIBUTION IN ALMA

Figure B.5: Block diagram of the Line Length Corrector system for ALMA. The FRM (Faraday Rotation Mirror), shown upper right, is located in the Band 9 cartridge of every front-end. It rotates the polarisation of the incoming light, and the resulting reflected signal is fed back though the buried optical fiber, to be compared with the outgoing signal. This allows the optical path length to be adjusted in a closed loop using the fiber stretchers.

angle by 90 degrees to the incident polarization. This means that the outgoing and returning light is orthogonal everywhere along the fiber between the PBS and the Faraday Mirror. Back at the PBS, the returning signal is sent to a third port where it is mixed with a sample of the Master Laser reference signal in a low-frequency photodetector. This results in an output at the 50 MHz offset frequency. This output is compared in a phase detector with a 50 MHz reference signal and the phase of the whole loop is kept constant by a servo driving the fiber stretchers.

The current stretchers can cover ranges up to 5 mm in two modes. A “slow” mode (about 10Hz) copes with the large deformations (about 3 mm, allowing for some headroom at the ends of the ranges) and a “fast” response mode (about 1 kHz) copes with the small range variations (about 0.1 mm). The LLCs are reset to mid-range at the start of every SB execution.

B.3.4 The First Local Oscillator (LO1)

The reference signal required to tune LO1 in the receivers is obtained as the difference of the wavelengths of two phase-locked infrared lasers, the Master and Slave lasers. The Master Laser (ML) has a fixed wavelength of 1556 nm and the tunable Slave Laser (SL) is offset from this; both are generated in the CLO (see Section B.3.2). The offset frequency can be anywhere in the range 27 to 38 and 65 to 122 GHz. The beat note from the two lasers constitutes the Photonic LO Reference; the LO1 reference signal is generated from this by photomixers located in the Warm Cartridge Assembly (WCA) of each receiver. This reference signal is used to drive a YIG (Yttrium Iron Garnet) oscillator operating at frequencies around 10–30 GHz (the exact range depending on the band), via a Phase Locked Loop (PLL) circuit. This produces LO1 for the SIS mixers via two sets of multipliers (see example Figure B.6 for Band 7). The same photonic reference signal is distributed to all antennas in the same sub-array. However, to correct for different delay rates required in different antennas, the First LO Offset Generator (FLOOG) in each antenna generates a small but variable (and different) offset frequency in the range 20–45 MHz which is also fed into each PLL. The FLOOGs for all the antennas are continuously tracked during an observation.
B.3.5 LO2 and the IF Processor Units and IF Switch

The output of each front-end cartridge is connected to an IF Switch unit (IFS) situated in the front end, which selects between bands, provides some amplification, and has variable attenuators to set the output levels. The four (or two) outputs from the IF switch unit are fed into two IF Processor units (IFP), one per orthogonal polarization. Figure B.7 shows a basic block diagram of one IF Processor (only one polarisation channel is shown). The Bands 3, 4, 6, 7 and 8 receivers are dual-sideband (2SB), where both the upper and lower sideband signals are provided separately and simultaneously. So there are four outputs from each receiver cartridge in these bands, two sidebands times two polarisations. Each output has an IF bandwidth of up to 4 GHz. For Bands 9 and 10, the receivers are double-sideband (DSB), where the mixer produces a downconverted output from signals in both USB and LSB. These bands have only two outputs, one per polarization, but the signal IF bandwidth of these DSB receivers is 8 GHz per output.

The IF processors divide the incoming IF bands from both sidebands into four 2 GHz basebands and downconvert them to the 2-4 GHz range using the second LO (LO2). Since each baseband is fed by a separate LO2, it is possible to locate them at different frequencies within the IF bandwidth of the receiver (see Chapter 6 and Table 6.2 for limitations). The LO2s are common to both mixer polarizations which means that both polarizations will have the same spectral setups.

The LO2s are digitally-tuned YIG oscillators with a range of 8-14 GHz. LO2 is generated from a harmonic of 125 MHz, plus a fine-tuned synthesiser (fts) of range 20-42.5 MHz, added or subtracted depending on the lock sideband selected by the software. Note that this does not give continuous LO2 coverage, and has to be compensated elsewhere in the LO system.
B.4. OTHER FUNCTIONS OF THE LO/IF

The IFP unit has 0.5 dB stepped attenuators and Total Power detectors for tuning/optimization of the IF power levels into the digital samplers; these levels are set up at the start of each scan. It is important to note that the switch network layout in the IFP means it is NOT possible to select IF configurations with one baseband in one sideband and three in the other (except for DSB receivers, where this is done using sideband selection). The IFP has anti-alias filters, one set of which is switchable depending on whether the IF range in use is in the upper or lower part of the IF band. As well as downconversion, the LO2s can also be used for sideband separation when combined with the first LO (Section B.1).

Figure B.7: Block diagram of one polarisation channel of the IF Processor. This has two IF inputs (at left), and feeds 4 baseband (BB) outputs to the digitisers (to the right of the diagram).

The IF processor also has anti-aliasing filters, which define the 2 GHz baseband width and remove out-of-band signals (Section B.3.6). This results in the higher noise levels on the upper and lower 50-100 MHz of channels in the TDM correlator mode (see Section 6.4). These filters cause a decrease in the effective IF range to approximately 1.875 GHz.

B.3.6 Digitization and Signal Transmission

The outputs of the IF Processor units are fed into the Data Transmission System Transmitter (DTX), that include digitizers and formatters to convert the signals to optical wavelengths for transmission via optical fibers. There are four DTX units per antenna, each one handling data for a given baseband pair (i.e., the same 2 GHz baseband from each of the two orthogonal polarizations). Each baseband is digitized by a separate digitizer at 4 GHz (i.e., Nyquist sampling for a 2 GHz bandwidth), quantizing each sample into 3 bits (8 levels) per polarization, so that a total of 6 bits must be transferred per baseband pair. The digitized signal is then transferred to the formatter part that packages the data in frames of equal size. The output of each DTX module is fed to three optical fibers, each transporting 2 bits, and the signal leaves the antenna after passing through a Fiber Optic Multiplexer (FOM). All DTX modules are fed with reference/timing signals from an associated Digital Clock (DGCK), which is also used to do the fine delay tracking.

The outputs of the DTX are sent, via the optical fibers, to the AOS Technical Building where the process is inverted (conversion from optical to digital signal) at the DRXs (Data Transmission System Receiver modules), before the signals are sent to the correlator. Delay corrections due to changes in the length of the optical fibers are done using metadata information to realign the frames sent from the transmitting side at the antenna (DTX) and the receiving side at the Technical Building (DRX). Figure B.8 shows a block diagram of a single DTX module.

B.4 Other functions of the LO/IF

In addition to frequency downconversion, the LO/IF performs several other tasks, detailed below.
B.4.1 Delay Corrections

ALMA handles delay corrections via the “Delay Server” software package. It computes the corrections for all the different components involved with a cadence of one minute and distributes them buffered. The three main components along the data flow chain where the corrections are applied are: the First LO Offset Generator (FLOOG), the Digital Clock (DGCK) and the correlator (see Figure B.1). Fringe tracking is done at the FLOOG by slightly offsetting the frequency of the LO1 signal. Currently, the delay handled by the FLOOG is in steps of 250 ps. The FLOOG is also used for phase and frequency switching for suppression and separation of sidebands, and for rejection of internally-generated interference, described in the next subsections.

Fine delay corrections are handled by the DGCK that feeds the corrections into the four DTS modules in each antenna. The delay correction resolution of the fine delay is 1/16 of the 4GHz ADC clock. The bulk delay correction is handled by the Correlator in integer multiples of the 250 ps units. On top of these corrections, the correlator also handles the “residual” delay corrections at much higher temporal resolution (<250 ps/16) by applying a linear phase gradient across the passband after correlation. Also, the correlator applies relative delay corrections between all the basebands and polarizations of a given ALMA band receiver. Currently, the first baseband of the X polarization is used as reference.

B.4.2 Sideband Suppression - LO Offseting

Some of the ALMA receivers (e.g. Bands 3, 4, 6, 7, and 8) are inherently single sideband (SSB), either through having a mixer or quasioptic design which rejects the unwanted sideband. Their intrinsic sideband rejection is typically only about 10-15 dB, which, although adequate for rejection of the unwanted sky noise, is not enough to remove strong lines from the other (image) sideband. Others receivers (Bands 9, 10) are double sideband (DSB), and the relative response of the two sidebands may not be equal, significantly affecting calibration. Accordingly, additional schemes are necessary for more effective removal of the unwanted sideband (known as

...Although they have mixers to allow both sidebands to be observed separately and simultaneously...
sideband suppression), and for correlation of both sidebands independently (or sideband separation - see next section). Sideband suppression in ALMA is done using the FLOOG, and either LO2 (2LO offset) or a combination of LO2 and LO4 (3LO offsetting). A small frequency offset \( F_o \) is added to LO1 and subtracted from the other LOs, so that while the signal sideband remains at the same frequency, the image sideband is shifted \( 2F_o \) away from its nominal value. A different value of \( F_o \) is applied at every antenna (the offsets are defined using a Walsh pattern), so that all signal sidebands are at the same frequency, but all image sidebands are at slightly different frequencies and no longer correlate.

Note that each of the basebands has an independent LO2 and LO4. So by setting the sign of the offset in (LO2+LO4) differently, each baseband can be set up to observe in a different sideband.

For single-dish observing, such interferometric sideband rejection methods cannot be used, and a frequency scanning method is under development which will allow image rejection.

B.4.3 Interference Rejection - 180 degree Phase Switching

The FLOOG is additionally used to reject spurious signals prior to digitization by applying 180 deg phase switching according to orthogonal Walsh function patterns, with a pattern cycle time of 16 ms. The Walsh pattern is different on each antenna, and is demodulated by a sign change within the DTS; as a result, the wanted signals correlate, and the unwanted signals are canceled out. This rejects spurious signals generated in the system between the receiver and the sampler, and also suppresses sampler DC offsets. 180-Walsh is a default setup for all observations. For further details please see Section 5.5.4.

B.4.4 Sideband Separation - 90 degree Walsh Switching

It is possible to apply a 90 deg phase switch in the FLOOG and in the correlator processing, allowing correlation of the upper and lower sidebands separately. For Bands 9 and 10 (DSB) it will effectively double the bandwidth from 8 GHz to 16 GHz per polarization from the DSB receivers. For further details please see Section 5.5.4.
Appendix C

Acronym Dictionary

2SB Dual Sideband
AIA Atmospheric Imaging Assembly (SDO)
ACA Atacama Compact Array
ACD Amplitude Calibration Device
ACS ALMA Common Software
AFA ALMA Frontend Archive
ALMA Atacama Large Millimeter/Submillimeter Array
AoD Astronomer on Duty
AOS Array Operation Site
APDM ALMA Project Data Model
APE ALMA Production Environment
APP ALMA Phasing Project
APS ALMA Phasing System
AQUA ALMA Quality Assurance software
ARC ALMA Regional Center
ASA ALMA Science Archive
ASC ALMA Sensitivity Calculator
ASDM ALMA Science Data Model
ASIC Application Specific Integrated Circuit
ATC Advanced Technology Center, NAOJ
ATM Atmospheric Transmission at Microwaves
AZ Azimuth
BB Baseband
BDF Binary Data Format
BE Backend
BL Baseline
BLC BaseLine Correlator
BWFN Beam Width between First Nulls
BWSW Bandwidth Switching
CASA Common Astronomy Software Applications package
CCA Cold Cartridge Assemblies
CCC Correlator Control Computer
CDP Correlator Data Processor
CFRP Carbon Fiber Reinforced Plastic
CIP Correlator and Integration Processor
CLO Central Local Oscillator
CLT Chilean Local Time
<table>
<thead>
<tr>
<th>Acronym</th>
<th>Description</th>
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<tbody>
<tr>
<td>CORBA</td>
<td>Common Object Request Broker Architecture</td>
</tr>
<tr>
<td>CRD</td>
<td>Central Reference Distributor</td>
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<tr>
<td>CRG</td>
<td>Central Reference Generator</td>
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<tr>
<td>CSV</td>
<td>Commissioning and Science Verification</td>
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<tr>
<td>CW</td>
<td>Continuous Wave</td>
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<tr>
<td>DC</td>
<td>Direct Current</td>
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<tr>
<td>DEC</td>
<td>Declination</td>
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<tr>
<td>DFP</td>
<td>DTS-Rx and FFT Processor</td>
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<tr>
<td>DGC</td>
<td>Differential Gain Calibrator</td>
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<tr>
<td>DGCK</td>
<td>Digital Clock</td>
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<tr>
<td>DMG</td>
<td>Data Management Group within DSO</td>
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<td>DRX</td>
<td>Data Receiver module</td>
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<tr>
<td>DSB</td>
<td>Double Sideband</td>
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<tr>
<td>DSO</td>
<td>Division of Science Operations</td>
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<tr>
<td>DTS</td>
<td>Data Transmission System</td>
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<tr>
<td>DTS-Rx</td>
<td>Data Transmission System Receiver</td>
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<td>DTX</td>
<td>Data Transmitter module</td>
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<tr>
<td>EB</td>
<td>Execution Block</td>
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<tr>
<td>EHT</td>
<td>Event Horizon Telescope</td>
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<tr>
<td>EL</td>
<td>Elevation</td>
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<tr>
<td>EM</td>
<td>Electromagnetic</td>
</tr>
<tr>
<td>EPO</td>
<td>Education and Public Outreach</td>
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<tr>
<td>ES</td>
<td>Early Science</td>
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<tr>
<td>ESO</td>
<td>European Southern Observatory</td>
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<tr>
<td>FDM</td>
<td>Frequency Division Mode</td>
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<tr>
<td>FTP</td>
<td>Fourier Transform</td>
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<tr>
<td>FWHM</td>
<td>Full Width Half Maximum</td>
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<tr>
<td>FWHP</td>
<td>Full Width to Half Power</td>
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<tr>
<td>FWBN</td>
<td>Full Width Between the Nulls</td>
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<tr>
<td>FX</td>
<td>Fourier transform correlation type correlator</td>
</tr>
<tr>
<td>FXF</td>
<td>Filtering, Correlation, and Fourier transform type correlator</td>
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<tr>
<td>GMVA</td>
<td>Global Millimeter VLBI Array</td>
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<tr>
<td>GOUS</td>
<td>Group Observing Unit Set</td>
</tr>
<tr>
<td>GPS</td>
<td>Global Positioning System</td>
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<tr>
<td>GUI</td>
<td>Graphical User Interface</td>
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<tr>
<td>HA</td>
<td>Hour Angle</td>
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<tr>
<td>HEMT</td>
<td>High Electron Mobility Transistor</td>
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<tr>
<td>HPBW</td>
<td>Half Power Beam Width</td>
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<tr>
<td>IF</td>
<td>Intermediate Frequency</td>
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<tr>
<td>IFP</td>
<td>Intermediate Frequency Processor</td>
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<tr>
<td>IRAM</td>
<td>Institut de Radioastronomie Millimetrique</td>
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<tr>
<td>JAO</td>
<td>Joint ALMA Observatory</td>
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<tr>
<td>JCMT</td>
<td>James Clerk Maxwell Telescope</td>
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<tr>
<td>JPL</td>
<td>Jet Propulsion Laboratory</td>
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<tr>
<td>Abbreviation</td>
<td>Description</td>
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<tr>
<td>LAS</td>
<td>Largest Angular Resolution</td>
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<td>LFRD</td>
<td>Low Frequency Reference Distributor</td>
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<tr>
<td>LLC</td>
<td>Line Length Corrector</td>
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<tr>
<td>LO</td>
<td>Local Oscillator</td>
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<tr>
<td>LO1</td>
<td>First LO</td>
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<tr>
<td>LO2</td>
<td>Second LO</td>
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<tr>
<td>LO3</td>
<td>Digitizer Clock LO, Third LO</td>
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<tr>
<td>LO4</td>
<td>Tunable Filterbanks LO, Fourth LO</td>
</tr>
<tr>
<td>LORR</td>
<td>LO Reference Receiver</td>
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<tr>
<td>LS</td>
<td>Laser Synthesizer</td>
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<td>LSB</td>
<td>Lower Sideband</td>
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<tr>
<td>LSRK</td>
<td>Local Kinematics Standard of Rest</td>
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<tr>
<td>LTA</td>
<td>Long Term Accumulator</td>
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<tr>
<td>MCI</td>
<td>Monitor and Control Interface</td>
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<tr>
<td>MD</td>
<td>Mixed De-Tuned or Mixed De-Biased mode</td>
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<tr>
<td>MFS</td>
<td>Master Frequency Standard</td>
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<tr>
<td>MIT</td>
<td>Massachusetts Institute of Technology</td>
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<tr>
<td>ML</td>
<td>Master Laser</td>
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<td>MLD</td>
<td>Master Laser Distributor</td>
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<td>MOUS</td>
<td>Member Observing Unit Set</td>
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<td>MRS</td>
<td>Maximum Recoverable Scale</td>
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<td>MS</td>
<td>Measurement Set</td>
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<td>NAOJ</td>
<td>National Astronomy Observatory Japan</td>
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<td>NGAS</td>
<td>New Generation Archive System</td>
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<td>NRAO</td>
<td>National Radio Astronomy Observatory</td>
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<tr>
<td>NVSS</td>
<td>NRAO VLA Sky Survey</td>
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<tr>
<td>OMC</td>
<td>Operator Monitoring and Control</td>
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<tr>
<td>OMT</td>
<td>Ortho-mode Transducer</td>
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<tr>
<td>OSF</td>
<td>Operations Support Facility</td>
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<tr>
<td>OST</td>
<td>Observation Support Tool</td>
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<tr>
<td>OT</td>
<td>Observing Tool</td>
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<tr>
<td>OTF</td>
<td>On the Fly</td>
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<tr>
<td>OUS</td>
<td>Observing Unit Set</td>
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<tr>
<td>PBS</td>
<td>Polarization Beam Splitter</td>
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<tr>
<td>PDM</td>
<td>Propagation Delay Measure</td>
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<tr>
<td>PI</td>
<td>Principal Investigator</td>
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<tr>
<td>PLL</td>
<td>Phase Lock Loop</td>
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<td>PMG</td>
<td>Program Management Group within DSO</td>
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<tr>
<td>PRD</td>
<td>Photonic Reference Distribution</td>
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<tr>
<td>PSF</td>
<td>Point Spread Function</td>
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<tr>
<td>PWV</td>
<td>Precipitable Water Vapor</td>
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<tr>
<td>QA</td>
<td>Quality Assurance</td>
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<td>QA0</td>
<td>Quality Assurance Level 0</td>
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<td>Quality Assurance Level 1</td>
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<td>Quality Assurance Level 2</td>
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<td>Full Form</td>
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<td>SDO</td>
<td>Solar Dynamics Observatory</td>
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<td>SED</td>
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<td>Source Equivalent Flux Density</td>
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<td>SL</td>
<td>Slave Laser</td>
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<td>Signal-to-Noise Ratio</td>
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<td>Spectral Window</td>
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<td>spws</td>
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<td>Square Law Detector</td>
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<td>Very Long Baseline Interferometry</td>
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<td>VOM</td>
<td>VLBI Observing Mode</td>
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<td>VO</td>
<td>Virtual Observatory</td>
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<td>WCA</td>
<td>Warm Cartridge Assembly</td>
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<td>WVR</td>
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<tr>
<td>XF</td>
<td>Correlation-Fourier Transform Type Correlator</td>
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<tr>
<td>XML</td>
<td>eXtensible Markup Language</td>
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<tr>
<td>YIG</td>
<td>Yttrium-Iron Garnet Oscillator</td>
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